

Study of the semileptonic decay $D_s^+ \rightarrow f_0(980)e^+\nu$ and implications for $B_s^0 \rightarrow J/\psi f_0$

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Abstract

Using $e^+e^- \rightarrow D_s^- D_s^{*+}$ and $D_s^{*-} D_s^+$ interactions at 4170 MeV collected with the CLEO-c detector, we investigate the semileptonic decays $D_s^+ \rightarrow f_0(980)e^+\nu$, and $D_s^+ \rightarrow \phi e^+\nu$. By examining the decay rates as functions of the four-momentum transfer squared, q^2 , we measure the ratio $\left[\frac{d\mathcal{B}}{dq^2}(D_s^+ \rightarrow f_0(980)e^+\nu)\mathcal{B}(f_0 \rightarrow \pi^+\pi^-)\right] / \left[\frac{d\mathcal{B}}{dq^2}(D_s^+ \rightarrow \phi e^+\nu)\mathcal{B}(\phi \rightarrow K^+K^-)\right]$ at q^2 of zero to be $(42 \pm 11)\%$. This ratio has been predicted to equal the rate ratio $[\mathcal{B}(B_s \rightarrow J/\psi f_0)\mathcal{B}(f_0 \rightarrow \pi^+\pi^-)] / [\mathcal{B}(B_s \rightarrow J/\psi \phi)\mathcal{B}(\phi \rightarrow K^+K^-)]$, thus indicating that the CP eigenstate $J/\psi f_0$ could be useful for measuring CP violation via B_s mixing. Assuming a simple pole model for the form factor $|f_+(q^2)|$ in the $f_0 e^+\nu$ decay, we find a pole mass of $(1.7_{-0.7}^{+4.5} \pm 0.2)$ GeV. We also determine the f_0 mass and width as $(977_{-9}^{+11} \pm 1)$ MeV, and $(91_{-22}^{+30} \pm 3)$ MeV, respectively. In addition, we present updated results for $\mathcal{B}(D_s^+ \rightarrow f_0(980)e^+\nu)\mathcal{B}(f_0 \rightarrow \pi^+\pi^-) = (0.20 \pm 0.03 \pm 0.01)\%$, and $\mathcal{B}(D_s^+ \rightarrow \phi e^+\nu) = (2.36 \pm 0.23 \pm 0.13)\%$. Assuming that the f_0 wavefunction is a combination of strange and non-strange quark-antiquark components, we use our measurement for $\mathcal{B}(D_s^+ \rightarrow f_0(980)e^+\nu)$ to extract a value of the mixing angle that we find consistent with $|\bar{s}s\rangle$ dominance, adding to the mystery as to why the f_0 decays predominantly to two pions rather than two kaons.

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I. INTRODUCTION

In this article we present a study of semileptonic decay of the D_s^+ meson into $f_0(980)e^+\nu$ and also to $\phi e^+\nu$. The $f_0(980)$ meson, a scalar, though well established experimentally, has a relatively uncertain mass and width [1]; in addition there have been claims that the quark content may be a mixture of a traditional quark-antiquark with a four-quark system [2]. Semileptonic D_s^+ decays provide a pristine environment where the f_0 is produced by an isoscalar combination of s and \bar{s} quarks. Evidence for $f_0(980)$ in semileptonic decays was seen by BaBar via the $f_0(980) \rightarrow K^+K^-$ channel, where interference was observed by an S-wave with the dominant P-wave ϕ decay [3]. The first measurement of the $D_s^+ \rightarrow f_0(980)e^+\nu$, $f_0 \rightarrow \pi^+\pi^-$ product branching fraction was recently made by CLEO [4]. The semileptonic decay diagram is shown in Fig. 1. Here we investigate this mode using a data sample of 600 pb^{-1} , approximately double the original size. This larger sample is sufficient to allow us to determine the f_0 mass and width and measure the semileptonic decay form factor as a function of the invariant four-momentum transfer squared, q^2 , between the D_s and the f_0 .

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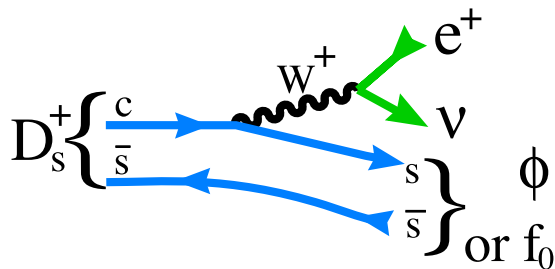


FIG. 1: The Feynman diagram for semileptonic D_s^+ decay into a ϕ or $f_0(980)$ meson.

CP violation measurements in the B_s system have concentrated on the final state $J/\psi\phi$, with $\phi \rightarrow K^+K^-$. For a review see Ref. [5]. Since this mode is not a CP eigenstate an angular analysis is required to separate the CP even and CP odd parts, and measure the CP violating phase $-2\beta_s$.

In what follows we will use a simple notation for the decay width of a particle multiplied by the branching fraction for the decay of one of its daughters. For example,

$$\Gamma(D_s^+ \rightarrow \pi^+ f_0, f_0 \rightarrow \pi^+\pi^-) \equiv \Gamma(D_s^+ \rightarrow \pi^+ f_0(980))\mathcal{B}(f_0(980) \rightarrow \pi^+\pi^-). \quad (1)$$

Stone and Zhang [6] have suggested that the $J/\psi\phi$ mode may also contain an S-wave K^+K^- system at the ϕ mass with a rate that could be as large as $\approx 5\%$ that of the ϕ . This S-wave, if it should be significant, would require additional parameters in the fit to extract $-2\beta_s$. They also suggest that the final state $J/\psi f_0$ may be a useful alternative; since it is a CP-eigenstate angular analysis is not required. Stone and Zhang estimated the branching ratio $B_s \rightarrow J/\psi f_0$ assuming equality of the ratios

$$R_{f/\phi} \equiv \frac{\Gamma(B_s^0 \rightarrow J/\psi f_0, f_0 \rightarrow \pi^+\pi^-)}{\Gamma(B_s^0 \rightarrow J/\psi\phi, \phi \rightarrow K^+K^-)} = \frac{\Gamma(D_s^+ \rightarrow f_0\pi^+, f_0 \rightarrow \pi^+\pi^-)}{\Gamma(D_s^+ \rightarrow \phi\pi^+, \phi \rightarrow K^+K^-)} \approx (20-30)\%. \quad (2)$$

The phase spaces for the B_s^0 and D_s^+ decays are quite similar, but the spin structure is not. In these D_s^+ decays both the ϕ and the f_0 are produced opposite a spin-0 pion, rather than

a spin-1 J/ψ . This consideration prompted Stone and Zhang to suggest that measuring the ratio of decay widths of semileptonic D_s^+ decays containing either an f_0 or a ϕ at four-momentum transfers, q^2 , equal to zero would give a superior prediction [6]. Specifically they propose that

$$R_{f/\phi} = \frac{\frac{d\Gamma}{dq^2}(D_s^+ \rightarrow f_0(980)e^+\nu, f_0 \rightarrow \pi^+\pi^-) |_{q^2=0}}{\frac{d\Gamma}{dq^2}(D_s^+ \rightarrow \phi e^+\nu, \phi \rightarrow K^+K^-) |_{q^2=0}}. \quad (3)$$

The point q^2 equal to zero is chosen to maximize the allowed phase space in order to make it as close as possible to that available in $B_s^0 \rightarrow J/\psi\phi$ (or f_0) decay. In this paper we will present measurements of this ratio, the form factor in the $D_s^+ \rightarrow f_0 e^+\nu$ channel, the f_0 mass and width, and update the previously published CLEO branching fractions for these two semileptonic decay modes [4].

II. EXPERIMENTAL METHOD

A. Selection of D_s Candidates

The CLEO-c detector [7] is equipped to measure the momenta and directions of charged particles, identify them using specific ionization (dE/dx) and Cherenkov light (RICH) [8], detect photons and determine their directions and energies.

In this study we use 600 pb⁻¹ of data produced in e^+e^- collisions using the Cornell Electron Storage Ring (CESR) and recorded near a center-of-mass energy (E_{CM}) of 4.170 GeV. At this energy the e^+e^- annihilation cross-section into $D_s^- D_s^{*+}$ plus $D_s^{*-} D_s^+$ is approximately 1 nb [9].

In this analysis we fully reconstruct a sample of D_s^- in several ‘‘tag’’ modes and then find candidate semileptonic decays in this sample. Mention of any specific decay implies the use of its charge-conjugate as well. The tag selection is identical to that used in our $D_s^+ \rightarrow \mu^+\nu$ paper, that can be consulted for details [10]. Briefly, we select candidates on the basis of their beam-constrained invariant mass. Then we detect an additional photon candidate from the D_s^* decay, and construct the missing mass squared, MM^{*2} recoiling against the photon and the D_s^- tag

$$\text{MM}^{*2} = (E_{\text{CM}} - E_{D_s} - E_\gamma)^2 - (\mathbf{p}_{\text{CM}} - \mathbf{p}_{D_s} - \mathbf{p}_\gamma)^2, \quad (4)$$

where E_{CM} (\mathbf{p}_{CM}) is the center-of-mass energy (momentum), E_{D_s} (\mathbf{p}_{D_s}) is the energy (momentum) of the fully reconstructed D_s^- tag, and E_γ (\mathbf{p}_γ) is the energy (momentum) of the additional photon. In performing this calculation we use a kinematic fit that constrains the decay products of the D_s^- to the known D_s mass and conserves overall momentum and energy. All photon candidates in the event are used, except for those that are decay products of the D_s^- tag candidate. Regardless of whether or not the photon forms a D_s^* with the tag, for real $D_s^* D_s$ events MM^{*2} should peak at the D_s^* mass-squared.

We list the number of signal events in each mode in Table I by finding the number of events within ± 17.5 MeV of the D_s mass. For ease of further analysis we sum all tag modes together, as shown in Fig. 2(a).

The MM^{*2} distributions for events in the D_s^- invariant mass signal region (± 17.5 MeV from the D_s mass) are shown in Fig. 2(b). In order to find the number of tags used for

TABLE I: Tagging modes and numbers of signal and background events, within ± 17.5 MeV of the D_s^- mass for each mode, determined from two-Gaussian function fits to the invariant mass plots, and the number of tags in each mode including the γ from the $D_s^* \rightarrow \gamma D_s$ transition, within an interval $3.782 < \text{MM}^{*2} < 4.0$ GeV², as determined from fits of the MM^{*2} distributions (see text) to a signal Crystal Ball function (see text) and two 5th order Chebychev background polynomial functions.

Mode	Invariant Mass		MM ^{*2}	
	Signal	Background	Signal	Background
$K^+ K^- \pi^-$	26534±274	25122	16087±373	39563
$K_S K^-$	6383±121	3501	4215±228	6297
$\eta \pi^-; \eta \rightarrow \gamma \gamma$	2993 ± 156	5050	2005±145	5016
$\eta' \pi^-; \eta' \rightarrow \pi^+ \pi^- \eta, \eta \rightarrow \gamma \gamma$	2293±82	531	1647±131	1565
$K^+ K^- \pi^- \pi^0$	11649±754	78588	6441±471	89284
$\pi^+ \pi^- \pi^-$	7374±303	60321	5014±402	43286
$K^{*-} K^{*0}; K^{*-} \rightarrow K_S^0 \pi^-, K^{*0} \rightarrow K^+ \pi^-$	4037±160	10568	2352±176	12088
$\eta \rho^-; \eta \rightarrow \gamma \gamma, \rho^- \rightarrow \pi^- \pi^0$	5700±281	24444	3295±425	24114
$\eta' \pi^-; \eta' \rightarrow \rho^0 \gamma,$	3551±202	19841	2802±227	17006
Sum	70514 ± 963	227966	43859±936	238218

further analysis we perform a two-dimensional binned maximum likelihood fit of the MM^{*2} distribution and the invariant mass distribution in the interval ± 60 MeV from the D_s mass and $3.5 < \text{MM}^{*2} < 4.25$ GeV². The background has two components, both described by 5th order Chebyshev polynomials in MM^{*2} ; the first comes from the background under the invariant mass peak, defined by the sidebands, and the second is due to multiple photon combinations. In both cases we allow the parameters to float.

We find a total of $43859 \pm 936 \pm 877$ events within the interval $3.782 < \text{MM}^{*2} < 4.0$ GeV² and having an invariant mass within ± 17.5 MeV of the D_s mass, where the total number of events is the sum of the yields from the fits to each mode as shown in Table I. The first uncertainty in the total is statistical and the second is systematic.

B. Signal Reconstruction

We next describe the reconstruction of $D_s^+ \rightarrow f_0 e^+ \nu$, $f_0 \rightarrow \pi^+ \pi^-$, and also $D_s^+ \rightarrow \phi e^+ \nu$, $\phi \rightarrow K^+ K^-$. We select events within the MM^{*2} region shown in Fig. 2(b) for further analysis. We note that the limits are rather wide. We use this selection because the background in the signal side is rather small and the errors are minimized by taking as many tags as possible.

Candidate events are selected that contain a charged track of opposite sign to the D_s^- tag that is positively identified as a positron. Electrons and positrons are identified on the basis of a likelihood ratio constructed from three inputs: the ratio between the energy deposited in the calorimeter and the momentum measured in the tracking system, the specific ionization dE/dx measured in the drift chamber, and RICH information [11]. Our selection efficiency averages 0.95 in the momentum region 0.3-1.0 GeV, and 0.71 in the region 0.2-0.3 GeV. The

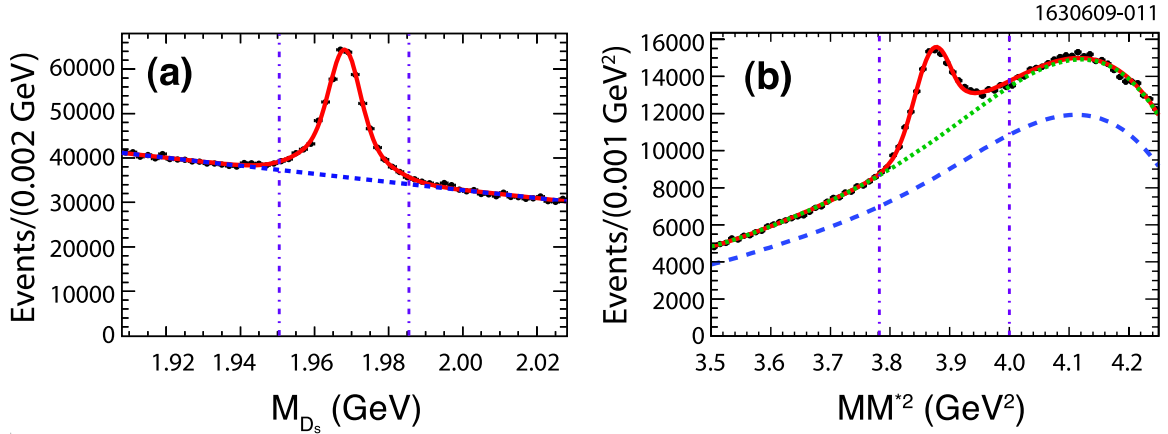


FIG. 2: (a) Invariant mass of D_s^- candidates summed over all decay modes and fit to a two-Gaussian signal shape plus a straight line for the background. The vertical dot-dashed lines indicate the ± 17.5 MeV definition of the signal region. (b) The MM^{*2} distribution summed over all modes. The curves are fits to the number of signal events using the Crystal Ball function and two 5th order Chebychev background functions; the dashed curve shows the background from fake D_s^- tags, while the dotted curve in (b) shows the sum of the backgrounds from multiple photon combinations and fake D_s^- tags. The vertical dashed lines show the region of events selected for further analysis.

average fractions of charged pions and kaons incorrectly identified as positrons averaged over the relevant momentum range are approximately 0.1%. We also require an additional pair of tracks with opposite charge that are both identified as pions or kaons using the dE/dx and RICH systems.

Since we are searching for events containing a single missing neutrino the missing mass squared for the f_0 mode, MM^2 , evaluated by taking into account the observed e^+ , π^+ , π^- , D_s^- , and γ should peak at zero; the MM^2 is computed as

$$MM^2 = (E_{\text{CM}} - E_{D_s} - E_\gamma - E_e - E_{\pi^+} - E_{\pi^-})^2 - (\mathbf{p}_{\text{CM}} - \mathbf{p}_{D_s} - \mathbf{p}_\gamma - \mathbf{p}_e - \mathbf{p}_{\pi^+} - \mathbf{p}_{\pi^-})^2, \quad (5)$$

where E_e (\mathbf{p}_e) are the energy (momentum) of the candidate positron, E_π (\mathbf{p}_π) are the energy (momenta) of the candidate pions, and all other variables are the same as defined in Eq. (4). A similar equation applies for the ϕ mode with the pions replaced by kaons.

We also make use of a set of kinematical constraints and fit each event to two hypotheses, one of which is that the D_s^- tag is the daughter of a D_s^{*-} , and the other that the D_s^{*+} decays into γD_s^+ with the D_s^+ subsequently decaying into either $\pi^+\pi^-e^+\nu$, or $K^+K^-e^+\nu$. The kinematical constraints, in the e^+e^- center-of-mass frame, are

$$\begin{aligned} \mathbf{p}_{D_s} + \mathbf{p}_{D_s^*} &= 0, \\ E_{\text{CM}} &= E_{D_s} + E_{D_s^*}, \\ E_{D_s^*} &= \frac{E_{\text{CM}}}{2} + \frac{M_{D_s^*}^2 - M_{D_s}^2}{2E_{\text{CM}}} \text{ or } E_{D_s} = \frac{E_{\text{CM}}}{2} - \frac{M_{D_s^*}^2 - M_{D_s}^2}{2E_{\text{CM}}}, \text{ and} \\ M_{D_s^*} - M_{D_s} &= 143.8 \text{ MeV}. \end{aligned} \quad (6)$$

In addition, we constrain the invariant mass of the D_s^- tag to the known D_s mass. This gives us a total of seven constraints. The missing neutrino four-vector needs to be determined, so we are left with a three-constraint fit. We perform an iterative fit that minimizes χ^2 . As we do not want to be subject to systematic uncertainties that depend on understanding the absolute scale of the errors, we do not make a χ^2 cut but simply choose the photon and the decay sequence in each event with the minimum χ^2 .

We model the MM^2 signal distributions for both $f_0 e^+ \nu$ and $\phi e^+ \nu$ final states as the sum of two Crystal Ball (CB) functions centered at zero [12] plus a wide Gaussian shape that serves to model the tails. The Monte Carlo simulations are shown in Fig. 3. The results are summarized in Table II.

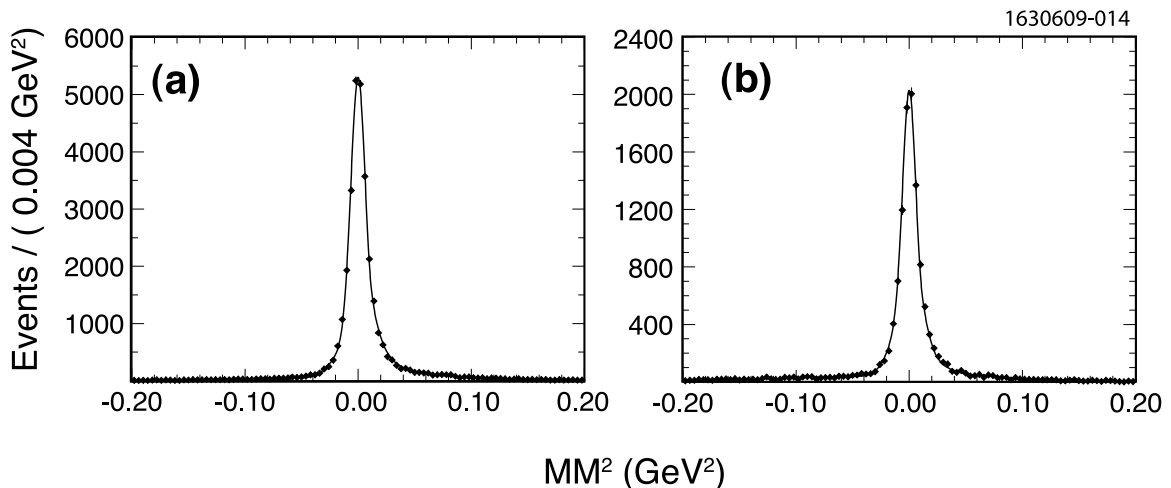


FIG. 3: MM^2 distributions from Monte Carlo simulation for (a) $D_s^+ \rightarrow f_0 e^+ \nu$ and (b) $D_s^+ \rightarrow \phi e^+ \nu$. The curves are the sum of two CB functions with means fixed at zero, n fixed at 1.8 and α values constrained to be equal, and a single Gaussian shape. (The parameters are listed in Table II.)

TABLE II: Parameters of signal fits to two CB functions plus a Gaussian shape; “ \mathcal{F} ” indicates the fraction of the yield in each fit component. The parameter n used in both CB functions is fixed at 1.8. The parameter α is kept the same for both CB functions. The mean and width of the Gaussian is kept the same for both decay modes. All means and σ ’s are in units of GeV^2 .

	1 st CB function			2 nd CB function		Gaussian		
	$\mathcal{F}(\%)$	σ_1	α	$\mathcal{F}(\%)$	σ_2	$\mathcal{F}(\%)$	mean	σ
$\pi^+ \pi^-$	63	$(6.30 \pm 0.15) \times 10^{-3}$	1.417 ± 0.014	30	$(1.67 \pm 0.07) \times 10^{-2}$	7	-0.0253	0.0992
$K^+ K^-$	45	$(5.61 \pm 0.35) \times 10^{-3}$	1.610 ± 0.003	37	$(1.32 \pm 0.09) \times 10^{-2}$	18	-0.0253	0.0992

We proceed by performing a simultaneous fit to the D_s^- invariant mass, using a mass range ± 70 MeV from the nominal mass, and the MM^2 for $\pi^+ \pi^- e^+ \nu$ and $K^+ K^- e^+ \nu$ to the Monte Carlo generated fitting functions letting only the normalizations float.¹ The resulting

¹ This is the same procedure as used in our $D_s^+ \rightarrow \mu^+ \nu$ analysis.

MM² distributions are shown in Fig. 4. Here and in subsequent analyses we require that the $\pi^+\pi^-$ invariant mass is above 0.6 GeV, in order to eliminate the $K_s^0 e^+\nu$ channel, and limit the K^+K^- invariant mass to be below 1.08 GeV.

Our next step is to determine the range of MM² to select for further analysis. Thus, we consider the backgrounds in our MM² sample. In general the background arises from two sources: one from real D_s^+ decays and the other from the background under the single-tag signal peaks (fake D_s^-). For $\phi e^+\nu$ only the background from fake D_s^- is significant. For $f_0 e^+\nu$ we show in Fig. 4(a) the background from the invariant mass sidebands, suitably scaled. We also find that there are two sources of real D_s^+ background, namely $\eta' e^+\nu$ decays where the $\eta' \rightarrow \rho\gamma$, and a very small component of $\phi e^+\nu$, $\phi \rightarrow \pi^+\pi^-\pi^0$ when the π^0 is ignored. We also show these backgrounds as a function of MM². We use a second order polynomial to parameterize the fake D_s^- background, and the other two backgrounds are parameterized by the sum of two bifurcated Gaussian shapes determined by simulating the background processes. (A bifurcated Gaussian shape has a different widths below and above the mean.) We are able to fix the $\pi^+\pi^-\pi^0 e^+\nu$ rate as 0.31%, using our previous measurement of $D_s^+ \rightarrow \phi e^+\nu$, but allow the $\pi^+\pi^-\gamma e^+\nu$ yield to float in the fit as our measured branching fraction is only accurate to $\pm 35\%$ [4].

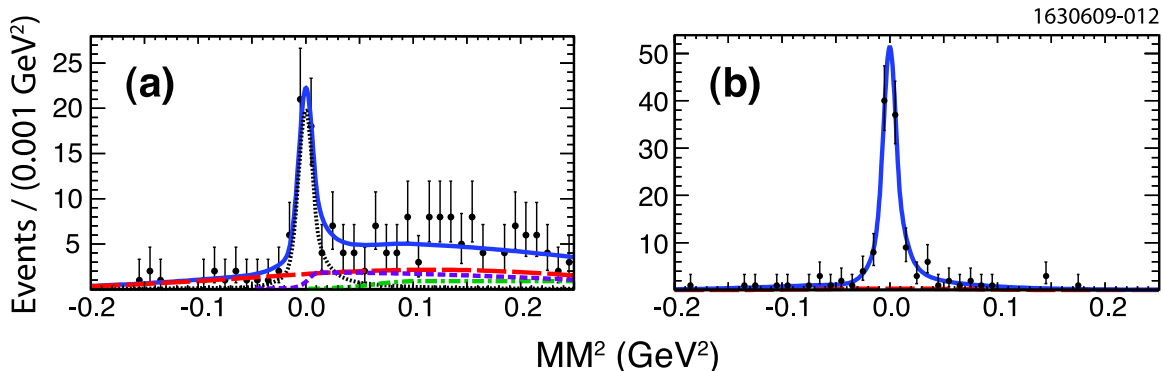


FIG. 4: (a) The MM² distribution for $\pi^+\pi^-e^+\nu$. The dotted curve shows the signal, the long-dashed distributions are from sidebands of the D_s^- candidate invariant mass distributions while the shorter dashed curve in (a) shows the $\eta' \rightarrow \pi^+\pi^-\gamma$ background level, while the dashed-dot curve shows the (small) $\pi^+\pi^-\pi^0$ background from $\phi e^+\nu$. The solid curve shows the total. (b) The MM² distribution for $K^+K^-e^+\nu$, where the total and the small side-band background is shown.

For further analysis we restrict ourselves to the interval $-0.04 < \text{MM}^2 < 0.04 \text{ GeV}^2$, which is 88.4% efficient on $f_0 e^+\nu$ and is 84.4% efficient on $\phi e^+\nu$. In these regions we find 42.9 ± 6.7 $\pi^+\pi^-e^+\nu$ signal events, 7.7 ± 2.0 $\eta' e^+\nu$ events, 0.5 $\phi e^+\nu$ events (fixed), and 13.4 ± 0.9 fake D_s^- background. There are 107.0 ± 9.8 $K^+K^-e^+\nu$ events of which 2.5 ± 0.7 are fake D_s^- background.

C. K^+K^- and $\pi^+\pi^-$ Invariant Mass Distributions

We next include the di-hadron mass as part of the joint fit to the D_s^- mass, and MM². Figure 5 shows the $\pi^+\pi^-$ invariant mass distribution for $\pi^+\pi^-e^+\nu$ events, and the K^+K^- invariant mass distribution in our selected sample of $K^+K^-e^+\nu$ events, within the interval

$-0.04 < MM^2 < 0.04 \text{ GeV}^2$. Both distributions are dominated by a single di-hadron resonance, either the $f_0(980)$ in the $\pi^+\pi^-$ mode or the ϕ in the K^+K^- mode. The backgrounds are included in the fit in the same manner as for the MM^2 distributions. We next find the size of the f_0 and ϕ signals.

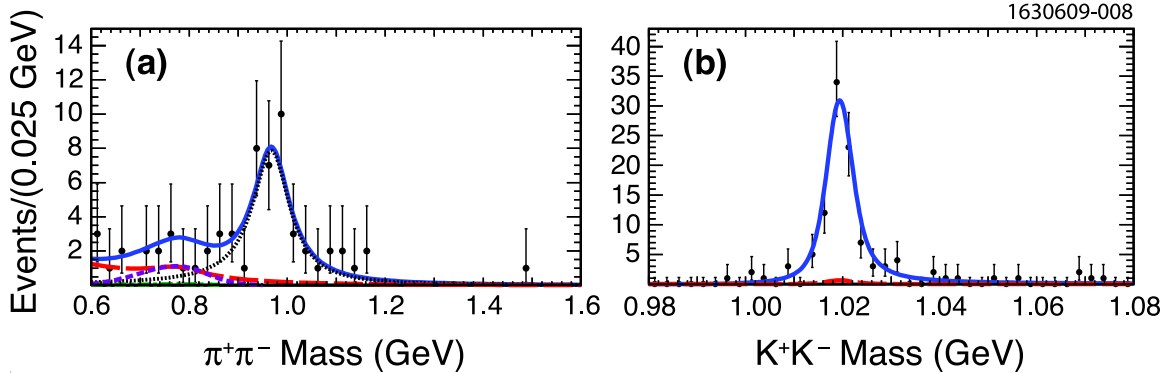


FIG. 5: Invariant mass distributions for (a) $\pi^+\pi^-$ and (b) K^+K^- for semileptonic decays.. The dotted curve shows the signal, the long-dashed distributions are from sidebands of the D_s^- candidate invariant mass distributions while the shorter dashed curve in (a) shows the $\eta' \rightarrow \pi^+\pi^-\gamma$ background level, and the dashed-dot curve in (a) shows the (small) $\pi^+\pi^-\pi^0$ background from $\phi e^+\nu$. The solid curves show the totals.

For the $K^+K^-e^+\nu$, the signal fitting function consists of the sum of a P-wave Breit-Wigner shape whose mass and width are fixed to the values reported by the Particle Data Group (PDG) [1] convoluted with our Gaussian experimental resolution, and a fixed background function determined from the D_s^- sidebands described by the same ϕ signal shape added to a wide component that is exponentially decreasing with mass.

For the $\pi^+\pi^-$ mass distribution we include three separate shapes to fit the backgrounds. (1) The $\pi^+\pi^-\gamma$ distribution is described by a bifurcated-Gaussian function; (2) the $\pi^+\pi^-\pi^0$ distribution is described by a bifurcated-Gaussian shape plus a Breit-Wigner shape describing the ρ , here our knowledge of the $\phi e^+\nu$ decay rate allows us to fix the magnitude of this component; (3) the shape describing non- D_s^- background is an exponential plus a Breit-Wigner function describing the ρ , the size of this component is determined from the sidebands of the D_s^- signal.

The mass and width of the $f_0(980)$ are not well determined. The PDG makes an estimate of the mass of $(980 \pm 10) \text{ MeV}$, but does not explicitly average any measurements. They also state that the width is very model dependent and give a range of 40-100 MeV, again not forming an average [1]. Here we provide measurements in an extremely clean environment. We fit the $\pi^+\pi^-$ mass distribution, $\sqrt{s} \equiv M_{\pi\pi}$ with a Breit-Wigner type resonance form with the threshold behavior suggested in [13]

$$P(M_{\pi\pi}^2) = \frac{M_{f_0}\Gamma(s)}{(s - M_{f_0}^2)^2 + M_{f_0}^2\Gamma^2(s)}, \quad (7)$$

where M_{f_0} denotes the f_0 mass and $\Gamma(s)$ the width which is a function of the mass having

the explicit dependence

$$\Gamma(s) = \Gamma_0 \frac{\sqrt{(s - 4M_\pi^2)} M_{f_0}}{\sqrt{(M_{f_0}^2 - 4M_\pi^2)} M_{\pi\pi}}, \quad (8)$$

where M_π is the charged pion mass, and Γ_0 the width in the limit where the pion becomes massless.

We find that $M_{f_0} = (968 \pm 9)$ MeV, and $\Gamma_0 = (92_{-21}^{+28})$ MeV. Later in this paper we will slightly refine these values of the mass and width, which we use for subsequent analysis. We note that performing this fit using a normal Breit-Wigner function gives $M_{f_0} = (966 \pm 9)$ MeV, $\Gamma = (89_{-20}^{+26})$ MeV.

III. BRANCHING FRACTIONS AND q^2 DISTRIBUTIONS

A. Main Experimental Results

The main aims of this paper are to provide the branching fraction and q^2 dependence for the $f_0 e^+ \nu$ mode and to measure the relative rates at $q^2 = 0$ of the f_0 and ϕ modes. The BaBar collaboration previously measured $\mathcal{B}(D_s^+ \rightarrow \phi e^+ \nu) = (2.61 \pm 0.03 \pm 0.08 \pm 0.15)\%$,² and the form factors for this decay given in the Appendix. They also found that the S-wave contribution to the $K^+ K^- e^+ \nu$ rate is a small fraction $(0.22_{-0.08}^{+0.12})\%$ of the total rate in the mass interval between 1.01 and 1.03 GeV [3].

In general, for a decay of a pseudoscalar D_s^+ to a scalar meson such as the $f_0(980)$, the only Lorentz invariants in the problem are the invariant square P^2 of the summed four-momentum of the D_s and f_0 , and the invariant square q^2 of the four-momentum transfer between the D_s and the f_0 . The hadronic current describing the decay can be expressed as

$$\langle f_0 | \bar{s} \gamma^\mu (1 - \gamma^5) c | D_s \rangle = f_+(q^2) P^\mu + f_-(q^2) q^\mu, \quad (9)$$

where $f_\pm(q^2)$ are arbitrary form factors. Since this transition is between a pseudoscalar D_s^+ and a scalar f_0 , only the axial-vector part of the current contributes. The term containing $f_-(q^2)$ gets multiplied by the positron mass squared, and becomes vanishingly small [14].

The decay rate can be written as

$$\frac{d\Gamma(D_s^+ \rightarrow f_0 e^+ \nu)}{dq^2} = \frac{G_F^2 |V_{cs}|^2}{24\pi^3} p_{f_0}^3 |f_+(q^2)|^2, \quad (10)$$

where G_F is the Fermi constant, p_{f_0} is the magnitude of the three-momentum of the f_0 in the D_s^+ rest frame, and $|V_{cs}|$ is a Cabibbo-Kobayashi-Maskawa matrix element equal to ≈ 0.97 [1]. Therefore, we can use our data to determine $|f_+(q^2)|$.

Equation 10 is strictly correct only for zero width. It can be modified, however, to correct for the finite width mesons as described by Isgur *et al.* [15]. Here we integrate the product of $P(s)$ as defined in Eq. 7 with the right-hand side of Eq. 10 over $dM_{\pi\pi}$.

² The first error is statistical, the second experimental systematic, and the third due to the absolute D_s^+ branching fraction scale.

The form factors for $\phi e^+\nu$ are somewhat more complicated as the hadronic transition involves a transformation from a pseudoscalar D_s to a vector ϕ . The formulas for this case are given in the Appendix. The form factors for this decay have been measured rather precisely by BaBar [3], and we use these for further analysis of this mode. Our aim is to determine our own rate at q^2 equals zero, so as to cancel systematic errors in the ratio of rates for $f_0 e^+\nu/\phi e^+\nu$.

We separate the data into five q^2 intervals and perform fits to the mass distributions. These are shown in Fig. 6 for $\pi^+\pi^-e^+\nu$, and in Fig. 7 for $K^+K^-e^+\nu$. The resulting signal yields are given in Table III. For the $\pi^+\pi^-e^+\nu$ mode, the fitting function consists of a sum of a signal Breit-Wigner with mass and width fixed to the values we found by summing all q^2 intervals for the f_0 , a fixed background function from the D_s^- sidebands, a floating background shape to account for $D_s^+ \rightarrow \eta' e^+\nu$, $\eta' \rightarrow \pi^+\pi^-\gamma$, and a fixed background function to account for $D_s^+ \rightarrow \phi e^+\nu$, $\phi \rightarrow \pi^+\pi^-\pi^0$. The $K^+K^-e^+\nu$ shape is defined in section II C and is kept the same in each q^2 interval.

TABLE III: Number of events ($\#$) and efficiencies ϵ in q^2 intervals.

q^2 interval (GeV ²)	$\#f_0 e^+\nu$	$\epsilon(f_0)$ (%)	$\#f_0 e^+\nu$ corrected	$\#\phi e^+\nu$	$\epsilon(\phi)$ (%)	$\#\phi e^+\nu$ corrected
0-0.2	14.6±3.9	45.5	32.1±8.6	30.5±5.1	25.9	118±20
0.2-0.4	12.3±3.5	50.2	24.5±7.0	30.5±5.3	22.9	133±23
0.4-0.6	12.4±3.5	53.6	23.1±6.5	22.3±4.4	21.5	104±20
0.6-0.8	4.2±2.2	56.9	7.4±3.9	21.5±4.7	15.6	138±30
0.8-2.0	0±1.0	55.9	0±1.8	1.6±1.1	9.8	16±11
Sum	43.5±6.7		87.1±13.5	106.4±9.8		508±49

The product branching fractions can be extracted by dividing the efficiency-corrected event sums in each mode by the number of D_s^- tags. We find

$$\begin{aligned} \mathcal{B}(D_s^+ \rightarrow f_0(980)e^+\nu, f_0 \rightarrow \pi^+\pi^-) &= (0.20 \pm 0.03 \pm 0.01)\%, \\ \mathcal{B}(D_s^+ \rightarrow \phi e^+\nu, \phi \rightarrow K^+K^-) &= (1.16 \pm 0.11 \pm 0.06)\%, \end{aligned} \quad (11)$$

where the first error is statistical and the second systematic (see section III B). These values are consistent with, though the f_0 mode is somewhat larger than, our previously published values of $(0.13 \pm 0.04 \pm 0.01)\%$ ³ and $(1.13 \pm 0.19 \pm 0.06)\%$ for the f_0 and ϕ modes, respectively. We can compare directly with the branching ratio measurement of the ϕ mode of BaBar [3] by dividing by the $\phi \rightarrow K^+K^-$ branching fraction of $(49.2 \pm 0.6)\%$ [1]. Our measurement then becomes

$$\mathcal{B}(D_s^+ \rightarrow \phi e^+\nu) = (2.36 \pm 0.23 \pm 0.13)\%, \quad (12)$$

while the BaBar measurement is $(2.61 \pm 0.03 \pm 0.08 \pm 0.15)\%$.

Figure 8 shows the q^2 distributions for the $f_0 e^+\nu$ and $\phi e^+\nu$ channels. We fit the $\phi e^+\nu$ channel using the form factors determined by the BaBar collaboration letting only the normalization float [3]. The fit yields 478 ± 45 events compared with 508 ± 49 from summing the data.

³ Our result is larger here because the previous analysis assumed the f_0 natural width was 50 MeV.

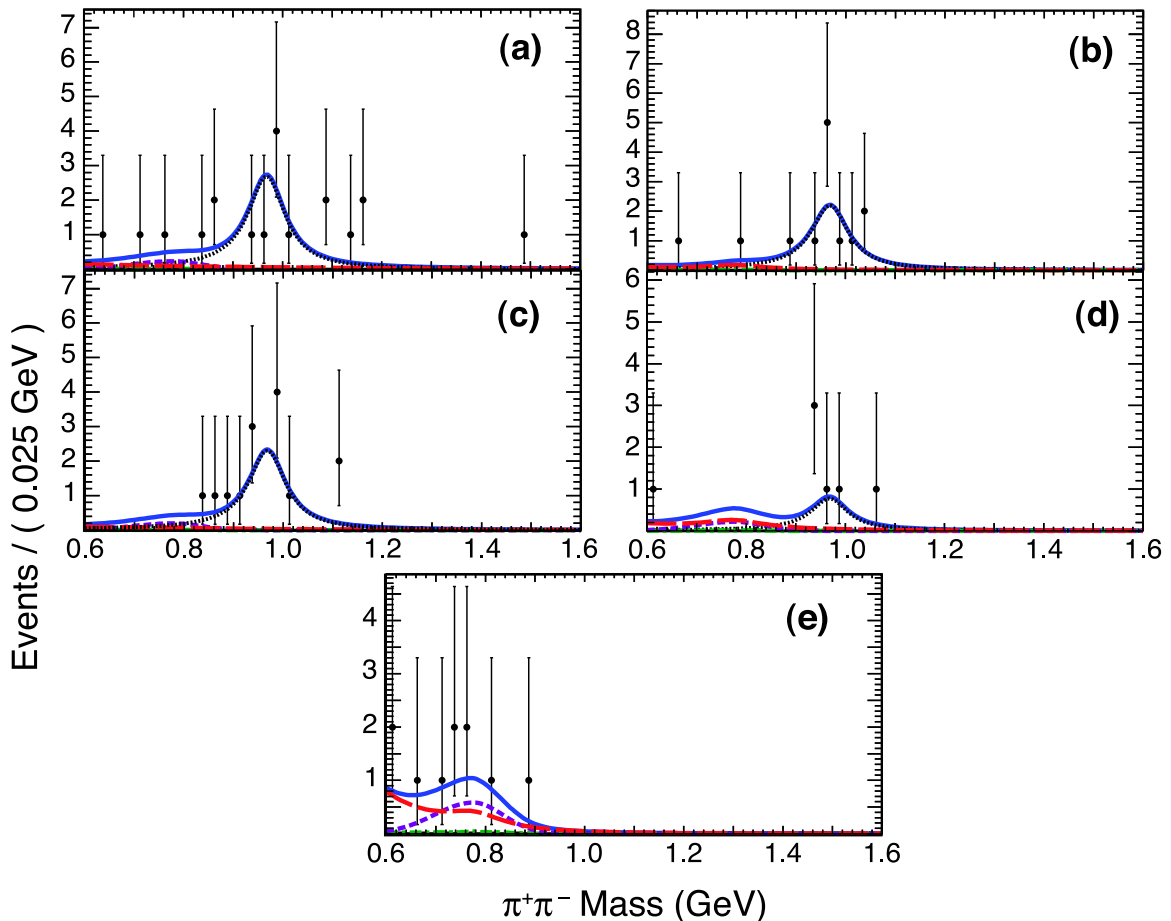


FIG. 6: The invariant $\pi^+\pi^-$ mass in the semileptonic mode for the different q^2 intervals in units of GeV^2 : (a) 0-0.2, (b) 0.2-0.4, (c) 0.4-0.6, (d) 0.6-0.8 and (e) 0.8-2.0. The fits to the data are described in the text. The dotted curves show the signal, the long-dashed distributions are from sidebands of the D_s^- candidate invariant mass distributions, while the shorter dashed curves show the $\eta' \rightarrow \pi^+\pi^-\gamma$ background level, and the dashed-dot curves show the (small) $\pi^+\pi^-\pi^0$ background from $\phi e^+\nu$. The solid curves show the total.

For $f_0 e^+\nu$ we fit the q^2 distribution to Eq. 10 using a simple pole model for the shape of the form factor, $|f_+(q^2)| = 1/(1 - q^2/M_{\text{pole}}^2)$. (We exclude the highest q^2 point from the fit.) This procedure gives a value of $(1.7_{-0.7}^{+4.5})$ GeV for M_{pole} . We then vary our measured f_0 width by one standard deviation and repeat the procedure to find the systematic uncertainty, resulting in a ± 0.2 GeV systematic error. Using the central value of the pole mass we find that the fitted yield for $f_0 e^+\nu$ is 86.6 ± 13.0 corrected events compared with 87.1 ± 13.5 corrected events found by summing over the q^2 intervals.

We now use these fits to determine $R_{f/\phi}$. At zero q^2 we find 38.8 ± 9.3 corrected $f_0 e^+\nu$ events, and 92.6 ± 8.7 $\phi e^+\nu$ events. Thus

$$R_{f/\phi} = (42 \pm 11)\%. \quad (13)$$

Here any errors in the event tagging or reconstruction efficiencies cancel. The systematic

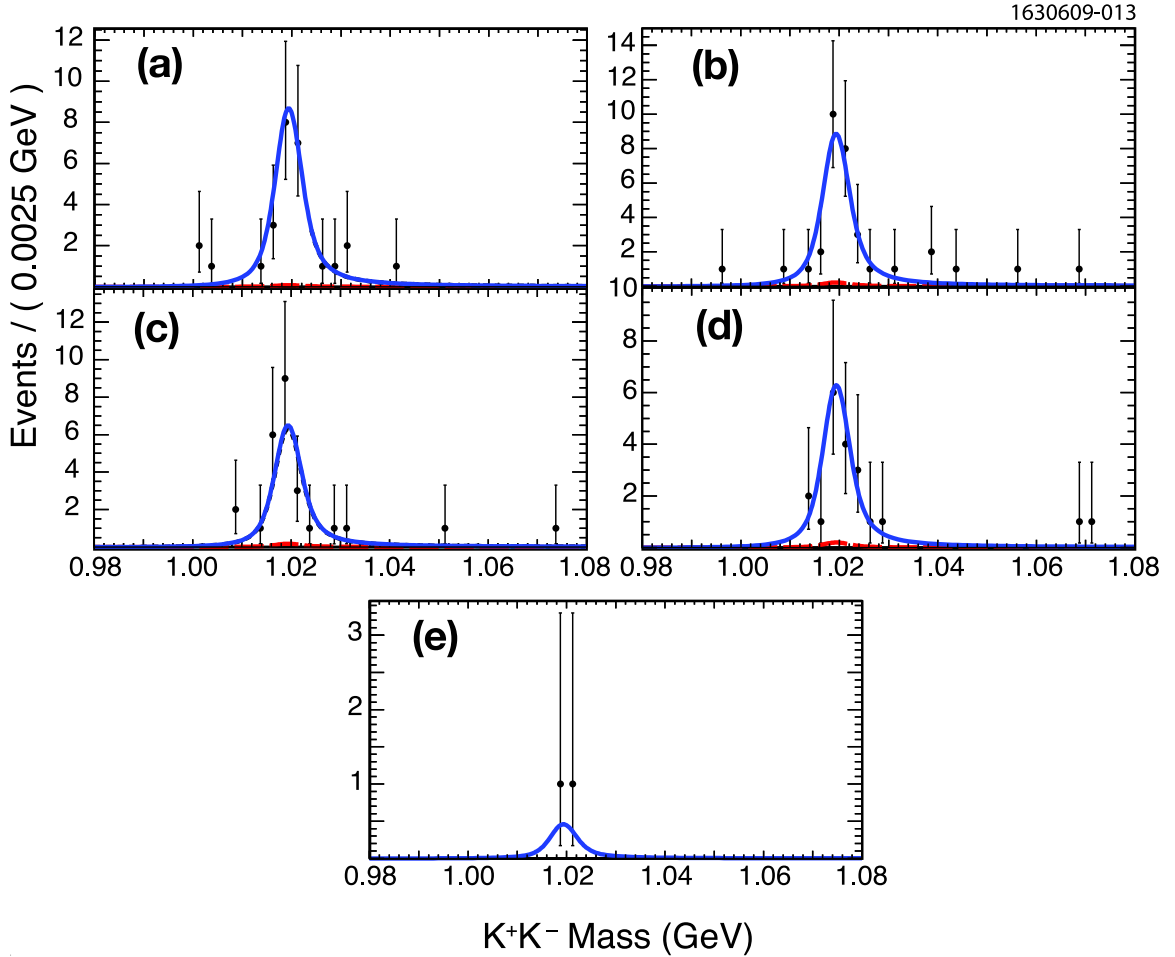


FIG. 7: The invariant K^+K^- mass in the semileptonic mode for the different q^2 intervals in units of GeV^2 : (a) 0-0.2, (b) 0.2-0.4, (c) 0.4-0.6, (d) 0.6-0.8 and (e) 0.8-1.0. The solid curve show the total and the long-dashed distributions are from sidebands of the D_s^- candidate invariant mass distributions.

error due to the uncertainties on the BaBar form factors are much smaller than the statistical errors.

B. Systematic Errors

The systematic errors on these branching fractions are given in Table IV. The error on track finding is determined from a detailed comparison of the simulation with double tag events where the entire event can be reconstructed even if one track is missed [16]. The particle identification uncertainty on the two hadronic tracks is listed as twice the uncertainty on the identification of one hadron. The error on the number of tags of $\pm 2\%$ is assigned by varying the fitting functions and ranges. Additional systematic errors arising from the background estimates are at the 1% level. Final state radiation effects are included in our simulations and the resulting accumulated uncertainty from the two-hadrons and the lepton is on the order of 1%.

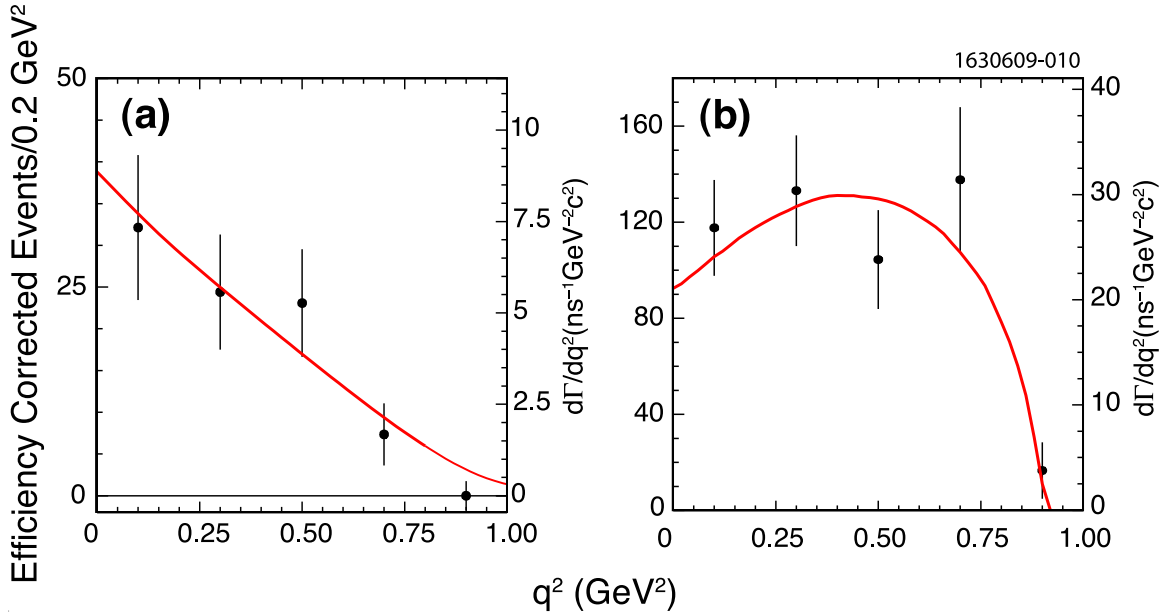


FIG. 8: The q^2 distributions for (a) $D_s^+ \rightarrow f_0 e^+ \nu$ fit to a function allowing a varying pole mass and (b) $D_s^+ \rightarrow \phi e^+ \nu$ fit to a function with form factors fixed to those measured by BaBar [3].

We have previously checked the resolution in MM^2 by fitting to a sample of $D_s^+ \rightarrow K^+ K^0$ candidates where we first found the K^0 and then ignored it [10]. The resolution agreed with the simulation to 1.7%. Here, because we have a positron in the final state instead of a K^+ we use an uncertainty of $\pm 3.4\%$ to allow for uncertainties associated with radiation in detector material. We have made wide selection cuts to minimize the effect of any resolution errors on our extracted rates.

We note that there is no enhancement in our ability to find tags in $f_0 e^+ \nu$ or $\phi e^+ \nu$ events (tag bias) as compared with events where the D_s^+ decays generically. Using a Monte Carlo simulation for each tag mode independently and then average the results based on the known tag fractions, we find a correction factor of $(0 \pm 1\%)$, which we assign as a systematic error.

TABLE IV: Systematic errors on determination of the branching fractions.

Error Source	Size (%)
Track finding	2.1
Hadron identification	2.0
Electron identification	1.0
MM^2 width	3.4
Background	1.0
Number of tags	2.0
Tag bias	1.0
Total	5.2

IV. FURTHER RESULTS AND IMPLICATIONS

A. Strange Mixing Angle of the $f_0(980)$

In one class of models, the $f_0(980)$ is thought to be a $J^P = 0^+$ state described in the quark model as composed of quark-antiquark states which form a mixture of strange and non-strange components characterized by a mixing angle θ defined in terms of the wave-function as described by Aliev and Savci [17]

$$|f_0\rangle = \cos\theta |\bar{s}s\rangle + \sin\theta \frac{1}{\sqrt{2}} |(\bar{u}u + \bar{d}d)\rangle. \quad (14)$$

Since the central value of the mass is below threshold for decay into a kaon pair, the branching ratio to physical $\bar{s}s$ states is suppressed. BES has extracted relative branching ratios using $\psi(2S) \rightarrow \gamma\chi_{c0}$ decays where the $\chi_{c0} \rightarrow f_0(980)f_0(980)$, and either both $f_0(980)$'s decay into $\pi^+\pi^-$ or one into $\pi^+\pi^-$ and the other into K^+K^- [18]. From their results we obtain

$$\frac{\mathcal{B}(f_0 \rightarrow K^+K^-)}{\mathcal{B}(f_0 \rightarrow \pi^+\pi^-)} = (25_{-11}^{+17})\%. \quad (15)$$

Assuming that the $\pi\pi$ and KK decays are dominant we can also extract

$$\mathcal{B}(f_0 \rightarrow \pi^+\pi^-) = (50_{-9}^{+7})\%, \quad (16)$$

where we have assumed that the only other decays are to $\pi^0\pi^0$, 1/2 of the $\pi^+\pi^-$ rate, and to neutral kaons, equal to charged kaons. This approach has been used by B. El-Bennich *et al.* [19]. It is possible, however, that the decays $f_0 \rightarrow \pi^+\pi^-\pi^+\pi^-$ or $\pi^+\pi^-\pi^0\pi^0$ are significant. If so, then the branching fraction of f_0 into two pions will be smaller,

Using the above branching fraction we extract $\mathcal{B}(D_s^+ \rightarrow f_0 e^+\nu) = (0.40 \pm 0.06 \pm 0.06)\%$, where the error on the $f_0 \rightarrow \pi^+\pi^-$ rate has been included in the systematic error, and, in fact, is the dominant contribution.

We can use this rate to estimate the mixing angle $\cos\theta$. Aliev and Savci using QCD sum rules predict that $\mathcal{B}(D_s^+ \rightarrow f_0 e^+\nu) = \cos^2\theta \times (0.41)\%$. This provides us with a value of

$$\cos^2\theta = 0.98_{-0.21}^{+0.02}, \quad (17)$$

where we have truncated the positive uncertainty so that the value does not exceed unity. This determination gives a somewhat larger but consistent evaluation with other methods [19]. If the f_0 decay rate into four pions is significant, our value of the mixing angle will be somewhat larger.

The near unity value for $\cos^2\theta$ poses a conundrum: If the f_0 is dominantly an $|\bar{s}s\rangle$ state, why does it decay predominantly into $\pi\pi$? A solution may be that the quark content of the f_0 is not composed of just a quark-antiquark state, but may contain a some four-quark content [2].

B. The $f_0(980)$ Mass and Width

We now return to the determination of the f_0 mass and width. In the fit described above we used a relativistic Breit-Wigner function. We did not, however, take into account that

the phase space for the larger f_0 masses is somewhat smaller, due to the finite D_s mass, than for smaller f_0 masses. The changes to the invariant mass spectrum of the f_0 have been parameterized by Dosch *et al.* [13]. The double differential decay rate in terms of $s \equiv M_{\pi\pi}^2$ is

$$\frac{d^2\Gamma(s, q^2)}{dsdq^2} = \frac{G_F^2 |V_{cs}|^2}{192\pi^4 M_{D_s}^3} \lambda^{3/2}(M_{D_s}^2, s, q^2) |f_+(q^2)|^2 P(s), \quad (18)$$

where $\lambda(x, y, z) = x^2 + y^2 + z^2 - 2xy - 2xz - 2yz$ and is equal to $4M_{D_s}^2$ times the square of the momentum of the f_0 in the D_s rest frame, and $P(s)$ is the relativistic Breit-Wigner shape given in equations 7 and 8. The resulting, slightly shifted mass spectrum, then is given by integrating the above expression over q^2

$$\frac{d\Gamma}{d\sqrt{s}} = 2\sqrt{s} \int_0^{(M_{D_s} - \sqrt{s})^2} \frac{d^2\Gamma(s, q^2)}{dsdq^2} dq^2. \quad (19)$$

We perform this integral using our fitted form factor shape for $|f_+(q^2)|$ and our original fitted parameters for the relativistic Breit-Wigner function. This generates a new shape for the mass spectrum, that we fit to the data. In principle this procedure can be iterated many times; here we find that only one iteration is sufficient. The resulting curve fitted to the data is shown in Fig. 9. The systematic errors on the mass are small compared to the statistical error. The masses of known particles, K_S , Λ , ϕ , D^0 are reproduced in other CLEO analyses to better than 1 MeV. Changing the background level by ± 1 standard deviation changes the mass by a negligible amount. For the width we have identified two sources of systematic error, one the resonance shape and the other the form factor shape. Changing the Breit-Wigner shape from relativistic to non-relativistic causes a change of ± 2 MeV. We also vary the form factor by \pm one standard deviation resulting in a change of about ± 2 MeV. Our final values for the mass and width are

$$M_{f_0} = (977_{-9}^{+11} \pm 1) \text{ MeV}, \quad \text{and} \quad \Gamma_0 = (91_{-22}^{+30} \pm 3) \text{ MeV}. \quad (20)$$

We do not fit our data to a function that allows for the opening of a threshold in another channel, because there is no obvious change in the shape of the dipion mass spectrum near 1 GeV, and we do not have enough data to specify the parameters of such a fit.

V. CONCLUSIONS

We present an updated result for the first measurement of

$$\mathcal{B}(D_s^+ \rightarrow f_0(980)e^+\nu)\mathcal{B}(f_0 \rightarrow \pi^+\pi^-) = (0.20 \pm 0.03 \pm 0.01)\%. \quad (21)$$

Assuming a simple pole model for the form factor, $|f_+(q^2)|$, we estimate the pole mass as $(1.7_{-0.7}^{+4.5} \pm 0.2)$ GeV.

In the final state the only hadron present is the f_0 . This provides a particularly clean environment that allows us to measure the mass and width as $(977_{-9}^{+11} \pm 1)$ MeV, and $(91_{-22}^{+30} \pm 3)$ MeV, respectively.

We update our measurement of

$$\mathcal{B}(D_s^+ \rightarrow \phi e^+\nu) = (2.36 \pm 0.23 \pm 0.13)\%. \quad (22)$$

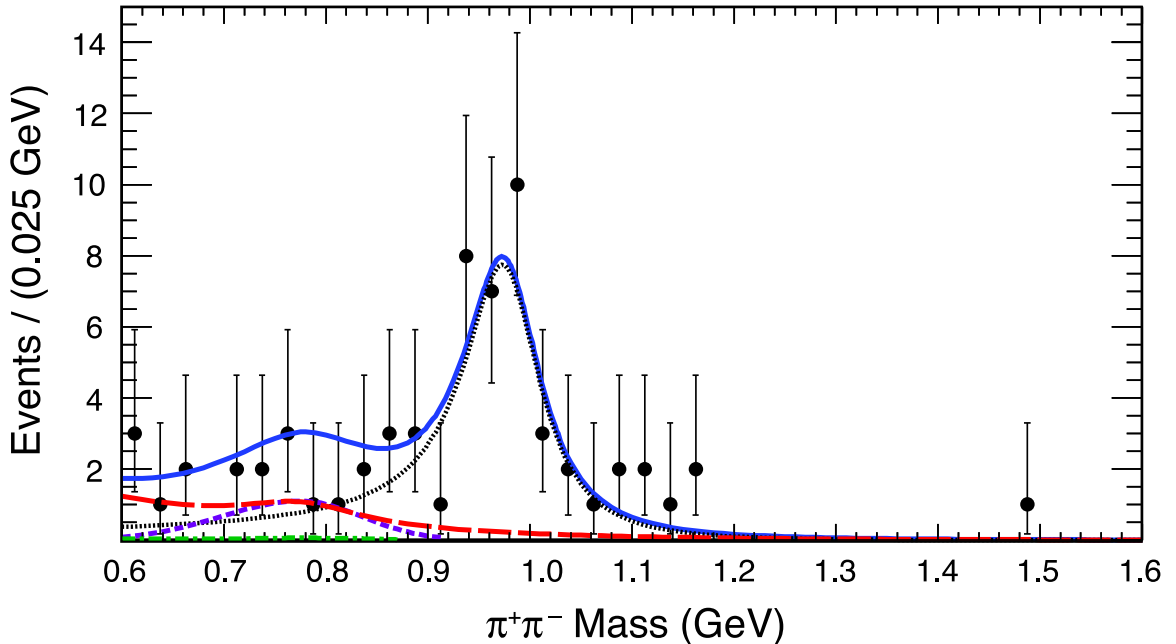


FIG. 9: Invariant mass distribution for $\pi^+\pi^-$ in the semileptonic mode. The dotted curve shows the signal shape, a relativistic Breit-Wigner modified by phase space and form factor effects, the long-dashed distributions are from sidebands of the D_s^- candidate invariant mass distributions while the shorter dashed curve in shows the $\eta' \rightarrow \pi^+\pi^-\gamma$ background level, and the dashed-dot curve shows the (small) $\pi^+\pi^-\pi^0$ background from $\phi e^+\nu$. The solid curve shows the total.

We measure the ratio of decay rates at $q^2 = 0$

$$R_{f/\phi} \equiv \frac{\frac{d\Gamma}{dq^2}(D_s^+ \rightarrow f_0(980)e^+\nu, f_0 \rightarrow \pi^+\pi^-) |_{q^2=0}}{\frac{d\Gamma}{dq^2}(D_s^+ \rightarrow \phi e^+\nu, \phi \rightarrow K^+K^-) |_{q^2=0}} = (42 \pm 11)\% . \quad (23)$$

This ratio has been predicted by Stone and Zhang to equal [6]

$$R_{f/\phi} = \frac{\Gamma(B_s \rightarrow J/\psi f_0, f_0 \rightarrow \pi^+\pi^-)}{\Gamma(B_s \rightarrow J/\psi \phi, \phi \rightarrow K^+K^-)} . \quad (24)$$

Our measurement indicates that the $B_s \rightarrow J/\psi f_0$ channel may indeed be a useful place to measure CP violation in the B_s system in that the rate can be $\sim 40\%$ that for $J/\psi \phi$ mode, especially since an angular analysis is not necessary, as the $J/\psi f_0$ mode is a CP eigenstate.

Finally, we show that viewing the f_0 as a quark-antiquark state results in a mixing angle between $\bar{s}s$ and light quarks whose cosine is close to unity, consistent with other estimates [19]. Thus, the large rate into $\pi\pi$ may be the result of part of the wave-function being a four-quark state [2].

VI. ACKNOWLEDGMENTS

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APPENDIX

We briefly summarize the relationships between different rates and form factors in D_s decays into a vector particle (ϕ) here. There are more Lorentz invariants possible than in the pseudoscalar to scalar case, and three independent form factors, V , A_1 , and A_2 are involved in the case of $e^+\nu$ decays. The matrix element can be written as the sum of a vector plus axial-vector current [14]:

$$\langle \phi | \bar{s} \gamma_\mu c | D_s \rangle = 2 \frac{V(q^2)}{M_{D_s} + M_\phi} \epsilon_{\mu\nu\alpha\beta} p_{D_s}^\nu p_\phi^\alpha \epsilon^{*\beta}, \quad \text{and} \quad (25)$$

$$\begin{aligned} \langle \phi | \bar{s} \gamma_\mu \gamma_5 c | D_s \rangle &= i (M_{D_s} + M_\phi) \left(\epsilon_\mu^* - \frac{\epsilon^* \cdot q}{q^2} q_\mu \right) A_1(q^2) \\ &\quad - i \frac{\epsilon^* \cdot q}{M_{D_s} + M_\phi} \left(P^\mu - \frac{M_{D_s}^2 - M_\phi^2}{q^2} q^\mu \right) A_2(q^2) \end{aligned} \quad (26)$$

where ϵ^* is a polarization of the final state meson. (An additional term proportional to A_0 has been dropped as it gets multiplied by the square of the electron mass.)

First of all the differential decay rate can be written in terms of the helicity amplitudes as [20]

$$\frac{d\Gamma}{dq^2} = \frac{G_F^2 |V_{cs}|^2 p_\phi q^2}{96\pi^3 M_{D_s}^2} (|H_+|^2 + |H_-|^2 + |H_0|^2). \quad (27)$$

The helicity amplitudes are expressed in terms of three form factors in the limit of zero lepton mass as [21]

$$H_0(q^2) = \frac{1}{2M_\phi \sqrt{q^2}} \left[(M_{D_s}^2 - M_\phi^2 - q^2) (M_{D_s} + M_\phi) A_1(q^2) - 4 \frac{M_{D_s}^2 p_\phi^2}{M_{D_s} + M_\phi} A_2(q^2) \right], \quad (28)$$

$$H_\pm(q^2) = (M_{D_s} + M_\phi) A_1(q^2) \mp \frac{2M_{D_s} p_\phi}{(M_{D_s} + M_\phi)} V(q^2). \quad (29)$$

The most elementary parameterization of the form factors considers the q^2 dependence as the behavior of a simple pole in the t -channel, yielding

$$V(q^2) = \frac{V(0)}{1 - q^2/M_V^2}; \quad A_i(q^2) = \frac{A_i(0)}{1 - q^2/M_A^2}, \quad \text{for } i = 1 \text{ and } 2. \quad (30)$$

The masses can be taken as $M_V = M_{D_s^*}$, and $M_A = M_{D_{s1}}$, although they can, in principle, be extracted from the data. The BaBar collaboration has assumed the value for M_V and measured $M_A = (2.28_{-0.18}^{+0.23} \pm 0.18)$ GeV; they also determined at $q^2 = 0$, $r_V = V(0)/A_1(0) = 1.849 \pm 0.060 \pm 0.095$; $r_2 = A_2(0)/A_1(0) = 0.763 \pm 0.071 \pm 0.065$ [3].

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- [1] C. Amsler *et al.*, Phys. Lett. B **667**, 1 (2008).
- [2] A. H. Fariborz, R. Jora, and J. Schechter, “Global Aspects of the Scalar Meson Puzzle,” arXiv:0902.2825v1 [hep-ph]; A. V. Anisovich *et al.*, “The Riddle of the $f_0(980)$ and $a_0(980)$: Are They the Quark-Antiquark States?” arXiv:hep-ph/0508260 and references contained therein; T. Branz, T. Gutsche, V. E. Lyubovitskij, Phys. Rev. D **79**, 014035 (2009) arXiv:0812.094 [hep-ph]; V. Baru, J. Haidenbauer, C. Hanhart, Yu. Kalashnikova, A. Kudryavtsev, Phys. Lett. B **586**, 53 (2004) arXiv:hep-ph/0308129.
- [3] B. Aubert *et al.* (BaBar), Phys. Rev. D **78**, 051101 (2008), arXiv:0807.1599v1 [hep-ex].
- [4] J. Yelton *et al.* (CLEO), “Absolute Branching Fraction Measurements for Exclusive D_s Semileptonic Decays,” arXiv:0903.0601v1 [hep-ex].
- [5] M. Artuso, E. Barberio, and S. Stone, PMC Physics A **3**, 3, (2009), arXiv:0902.3743v3 [hep-ph].
- [6] S. Stone and L. Zhang, Phys. Rev. D **79**, 074024 (2009), arXiv:0812.2832v4 [hep-ph].
- [7] D. Peterson *et al.*, Nucl. Instrum. and Meth. **A478**, 142 (2002); Y. Kubota *et al.* (CLEO Collaboration), Nucl. Instrum. and Meth. **A320**, 66 (1992).
- [8] M. Artuso *et al.*, Nucl. Instrum. Meth. **A554**, 147 (2005) [physics/0506132].
- [9] D. Cronin-Hennessy *et al.* (CLEO), “Measurement of Charm Production Cross Sections in e^+e^- Annihilation at Energies between 3.97 and 4.26 GeV,” arXiv:0801.3418 [hep-ex].
- [10] J. P. Alexander *et al.* (CLEO), Phys. Rev. D **79**, 052001 (2009), arXiv:0901.1216v2 [hep-ex].
- [11] T. E. Coan *et al.* (CLEO), Phys. Rev. Lett. **95**, 181802 (2005).
- [12] T. Skwarnicki, “A Study of the Radiative Cascade Transitions Between the Upsilon-Prime and Upsilon Resonances,” DESY F31-86-02 (thesis, unpublished) (1986); see also Eq. (4) in P. Rubin *et al.* (CLEO), Phys. Rev. D **73**, 112005 (2006), [hep-ex/0604043].
- [13] H.G. Dosch, E. M. Ferreira, F. S. Navarra, and M. Nielsen, Phys. Rev. D **65**, 114002 (2002).
- [14] M. Artuso, B. Meadows, and A. A. Petrov, Ann. Rev. Nucl. Part. Sci. **58**, 249 (2008), arXiv:0802.2934v1 [hep-ph].
- [15] N. Isgur, D. Scora, B. Grinstein and M. B. Wise, Phys. Rev. D **39**, 799 (1989). See Appendix D.
- [16] S. Dobbs *et al.* (CLEO), Phys. Rev. D **76**, 112001 (2007), see Appendix B.
- [17] T. M. Aliev and M. Savci, “Semileptonic decays of pseudoscalar mesons to scalar f_0 meson,” arXiv:hep-ph/0701108. A previous work, I. Bediaga, F. S. Navarra and M. Nielsen, Phys. Lett. B **579**, 59 (2006), predicts a smaller width with a 50% error, which leads to a smaller value of $\cos^2 \theta$ with a large error.
- [18] M. Ablikim *et al.* (BES), Phys. Rev. D **70**, 092002 (2004); M. Ablikim *et al.* (BES), Phys. Rev. D **72**, 092002 (2005).
- [19] See for example, B. El-Bennich, O. Leitner, J.-P. Dedonder, B. Loiseau, Phys. Rev. D **79**, 076004 (2009) arXiv:0810.5771 [hep-ph], and references contained therein; H. Y. Cheng, Phys. Rev. D **67**, 034024 (2003); W. Wang, Y. L. Shen, Y. Li, and C. D. Lu, Phys. Rev. D **74**, 114010 (2006).
- [20] F. J. Gilman and R. L. Singleton, Phys. Rev. D **41**, 142 (1990).
- [21] J. M. Cline, W. F. Palmer, and G. Kramer, Phys. Rev. D **40**, 793 (1989); G. Kopp *et al.*, Phys. Lett. B **256**, 271 (1991); P. Ball, Phys. Rev. D **48**, 3190 (1993).