

# Superconducting RF Cavity Measurement Formulae for an Exponential Decayed Pulse Incident Power

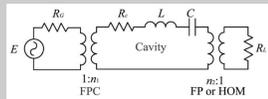
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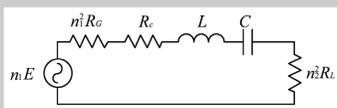
## Abstract

Experimental method for evaluating a Superconducting RF (SRF) cavity performance is through low power and high power measurements without a beam load. The equations for square incident power pulse are the most popular formulae for the pulsed measurements. In practice, incident power may not be exactly square pulse. To understand cavity behavior and performance more accurately, in this paper, the SRF cavity's measurement equations for an exponential-decayed pulsed incident power are developed from a series equivalent circuit. The analytical result can be directly compared with the experimental data of SNS cavities obtained from the Cryomodule Test Facility (CMTF) at Jefferson Lab.

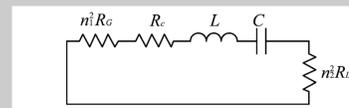
## EQUIVALENT CIRCUIT MODEL



(a) Equivalent circuit of a two-port cavity coupling system.



(b) Alternative form of the circuit of (a) for RF Switch on.



(c) Alternative form of the circuit of (a) for RF Switch off.

Figure 1: The equivalent circuits in transient state.

$R_G$  and  $R_L$  are the impedances of the signal source and the load, respectively.

$R_c$ ,  $L$ ,  $C$  are the resistance, inductance and capacitance of the SRF cavity.

$E$  is the equivalent signal source voltage with frequency of  $\omega$ .

The cavity voltage  $V_c$  is:  $V_c(\omega, t) = \frac{I(\omega, t)}{\omega C} = dE_{acc}(\omega, t)$

The cavity's emitted power  $P_e$ , dissipated power  $P_d$ , and transmitted power  $P_t$  are:

$$\begin{cases} P_e(\omega_0, t) = n_1^2 R_G |I(\omega_0, t)|^2 \\ P_d(\omega_0, t) = R_c |I(\omega_0, t)|^2 \\ P_t(\omega_0, t) = n_2^2 R_L |I(\omega_0, t)|^2 \end{cases}$$

Cavity's intrinsic quality factor  $Q_0$ :  $Q_0 = \omega_0 U(\omega_0, t) / P_d(\omega_0, t) = \omega_0 L / R_c = 1 / \omega_0 C R_c$   $\omega_0 = \sqrt{1/LC}$

External quality factors  $Q_e$ :  $Q_e = \omega_0 U(\omega_0, t) / P_t(\omega_0, t) = \omega_0 L / n_2^2 R_L = 1 / \omega_0 C n_2^2 R_L$   $Q_t$  of

Field Probe:  $Q_e$  of the FPC:  $Q_e = \omega_0 U(\omega_0, t) / P_e(\omega_0, t) = \omega_0 L / n_1^2 R_G = 1 / \omega_0 C n_1^2 R_G$

Cavity FPC coupling coefficient  $\beta_e$  and the FP coupling coefficient  $\beta_t$  are defined as:

$$\begin{cases} \beta_e = Q_0 / Q_e = P_e / P_d = n_1^2 R_G / R_c \\ \beta_t = Q_0 / Q_t = P_t / P_d = n_2^2 R_L / R_c \end{cases}$$

Cavity loaded quality factor  $Q_L$ :  $Q_L = \frac{\omega_0 L}{R_c + n_1^2 R_G + n_2^2 R_L} = \frac{Q_0}{1 + \beta_e + \beta_t}$

Differential equation relating current  $I(\omega, t)$ , and  $L \frac{dI(\omega, t)}{dt} + (n_1^2 R_G + R_c + n_2^2 R_L) I(\omega, t) + \frac{1}{C} \int I(\omega, t) dt = n_1 E(\omega, t)$  voltage  $E(\omega, t)$  in the circuit is:

## CAVITY MEASUREMENT FORMULAE

For a pulsed incident power which is exponentially decay with the form  $P_{in}(t) = P_{in} \exp(-2at)$ , the cavity measurement formulae are:

### A. RF Switch On

$$P_d(\omega, t) = \frac{4\beta_e(\alpha^2 + \omega^2) \exp(-2at) \left[ 1 + \exp\left[2a - \frac{\omega_0}{Q_L} t\right] - 2 \exp\left[\left(a - \frac{\omega_0}{2Q_L}\right) t\right] \cos(\Delta\omega t) \right]}{Q_0^2 \omega^2 \left[ \left(\frac{\alpha^2}{\omega\omega_0} - \frac{\alpha}{\omega Q_L} + \frac{\omega_0}{\omega} - \frac{\omega}{\omega_0}\right)^2 + \left(\frac{1}{Q_L} - \frac{2\alpha}{\omega_0}\right)^2 \right]} P_{in}$$

$$P_t(\omega, t) = \frac{4\beta_t \beta_e (\alpha^2 + \omega^2) \exp(-2at) \left[ 1 + \exp\left[2a - \frac{\omega_0}{Q_L} t\right] - 2 \exp\left[\left(a - \frac{\omega_0}{2Q_L}\right) t\right] \cos(\Delta\omega t) \right]}{Q_0^2 \omega^2 \left[ \left(\frac{\alpha^2}{\omega\omega_0} - \frac{\alpha}{\omega Q_L} + \frac{\omega_0}{\omega} - \frac{\omega}{\omega_0}\right)^2 + \left(\frac{1}{Q_L} - \frac{2\alpha}{\omega_0}\right)^2 \right]} P_{in} = \beta_t P_d(\omega, t)$$

$$P_e(\omega, t) = \frac{4\beta_e^2 (\alpha^2 + \omega^2) \exp(-2at) \left[ 1 + \exp\left[2a - \frac{\omega_0}{Q_L} t\right] - 2 \exp\left[\left(a - \frac{\omega_0}{2Q_L}\right) t\right] \cos(\Delta\omega t) \right]}{Q_0^2 \omega^2 \left[ \left(\frac{\alpha^2}{\omega\omega_0} - \frac{\alpha}{\omega Q_L} + \frac{\omega_0}{\omega} - \frac{\omega}{\omega_0}\right)^2 + \left(\frac{1}{Q_L} - \frac{2\alpha}{\omega_0}\right)^2 \right]} P_{in} = \beta_e P_d(\omega, t)$$

$$E_{acc}(\omega, t) = \frac{4\beta_e (R/G) Q_0 (\alpha^2 + \omega^2) P_{in} \exp(-2at) \left[ 1 + \exp\left[2a - \frac{\omega_0}{Q_L} t\right] - 2 \exp\left[\left(a - \frac{\omega_0}{2Q_L}\right) t\right] \cos(\Delta\omega t) \right]}{d^2 Q_0 \omega^2 \left[ \left(\frac{\alpha^2}{\omega\omega_0} - \frac{\alpha}{\omega Q_L} + \frac{\omega_0}{\omega} - \frac{\omega}{\omega_0}\right)^2 + \left(\frac{1}{Q_L} - \frac{2\alpha}{\omega_0}\right)^2 \right]} = \sqrt{\frac{(R/G) \omega_0^2 Q_0 P_{in}(t)}{d^2 \omega^2}} = \sqrt{\frac{(R/G) \omega_0^2 Q_0 P_{in}(t)}{d^2 \omega^2}} = \sqrt{\frac{(R/G) \omega_0^2 Q_0 P_{in}(t)}{d^2 \omega^2}}$$

$$U(\omega, t) = \frac{2\beta_e (\alpha^2 + \omega^2) Q_0 (1 + \beta_e + \beta_t) \exp(-2at) \left[ 1 + \exp\left[2a - \frac{\omega_0}{Q_L} t\right] - 2 \exp\left[\left(a - \frac{\omega_0}{2Q_L}\right) t\right] \cos(\Delta\omega t) \right]}{Q_0^2 \omega_0 \omega^2 \left[ \left(\frac{\alpha^2}{\omega\omega_0} - \frac{\alpha}{\omega Q_L} + \frac{\omega_0}{\omega} - \frac{\omega}{\omega_0}\right)^2 + \left(\frac{1}{Q_L} - \frac{2\alpha}{\omega_0}\right)^2 \right]} P_{in} = \left(1 + \frac{\omega_0^2}{\omega^2}\right) \frac{Q_0 (1 + \beta_e + \beta_t)}{2\omega_0} P_d(\omega, t)$$

The reflected power  $P_r$ :  $P_r(\omega, t) = P_{in} \exp(-2at) - P_d(\omega, t) - P_t(\omega, t) - \frac{dU}{dt}(\omega, t)$

### B. RF Switch Off

Supposing the pulse length is  $\tau_0$ , after the RF switch off, then:

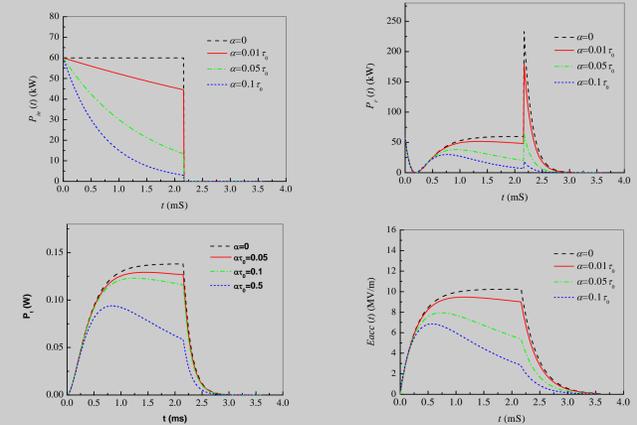
$$\begin{cases} P_d(t) = P_d(\omega_0, \tau_0) \exp\left[-\frac{\omega_0}{Q_L} (t - \tau_0)\right] \\ P_t(t) = P_t(\omega_0, \tau_0) \exp\left[-\frac{\omega_0}{Q_L} (t - \tau_0)\right] = \beta_t P_d(t) \\ P_e(t) = P_e(\omega_0, \tau_0) \exp\left[-\frac{\omega_0}{Q_L} (t - \tau_0)\right] = \beta_e P_d(t) \end{cases}$$

$$U(t) = \int_{-\infty}^{+\infty} \frac{dU}{d\tau}(\tau) d\tau = L |I(\omega, \tau_0)|^2 \exp\left[-\frac{\omega_0}{Q_L} (t - \tau_0)\right] = \frac{Q_0 (1 + \beta_e + \beta_t)}{\omega_0} P_d(t)$$

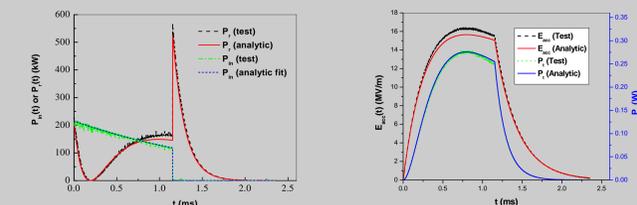
$$P_r(t) = -P_d(t) - P_t(t) - \frac{dU}{dt}(t) = P_e(t) = \beta_e P_d(t)$$

$$E_{acc}(t) = \sqrt{\frac{4\beta_e (R_{acc}/Q) Q_0 P_{in}}{d^2 (1 + \beta_e + \beta_t)}} \left[ 1 - \exp\left[-\frac{\omega_0}{2Q_L} \tau_0\right] \right] \exp\left[-\frac{\omega_0}{2Q_L} (t - \tau_0)\right] = E_{acc}(\omega, \tau_0) \exp\left[-\frac{\omega_0}{2Q_L} (t - \tau_0)\right]$$

$$= \sqrt{\frac{(R_{acc}/Q) Q_0 P_{in}(t)}{d}} = \sqrt{\frac{(R_{acc}/Q) Q_0 P_{in}(t)}{d}} = \sqrt{\frac{(R_{acc}/Q) Q_0 P_{in}(t)}{d}} = \sqrt{\frac{(R_{acc}/Q) Q_0 P_{in}(t)}{d}}$$



The SNS medium beta cavity's reflected power  $P_r(t)$ , transmitted power  $P_t(t)$ , accelerating gradient  $E_{acc}(t)$  and incident power  $P_{in}(t)$  versus different decay  $\alpha$ . Here incident power  $P_{in} = 60$  kW with pulse length of  $15\tau_0$  (here  $\tau_0 = Q_L / \omega_0 = 0.1435$  ms).



The SNS medium beta cavity M082 test data in JLab CMTF and comparison with the analytic fitting data in the  $\alpha = 0.0382\omega_0/Q_L = 266s^{-1}$  decay rate of incident pulse. The pulse length is 1.153 ms. The initial amplitude of incident power is  $P_{in} = 216$  kW. The  $E_{acc}(\text{Test})$  was on-line calculated by one-port measurement equations.

## CONCLUSION

The two-port RF cavity's equations, developed by a series lumped equivalent circuit, can accurately describe the cavity operation conditions. These equations can be used to exactly measure all superconducting RF cavity's parameters, except cavity intrinsic quality factor  $Q_0$  under special conditions. The two-port equations can be simplified into one-port cavity equations.

Cavity stored energy change  $dU/dt$  is the intrinsic cause of the reflected power  $P_r$ , emitted power  $P_e$  and transmitted power  $P_t$  change. In practice, the incident power  $P_{in}$  is not always the standard constant power or square wave pulse power. To evaluate cavity performance precisely, the cavity measurement equations for standard incident power must be modified. This paper modified the two-port measurement equations as example. The discussion and equations in this paper demonstrate the LCR equivalent circuit can be utilized to model and predict the electrical behavior of the superconducting RF cavity.