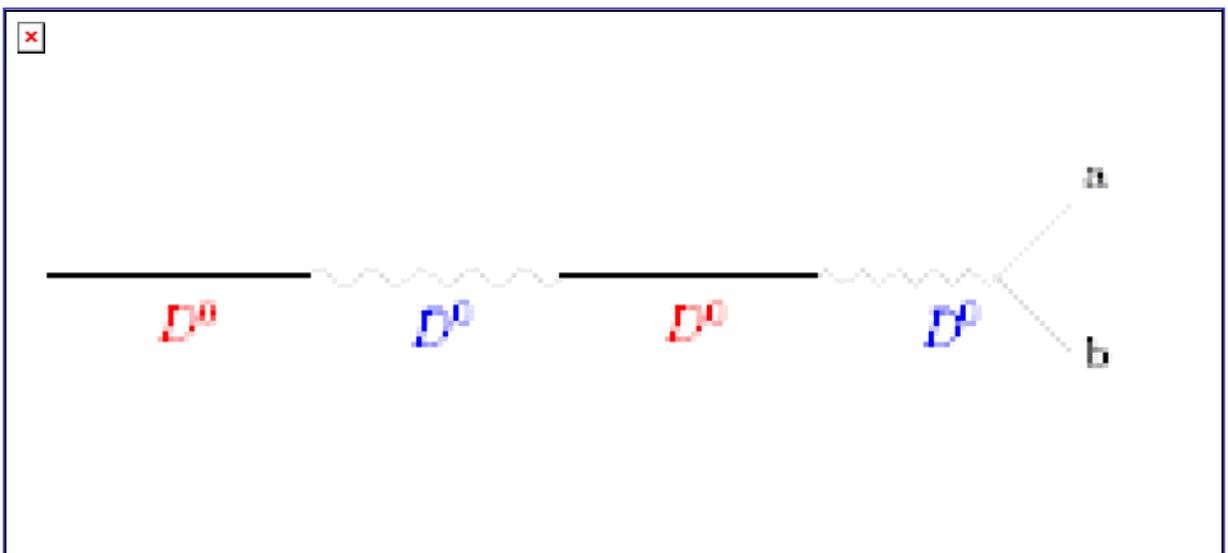


# Charm Mixing - Theory

[CP-Conserving only]

Gene Golowich  
Univ. of Massachusetts

CHARM-07  
Cornell University  
5-8 August 2007



# Our Thoughts on $\gamma_D$ and $x_D$

This talk will refer to:

1]  $\gamma_D$  : PRL 98 (2007) 181808-1

New Physics contributions to the lifetime difference in  $D^0$ - $\bar{D}^0$  mixing

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2]  $x_D$  : arXiv:0705.3650 [hep-ph]

Implications of  $D^0$ - $\bar{D}^0$  Mixing for New Physics

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# Status of $D^0$ Mixing

At the time of our paper on  $x_D$ :

$$x = \frac{\Delta M}{\Gamma} \quad x_D = (8.7^{+3.0}_{-3.4}) \cdot 10^{-3}$$

$$y = \frac{\Gamma_1 - \Gamma_2}{2\Gamma} \quad y_D = (6.6 \pm 2.1) \cdot 10^{-3}$$

The above  $x_D$  is a **2.4  $\sigma$**  effect.

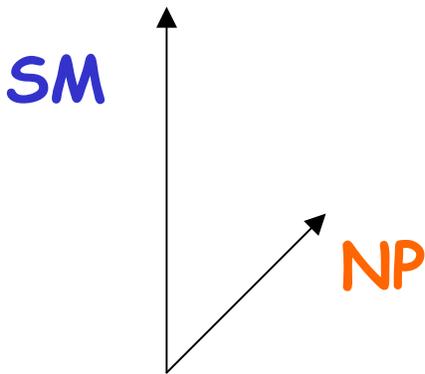
PRL discovery criteria are:

- a) 'Observation':  **$>5\sigma$**
- b) 'Evidence':  **$3\sigma$ -to- $5\sigma$**
- c) 'Measurement':  **$<3\sigma$**

## Basic Strategy for $\chi_D$

Observed Signal at roughly 1% level.  
To some, this is 'large' for SM.  
Or is it?

Our premise is to study **both** SM, NP.  
We do not know the relative phase.



So we compare pure NP signals to  
values  $\chi_D = (3.0 \rightarrow 15.0) \cdot 10^{-3}$

# Standard Model

## Quark-level Analysis

Operator Product Expansion

QCD Perturbation Theory

Expansion in  $m_s/m_c$

Evaluation of B-parameters

## Hadron-level Analysis

Focus on  $\gamma_D$

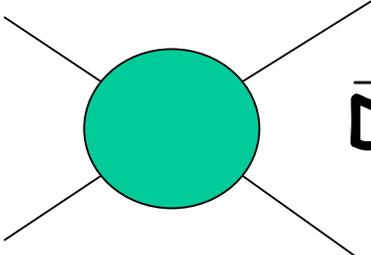
Direct Involvement of Data/Models

Role of SU(3) Breaking

Possible Large Effect

# Charm Mixing and the OPE\*

Expand in increasing operator dimension:

$$D^0 \quad \text{---} \quad \text{---} \quad \text{---} \quad \text{---} \quad \bar{D}^0 = \sum_n O_n^{(d=6)}$$

$$+ \sum_n O_n^{(d=9)} + \dots$$

**D=6:** Two local 4F operators

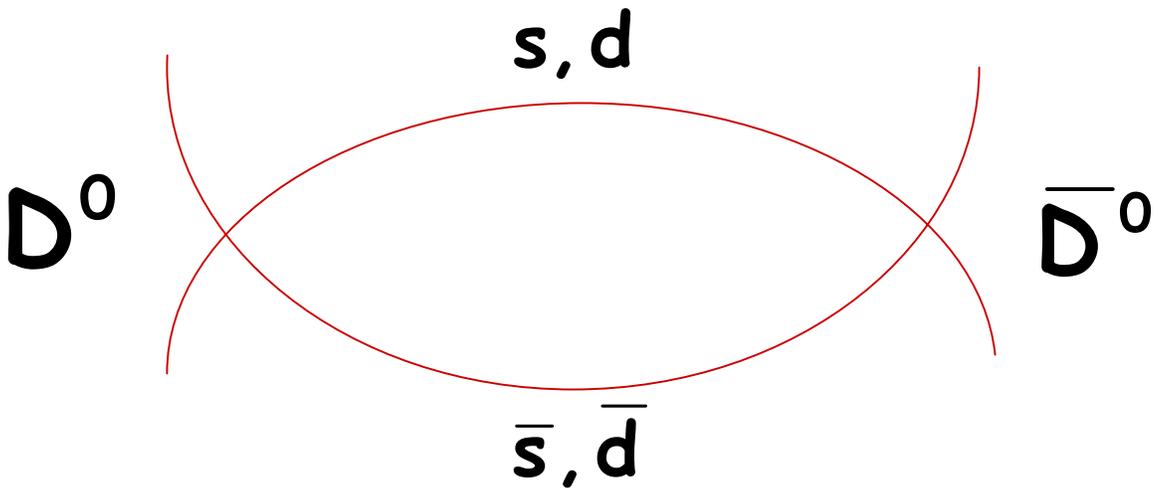
**D=9:** Fifteen local 6F operators

Etc

\*[Georgi PL B297 (1992) 353]

## Dimension Six

Ignore b quark. Sum over  $s\bar{s}, d\bar{d}, s\bar{d} + d\bar{s}$  intermediate states.



Expand in powers of

$$z = \frac{m_s^2}{m_c^2} \approx 0.006$$

# Flavor Cancellations

$\Delta\Gamma$  at  $d=6$  ( $m_d=0$ ):

$z^0$        $z^1$        $z^2$

$$s\bar{s} \quad \frac{1}{2}$$

$$d\bar{d}$$

$$s\bar{d} + d\bar{s}$$

**Total**

# Flavor Cancellations

$\Delta\Gamma$  at  $d=6$  ( $m_d=0$ ):

$z^0$        $z^1$        $z^2$

$$s\bar{s} \quad \frac{1}{2}$$

$$d\bar{d} \quad \frac{1}{2}$$

$$s\bar{d} + d\bar{s}$$

**Total**

# Flavor Cancellations

$\Delta\Gamma$  at  $d=6$  ( $m_d=0$ ):

$z^0$        $z^1$        $z^2$

$$s\bar{s} \quad \frac{1}{2}$$

$$d\bar{d} \quad \frac{1}{2}$$

$$s\bar{d} + d\bar{s} \quad -1$$

$$\text{Total} \quad 0$$

# Flavor Cancellations

$\Delta\Gamma$  at  $d=6$  ( $m_d=0$ ):

	$z^0$	$z^1$	$z^2$
$s\bar{s}$	$\frac{1}{2}$	$-3z$	
$d\bar{d}$	$\frac{1}{2}$	$0$	
$s\bar{d} + d\bar{s}$	$-1$	$3z$	
<b>Total</b>	<b>0</b>	<b>0</b>	

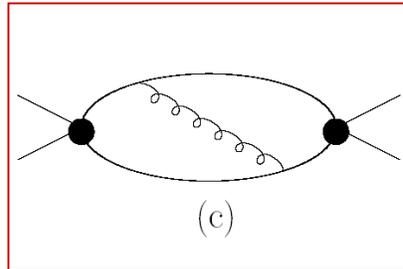
# Flavor Cancellations

$\Delta\Gamma$  at  $d=6$  ( $m_d=0$ ):

	$z^0$	$z^1$	$z^2$
$s\bar{s}$	$\frac{1}{2}$	$-3z$	$3z^2$
$d\bar{d}$	$\frac{1}{2}$	$0$	$0$
$s\bar{d} + d\bar{s}$	$-1$	$3z$	$-3z^2$
<b>Total</b>	<b>0</b>	<b>0</b>	<b>0</b>

# Allowing for QCD\*

Expand in  $\alpha_s$ :



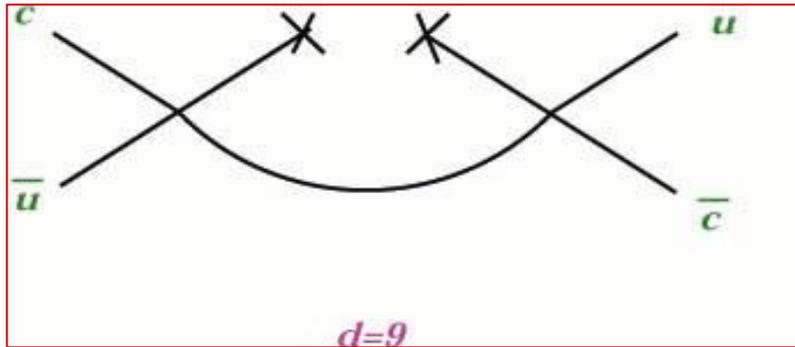
	$x$	$y$	Comment
$\alpha_s^0$ (LO)	$z^2$	$z^3$	$x^{(LO)} \gg y^{(LO)}$
$\alpha_s^1$ (NLO)	$z^2$	$z^2$	$x^{(NLO)} > y^{(NLO)}$

Main LO + NLO Result:  $x \cong y \approx 10^{-6}$   
(And find NLO > LO)

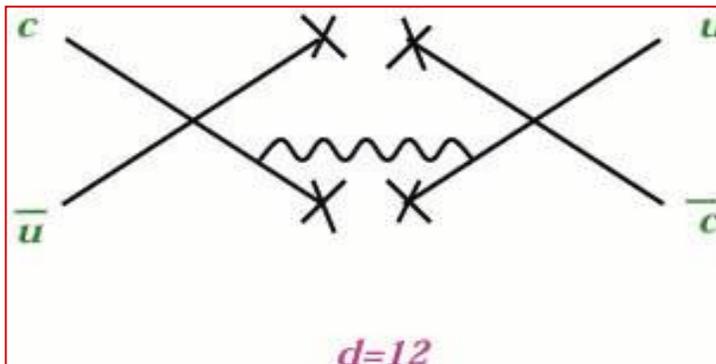
\*EG & Petrov PLB 625 (2005) 53

# Higher Terms in the OPE\*

$$D = 9 \ (\propto z^{3/2})$$



$$D = 12 \ (\propto z)$$



\*[Ohl, Ricciardi & Simmons NP B403 (1993) 605]

# Quark-Level Summary

## Triple Expansion:

1. Operator dimensions  $d = 6, 9, 12, \dots$
2. QCD factors  $\alpha_s/4\pi$
3. Mass ratio  $z = (m_s/m_c)^2$

## Status:

1. Scale thus far  $\ll 1\%$
2. Scale if all terms included, unknown.

Keep trying!

Bigi & Uraltsev, NP B592 (2001) 92

Gagik, Golowich, Petrov (in progress)

## Hadron-level ( $\Delta\Gamma$ )

$$\gamma_D = \frac{1}{2M_D\Gamma_D} \text{Im} \langle \bar{D}^0 | i \int d^4x T(H_w(x)H_w(0)) | D^0 \rangle$$

Insert hadronic int. states:  $\sum_n |n\rangle\langle n|$

Require matrix elements  $\langle n | H_w | D^0 \rangle$

**1. Use a model:**  $\gamma_D \sim 10^{-3}$

Naples Group, PRD 51 (1995) 3478

# Hadron-level ( $\Delta\Gamma$ ) cont.

## 2. Use data

(a) Early Work [UMass PRD 33 (1985) 178]

Choose  $n = P^+P^-$

SU(3) Limit: Zero via cancellation

SU(3) breaking important?

Preliminary finding: ' $\gamma_D$  large'

(b) Recent Work [FGLNP PRD 69 (2004) 114021]

Theorem: SU(3) breaking 2<sup>nd</sup> order

So maybe SU(3) breaking not large

But 4P sector cannot cancel.

Conclude ' $\gamma_D \sim 10^{-2}$  possible'

Quite possibly correct.

More persuasive than compelling

Uncontrollable uncertainties,

# New Physics in $D^0$ Mixing

$Y_D$

Intermediate states **on-shell**.

Thus only light particles propagate.

Can there be **any** NP effects?

Derive a 'master formula'.

$X_D$

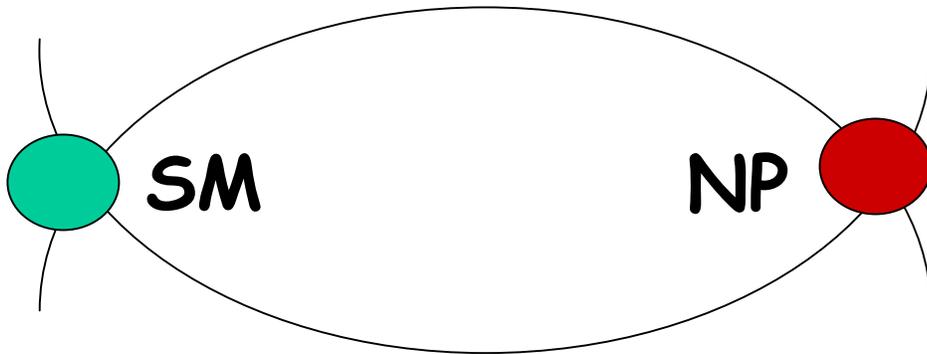
Intermediate states off-shell.

**Many** possible NP candidates.

Which one to consider?

How to organize?

# New Physics and $\Delta\Gamma$



NP **can** affect  $\Delta\Gamma$ !

Via the  $\Delta C = \pm 1$  interaction vertex.

Processes like  $c\bar{u} \rightarrow q_1\bar{q}_2$

Golowich, Pakvasa, Petrov, PRL 98 (2007) 181801.

[Comment]: Chen, Geng, Nam, PRL 99 (2007) 019101.

[Comment<sup>2</sup>]: Yeghiyan, arXiv 0707.3285 [hep-ph].

# The Master Formula

INPUT

$$\mathbf{H}_{NP}^{\Delta C=-1} = \sum_{q,q'} \mathbf{D}_{qq'} [\bar{\mathcal{C}}_1(\mu) \mathbf{Q}_1 + \bar{\mathcal{C}}_2(\mu) \mathbf{Q}_2]$$

$$\mathbf{Q}_1 = \bar{u}_i \bar{\Gamma}_1 q'_j \bar{q}_j \bar{\Gamma}_2 c_i \quad \mathbf{Q}_2 = \bar{u}_i \bar{\Gamma}_1 q'_i \bar{q}_j \bar{\Gamma}_2 c_j$$

OUTPUT

$$\begin{aligned} y_D = & - \frac{4\sqrt{2}G_F}{M_D \Gamma_D} \sum_{q,q'} \mathbf{V}_{cq'}^* \mathbf{V}_{uq} \mathbf{D}_{qq'} (\mathbf{K}_1 \delta_{ik} \delta_{jl} \\ & + \mathbf{K}_2 \delta_{il} \delta_{jk}) \sum_{\alpha=1}^5 \mathbf{I}_\alpha(x, x') \langle \bar{D}^0 | \mathbf{O}_\alpha^{ijkl} | D^0 \rangle \end{aligned}$$

## Some Results for $y_D$

Model	$y_D$	Comment
RPV-SUSY	$6 \cdot 10^{-6}$	Squark Exch.
	$-4 \cdot 10^{-2}$	Slepton Exch.
Left-right	$-5 \cdot 10^{-6}$	'Manifest'.
	$-8.8 \cdot 10^{-5}$	'Nonmanifest'.
Multi-Higgs	$2 \cdot 10^{-10}$	Charged Higgs
Extra Quarks -	$10^{-8}$	Not Little Higgs

# New Physics and $x_D$

As the LHC era begins, many extras possible (21 models in GHPP)\*

- Extra **gauge bosons**  
(LR models, etc)
- Extra **scalars**  
(Multi-Higgs models, etc)
- Extra **fermions**  
(Little Higgs, etc)
- Extra **dimensions**  
(Universal extra dimensions, etc)
- Extra **global symmetries**  
(SUSY, etc)

\*GHPP: arXiv 0705.3650 [hep-ph]

# List of NP Models

Fourth Generation

$Q=-1/3$  Singlet Quark

$Q=+2/3$  Singlet Quark

Little Higgs

Generic  $Z'$

Family Symmetries

Left-Right Symmetries

Alternate L-R Symmetries

Vector Leptoquark Bosons

Fl-Cons Two-Higgs Dbt

Fl-Chnge Neutral Higgs I

Fl-Chnge Neutral Higgs II

Scalar Leptoquark Bosons

Higgless

Universal Extra Dims

Split Fermion

Warped Geometries

Minimal SUSY Standard

SUSY Alignment

SUSY with RPV

Split SUSY

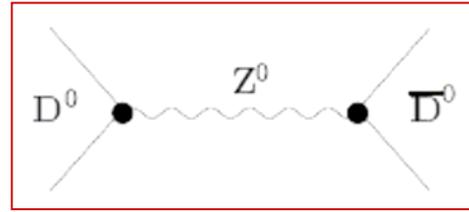
## Challenge to the Audience\*

Of the **21** NP models, how many turn out to yield contributions **too small** for  $D^0$  mixing at the observed  $10^{-2}$  level?

**Comment:** Note many NP models have been on the market for years (e.g., SUSY has been studied for over 30 yrs) and their parameter spaces have been steadily constrained.

\* If you have already seen the paper, please keep quiet.

# A NP Example



## Diagram

$Z^0$  tree amplitude

Two flavor-changing vertices

## Realizations:

Vector-like  $SU(2)$ -singlet quarks

$E(6)$ :  $Q = -e/3$

Little Higgs:  $Q = 2e/3$

**Find** 
$$x_D = \frac{G_F \lambda_{uc}^2}{\sqrt{2} M_D \Gamma_D} C_1(m_c) \langle \bar{D}^0 | Q_1 | D^0 \rangle$$

# Operator Basis

$$\begin{aligned} Q_1 &= \bar{u}_L \gamma_\mu c_L \bar{u}_L \gamma^\mu c_L \\ Q_2 &= \bar{u}_L \gamma_\mu c_L \bar{u}_R \gamma^\mu c_R \\ Q_3 &= \bar{u}_L c_R \bar{u}_R c_L \\ Q_4 &= \bar{u}_R c_L \bar{u}_R c_L \\ Q_5 &= \bar{u}_R \sigma_{\mu\nu} c_L \bar{u}_R \sigma^{\mu\nu} c_L \\ Q_6 &= \bar{u}_R \gamma_\mu c_R \bar{u}_R \gamma^\mu c_R \\ Q_7 &= \bar{u}_L c_R \bar{u}_L c_R \\ Q_8 &= \bar{u}_L \sigma_{\mu\nu} c_R \bar{u}_L \sigma^{\mu\nu} c_R \end{aligned}$$

# Operator Matrix Elements

In vacuum saturation, just **two** nonperturbative constants ( $B, B'_S$ )

$$\langle Q_1 \rangle = \frac{2}{3} f_D^2 M_D^2 B$$

$$\langle Q_2 \rangle = -\frac{1}{2} f_D^2 M_D^2 B - \frac{1}{3} f_D^2 M_D^2 B'_S$$

$$\langle Q_3 \rangle = \frac{1}{12} f_D^2 M_D^2 B + \frac{1}{2} f_D^2 M_D^2 B'_S$$

$$\langle Q_4 \rangle = -\frac{5}{12} f_D^2 M_D^2 B'_S$$

$$\langle Q_5 \rangle = f_D^2 M_D^2 B'_S$$

$$\langle Q_6 \rangle = \langle Q_1 \rangle$$

$$\langle Q_7 \rangle = \langle Q_4 \rangle$$

$$\langle Q_8 \rangle = \langle Q_5 \rangle$$

# RG Factor

Ex:  $Q_6 = \bar{u}_R \gamma^\mu c_R \bar{u}_R \gamma_\mu c_R$

Two scales:  $M \gg m_c$

Have  $C_6(M)$                       Need  $C_6(m_c)$

Integrate RG equation

Obtain  $C_6(m_c) = R[M, m_c] C_6(M)$

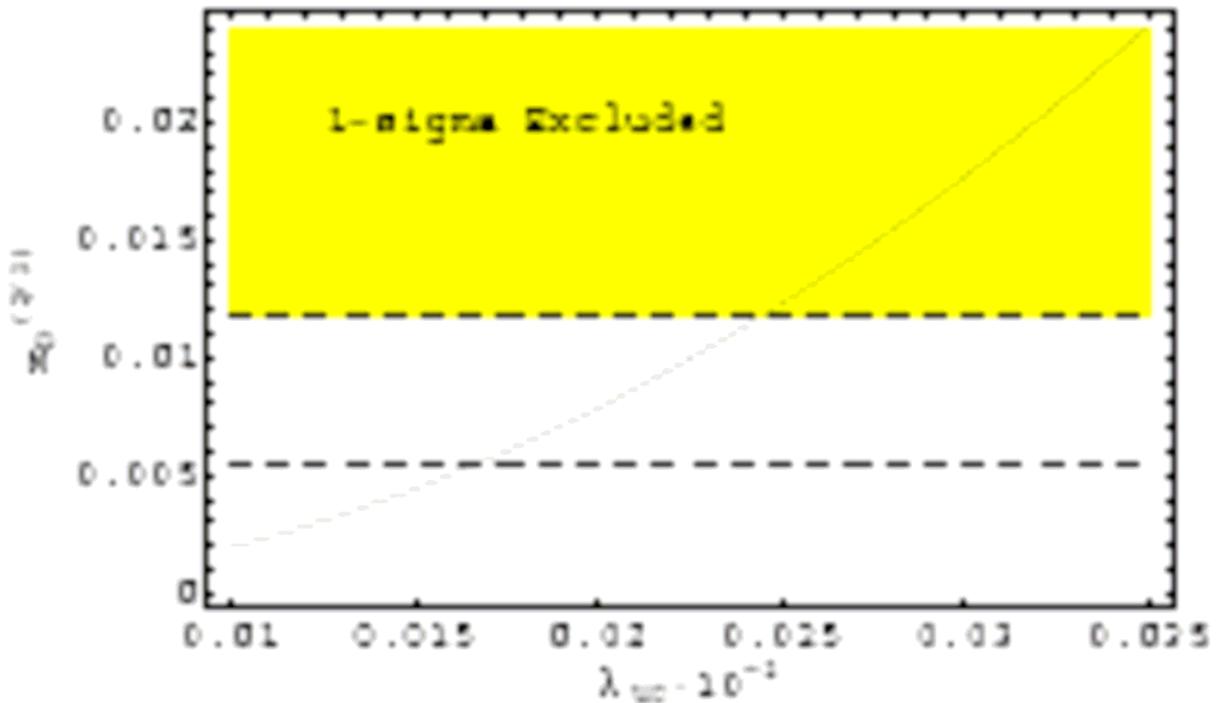
with

$$R[M, m_c] = r^{2/7}(M, m_t) r^{6/23}(m_t, m_b) r^{6/25}(m_b, m_c)$$

$$r(m_1, m_2) = a_s(m_1) / a_s(m_2)$$

# Limit on $Q=2/3$ Quark Singlet

Plot  $x_D^{(2/3)}$  vs  $\lambda_{uc}$ :



$D^0$ -mixing limit on  $\lambda_{uc}$  is about  $10^2$  better than that from  $4 \times 4$  CKM unitarity.

# Answer to Challenge

## Ineffective Models:

**Four** yield no constraints:

1. Split supersymmetry
2. Universal Extra Dimensions
3. Left-right symmetric
4. FC two-Higgs doublet

## Constrainable Models:

There are **17** which can, in principle, exceed the observed  $x_D$ . For these, we can get constraints on masses and mixing parameters.

# Split SUSY - Why So Small?

## What 'is' Split SUSY?:

- New variant of SUSY (2003-4)
- SUSY breaks at  $m_s \gg 1000 \text{ TeV}$
- Scalars (except Higgs) have mass  $\sim m_s$
- Fermions have usual weak scale mass

## Why So Small in $D^0$ Mixing?:

Large  $D^0$  mixing in SUSY involves squark (i.e. scalar quarks) amplitudes. But squark masses are **huge** in Split SUSY. Thus the mixing is suppressed.

# UEDs - Why So Small?

## What 'are' Universal Extra Dimensions?

- Variant (2000) of having  $\text{TeV}^{-1}$ -sized extra dimensions
- No branes in this approach
- All SM fields reside in the bulk
- Usually one extra dimension

## Why So Small in $D^0$ Mixing?:

Each SM field has an infinity of KK excitations. **GIM cancellations** affect all save a few b-quark KK terms, but these are CKM suppressed.

# Results of $x_D$ Analysis

Fourth Generation	$ V_{ub}' V_{cb}'  m_b' < 0.5 \text{ GeV}$
Q=-1/3 Singlet Quark	$s_2 m_s < 0.27 \text{ GeV}$
Q=+2/3 Singlet Quark	$ \lambda_{uc}  < 2.4 \cdot 10^{-4}$
Little Higgs	Tree: Same as Q=-1/3 Singlet Qk Box: Can reach observed $x_D$
Generic Z'	$M_{Z'}/C > 2.2 \cdot 10^3 \text{ TeV}$
Family Symmetries	$m_1/f > 1.2 \cdot 10^3 \text{ TeV}$
Left-Right Symmetries	No Constraint
Alternate L-R Symmetries	$M_R > 1.2 \text{ TeV}$ ( $m_{D1} = 0.5 \text{ TeV}$ ) $(\Delta m/m_{D1})/M_R > 0.4 \text{ TeV}^{-1}$
Vector Leptoquark Bosons	$M_{VLQ} > 55 (\lambda_{PP}/0.1) \text{ TeV}$
Fl-Cons Two-Higgs Doublet	No Constraint
Fl- Change Neutral Higgs I	$m_H/C > 2.4 \cdot 10^3 \text{ TeV}$
Fl-Change Neutral Higgs II	$m_H/ \Delta_{uc}  > 600 \text{ GeV}$
Scalar Leptoquark Bosons	See RPV SUSY
Higgless	$M > 100 \text{ TeV}$
Universal Extra Dimensions	No Constraint
Split Fermion	$M/ \Delta y  > 600 \text{ GeV}$
Warped Geometries	$M_1 > 3.5 \text{ TeV}$
Minimal SUSY Standard	$ (\delta_{12}^u)_{LR,LR}  < 0.035$ $ (\delta_{12}^u)_{LL,RR}  < 0.25$
SUSY Alignment	$M > 2 \text{ TeV}$
SUSY with RPV	$\lambda'_{12k} \lambda'_{11k}/m < 1.8 \cdot 10^{-4}/100 \text{ GeV}$
Split SUSY	No Constraint

# Concluding Remarks

## Experiment:

$x_D$  and  $y_D$  signals at 1% level. Great!

But more sensitivity desired.

Ultimately attain PRL criterion?

## SM Theory:

### Quarks:

To date, find  $x_D \cong y_D \cong 10^{-6}$  Tiny! But triple expansion not rapidly convergent.

### Hadrons:

Might be that  $x_D, y_D \sim 10^{-2}$  (!) but hadronic physics messy as always.

## NP Theory:

We have found which NP models **can** yield sizable  $x_D$  and  $y_D$  and which **cannot**. Charm mixing data yield useful constraints. A most welcome addition to the NP community!