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Twist duality for flux backgrounds of type II and heterotic String Theory from Generalized Complex Geometry

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arXiv:0903.0633 by D. A., R. Minasian, M. Petrini

01/05/2009, Cornell University, NY, USA

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• Dualities are a major tool to study string theory.

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• Dualities are a major tool to study string theory. Different string theories related by a web of dualities

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Flux backgrounds for compactifications.

 ${\scriptstyle \bullet}\,$  String theory  $\rightarrow$  Real world low energy physics

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### Flux backgrounds for compactifications.

• String theory  $\rightarrow$  Real world low energy physics  $\hookrightarrow$  Compactification:  $\mathbb{R}^{3,1} \times \mathcal{M}_{internal}$ 

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Try to preserve the minimal amount of 4D supersymmetry

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- String theory  $\rightarrow$  Real world low energy physics
  - $\hookrightarrow$  Compactification:  $\mathbb{R}^{3,1} \times \mathcal{M}_{internal}$
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  - $\hookrightarrow$  Usually led to  $\mathcal{M}_{internal} = Calabi-Yau (CY)$

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### Flux backgrounds for compactifications

String theory → Real world low energy physics
 → Compactification: ℝ<sup>3,1</sup> × M<sub>internal</sub>
 Try to preserve the minimal amount of 4D supersymmetry
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 Effective theory on it ⇒ the moduli problem...

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   → On what M<sub>internal</sub> to compactify ?

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- For type II SUGRA with fluxes, mathematical characterization of  $\mathcal{M}_{internal}$  given in terms of Generalized Complex Geometry (GCG):

math. DG/0209099 by N. Hitchin, math. DG/0401221 by M. Gualtieri

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### Flux backgrounds for compactifications.

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 $\mathcal{M}_{internal}$  preserving at least  $\mathcal{N} = 1$  are Generalized CY (GCY)

hep-th/0406137, hep-th/0505212 by M. Graña, R. Minasian, M. Petrini, A. Tomasiello

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# • More generally: GCG is a natural set-up to study type II SUSY flux backgrounds.

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• More generally: GCG is a natural set-up to study type II SUSY flux backgrounds. Useful for non-geometry, effective actions,

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• More generally: GCG is a natural set-up to study type II SUSY flux backgrounds.

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Propose Twist duality in this language.

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 $\hookrightarrow$  Introduce some GCG objects, and then perform Twist duality to map two solutions.

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• Review appearance of GCG in type II SUSY flux backgrounds.

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• Review appearance of GCG in type II SUSY flux backgrounds. More elements of GCG in type II.

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- Heterotic backgrounds treatment.

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The SUSY conditions for a SUGRA vacuum

• Type II SUGRA:  $\mathcal{N}_{10D} = 2$ Spectrum:  $g, \phi, H = dB, F_p, \psi^{1,2}_{\mu}, \lambda^{1,2}$ 

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• Find some SUSY Minkowski flux vacuum of it:

• vacuum: equations of motion (e.o.m.)

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- vacuum: equations of motion (e.o.m.)
- fluxes: Bianchi Identities (BI)
- SUSY: supersymmetry conditions

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# Generalized Complex Geometry and type II flux compactifications

The SUSY conditions for a SUGRA vacuum

• Type II SUGRA:  $\mathcal{N}_{10D} = 2$ Spectrum:  $g, \phi, H = dB, F_p, \ \psi_{\mu}^{1,2}, \lambda^{1,2}$ Compactification:  $\mathbb{R}^{3,1} \times \mathcal{M}_{\text{internal}}$  $\hookrightarrow$  Metric Ansatz :  $ds_{(10)}^2 = e^{2A(y)} \ \eta_{\mu\nu} dx^{\mu} dx^{\nu} + g_{\mu\nu}(y) dy^{\mu} dy^{\nu}$ 

- vacuum: equations of motion (e.o.m.)
- fluxes: Bianchi Identities (BI)
- SUSY: supersymmetry conditions
- Other constraints...

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## • Find some SUSY Minkowski flux vacuum of it:

- vacuum: equations of motion (e.o.m.)
- fluxes: Bianchi Identities (BI)
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One can show:

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SUSY conditions + BI  $\Rightarrow$  e.o.m.

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The SUSY conditions for a SUGRA vacuum

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## • Find some SUSY Minkowski flux vacuum of it:

- vacuum: equations of motion (e.o.m.)
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- SUSY: supersymmetry conditions
- Other constraints...

One can show:

SUSY conditions + BI  $\Rightarrow$  e.o.m. Main focus: solve the SUSY conditions

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## • SUSY conditions :

$$0 = \delta \psi_{\mu} = D_{\mu}\epsilon + \frac{1}{4}H_{\mu}\mathcal{P}\epsilon + \frac{1}{16}e^{\phi}\sum_{n} I_{2n}^{\mu}\gamma_{\mu}\mathcal{P}_{n}\epsilon$$
$$0 = \delta\lambda = \left(\partial \!\!\!/ \phi + \frac{1}{2}I_{n}^{\mu}\mathcal{P}\right)\epsilon + \frac{1}{8}e^{\phi}\sum_{n}(-1)^{2n}(5-2n)I_{2n}^{\mu}\mathcal{P}_{n}\epsilon$$

SUSY parameters of type II SUGRA:  $\epsilon = (\epsilon^1, \epsilon^2)$  .

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## • SUSY conditions :

$$= \delta \psi_{\mu} = D_{\mu}\epsilon + \frac{1}{4}H_{\mu}\mathcal{P}\epsilon + \frac{1}{16}e^{\phi}\sum_{n} I\!\!\!/_{2n}\gamma_{\mu}\mathcal{P}_{n}\epsilon$$
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SUSY parameters of type II SUGRA:  $\epsilon=(\epsilon^1,\epsilon^2)$  . Fluxes in the SUSY conditions

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## • SUSY conditions : CY condition

$$0 = \delta \psi_{\mu} = D_{\mu} \epsilon$$
$$0 = \delta \lambda = \left( \partial \phi \right) \epsilon$$

SUSY parameters of type II SUGRA:  $\epsilon=(\epsilon^1,\epsilon^2)$  . Fluxes in the SUSY conditions

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### • SUSY conditions :

$$= \delta \psi_{\mu} = D_{\mu}\epsilon + \frac{1}{4}H_{\mu}\mathcal{P}\epsilon + \frac{1}{16}e^{\phi}\sum_{n} I_{2n}^{r}\gamma_{\mu}\mathcal{P}_{n}\epsilon$$
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SUSY parameters of type II SUGRA:  $\epsilon = (\epsilon^1, \epsilon^2)$ . Fluxes in the SUSY conditions  $\Rightarrow$  GCG rewriting.

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### • SUSY conditions :

$$= \delta \psi_{\mu} = D_{\mu} \epsilon + \frac{1}{4} H_{\mu} \mathcal{P} \epsilon + \frac{1}{16} e^{\phi} \sum_{n} \mathcal{F}_{2n} \gamma_{\mu} \mathcal{P}_{n} \epsilon$$
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• SUSY parameters decomposed

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### • SUSY conditions :

$$= \delta \psi_{\mu} = D_{\mu} \epsilon + \frac{1}{4} H_{\mu} \mathcal{P} \epsilon + \frac{1}{16} e^{\phi} \sum_{n} \mathcal{F}_{2n} \gamma_{\mu} \mathcal{P}_{n} \epsilon$$
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SUSY parameters of type II SUGRA:  $\epsilon = (\epsilon^1, \epsilon^2)$ . Fluxes in the SUSY conditions  $\Rightarrow$  GCG rewriting.

• SUSY parameters decomposed, for  $\mathcal{N}_{4D} = 1$ :

$$\begin{aligned} \epsilon^1 &= \zeta \otimes \eta^1 + c.c. \\ \epsilon^2 &= \zeta \otimes \eta^2 + c.c. \end{aligned}$$

an internal pair  $(\eta^1, \eta^2)$ .

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### • SUSY conditions :

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an internal pair  $(\eta^1, \eta^2)$ . For a consistent reduction,

Two globally defined non-vanishing spinors on  $\mathcal{M}_{internal}$ :  $\eta^1, \eta^2$ 

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SUSY parameters of type II SUGRA:  $\epsilon = (\epsilon^1, \epsilon^2)$ . Fluxes in the SUSY conditions  $\Rightarrow$  GCG rewriting.

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Two globally defined non-vanishing spinors on  $\mathcal{M}_{\text{internal}}$ :  $\eta^1, \eta^2$ Consider in the following  $\eta^1 = \eta^2$  (SU(3) structure).

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## • SUSY conditions on $\mathcal{M}_{internal}$ :

	• SUSY conditions on $\mathcal{M}_{\text{internal}}$ :	
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## • SUSY conditions on $\mathcal{M}_{internal}$ :

No flux

With flux (GCG)

$$\begin{array}{l} D_m\eta^1=0\\ D_m\eta^2=0 \end{array}$$
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### • SUSY conditions on $\mathcal{M}_{internal}$ :

No flux

With flux (GCG)

 $D_m \eta^1 = 0$  $D_m \eta^2 = 0$ 

 $\eta^{1,2}$  spinors on T

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### • SUSY conditions on $\mathcal{M}_{internal}$ :

No flux

With flux (GCG)

 $D_m \eta^1 = 0$  $D_m \eta^2 = 0$ 

 $\eta^{1,2}$  spinors on TCY

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## • SUSY conditions on $\mathcal{M}_{internal}$ :

No flux

With flux (GCG)

 $D_m \eta^1 = 0$  $D_m \eta^2 = 0$  $\eta^{1,2} \text{ spinors on } T$ 

CY

$$\begin{array}{l} {\rm d}(e^{3A} \ \Psi_1) = 0 \\ {\rm d}(e^{2A} \ {\rm Re}(\Psi_2)) = 0 \\ {\rm d}(e^{4A} \ {\rm Im}(\Psi_2)) = e^{4A}e^{-B} \ * \lambda(\sum_p F_p) \end{array}$$

SUSY conditions

• SUSY conditions on  $\mathcal{M}_{internal}$ :

 $\begin{array}{l} D_m\eta^1=0\\ D_m\eta^2=0 \end{array}$  $\eta^{1,2}$  spinors on T

No flux

CY

With flux (GCG)

$$d(e^{3A} \Psi_1) = 0$$
  

$$d(e^{2A} \operatorname{Re}(\Psi_2)) = 0$$
  

$$d(e^{4A} \operatorname{Im}(\Psi_2)) = e^{4A}e^{-B} * \lambda(\sum_p F_p)$$
  

$$\Psi_{1,2} \text{ spinors on } T \oplus T^*$$

	• SUSY conditions on	$\mathcal{M}_{ ext{internal}}$ :
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Introduction GCG and type II SUSY conditions $T \oplus T^*$ bundle Twist Dual solutions Heterotic	$D_m \eta^1 = 0$ $D_m \eta^2 = 0$ $\eta^{1,2} \text{ spinors on } T$ $CY$	$\begin{aligned} \mathbf{d}(e^{3A} \ \Psi_1) &= 0\\ \mathbf{d}(e^{2A} \ \mathrm{Re}(\Psi_2)) &= 0\\ \mathbf{d}(e^{4A} \ \mathrm{Im}(\Psi_2)) &= e^{4A} e^{-B} \ast \lambda(\sum_p F_p)\\ \Psi_{1,2} \ \mathrm{spinors \ on} \ T \oplus T^*\\ \mathrm{GCY} \end{aligned}$

	• SUSY conditions on $\mathcal{M}_{internal}$ :		
David ANDRIOT	No flux	With flux (GCG)	
	$D_m \eta^1 = 0$	$d(e^{3A} \Psi_1) = 0$	
SUSY conditions	$D_m \eta^2 = 0$	$d(e^{2A} \operatorname{Re}(\Psi_2)) = 0$	
$T \oplus T^*$ bundle Twist	$n^{1,2}$ spinors on T	$d(e^{\dots} \operatorname{Im}(\Psi_2)) = e^{\dots}e^{-D^*} * \lambda(\sum_p F_p)$ $\Psi_{1,2} \text{ spinors on } T \oplus T^*$	
	CY	GCY	
	• First define different	ial forms (with $\eta_{-} = (\eta_{+})^*$ ) for $SU(3)$ :	
	$J_{\mu\nu} = -i\eta_{\perp}^{\dagger}\gamma_{\mu}$	$\mu_{\mu}\eta_{\pm}$ , $\Omega_{\mu\nu\rho} = -i\eta_{\pm}^{\dagger}\gamma_{\mu\nu\rho}\eta_{\pm}$ .	

	• SUSY conditions on $\mathcal{M}_{internal}$ :	
David ANDRIOT	No flux	With flux (GCG)
Introduction GCG and type II SUSY conditions $T \oplus T^*$ bundle	$D_m\eta^1=0\ D_m\eta^2=0$	$d(e^{3A} \Psi_1) = 0$ $d(e^{2A} \operatorname{Re}(\Psi_2)) = 0$ $d(e^{4A} \operatorname{Im}(\Psi_2)) = e^{4A}e^{-B} * \lambda(\sum_n F_n)$
Twist Dual solutions Heterotic Conclusions	$\eta^{1,2}$ spinors on $T$ CY • First define different $J_{\mu\nu} = -i\eta^{\dagger}\gamma_{\mu}$	$\Psi_{1,2} \text{ spinors on } T \oplus T^*$ GCY ial forms (with $\eta = (\eta_+)^*$ ) for $SU(3)$ : $\Omega_{\mu\nu\rho} = -i\eta^{\dagger} \gamma_{\mu\nu\rho} \eta_+ .$
	SUSY conditions ca	n be expressed in terms of forms !

	• SUSY conditions on	$\mathcal{M}_{ ext{internal}}$ :
David ANDRIOT	No flux	With flux (GCG)
Introduction GCG and type II SUSY conditions $T \oplus T^*$ bundle Twist Dual solutions Heterotic Conclusions	$D_m \eta^1 = 0$ $D_m \eta^2 = 0$ $\eta^{1,2} \text{ spinors on } T$ $CY$ • First define different	$\begin{aligned} \mathbf{d}(e^{3A} \ \Psi_1) &= 0\\ \mathbf{d}(e^{2A} \ \mathrm{Re}(\Psi_2)) &= 0\\ \mathbf{d}(e^{4A} \ \mathrm{Im}(\Psi_2)) &= e^{4A}e^{-B} * \lambda(\sum_p F_p)\\ \Psi_{1,2} \text{ spinors on } T \oplus T^*\\ \mathrm{GCY} \end{aligned}$ tial forms (with $\eta = (\eta_+)^*$ ) for $SU(3)$ :
	$J_{\mu u}=-i\eta^{\dagger}_{+}\gamma_{\mu}$	$_{\mu\nu}\eta_+ , \qquad \Omega_{\mu\nu\rho} = -i\eta^\dagger \gamma_{\mu\nu\rho}\eta_+ .$
	<ul> <li>SUSY conditions can be expressed in terms of forms !</li> <li>• Further define the bi-spinors: Φ<sub>+</sub> = η<sup>1</sup><sub>+</sub> ⊗ η<sup>2†</sup><sub>+</sub>, Φ<sub>-</sub> = η<sup>1</sup><sub>+</sub> ⊗ η<sup>2†</sup><sub>-</sub></li> </ul>	

	• SUSY conditions on $\mathcal{M}_{internal}$ :	
David ANDRIOT	No flux	With flux (GCG)
Introduction GCG and type II SUSY conditions $T \oplus T^*$ bundle	$D_m \eta^1 = 0$ $D_m \eta^2 = 0$	$d(e^{3A} \Psi_1) = 0$ $d(e^{2A} \operatorname{Re}(\Psi_2)) = 0$ $d(e^{4A} \operatorname{Im}(\Psi_2)) = e^{4A}e^{-B} * \lambda(\Sigma - F_2)$
Twist Dual solutions Heterotic Conclusions	$\eta^{1,2}$ spinors on $T$ CY • First define different	$\Psi_{1,2} \text{ spinors on } T \oplus T^*$ $\frac{\Psi_{1,2} \text{ spinors on } T \oplus T^*}{\text{GCY}}$ $\text{fial forms (with } \eta = (\eta_+)^*) \text{ for } SU(3):$
	$J_{\mu\nu} = -i\eta_{+}^{\dagger}\gamma_{\mu\nu}\eta_{+} , \qquad \Omega_{\mu\nu\rho} = -i\eta_{-}^{\dagger}\gamma_{\mu\nu\rho}\eta_{+} .$ SUSY conditions can be expressed in terms of forms ! • Further define the bi-spinors:	

 $\Phi_+ = \eta_+^1 \otimes \eta_+^{2\dagger}, \quad \Phi_- = \eta_+^1 \otimes \eta_-^{2\dagger}$ 

 $\Phi_{\pm} = \text{polyforms}, \text{ for } SU(3):$ 

	• SUSY conditions on $\mathcal{M}_{internal}$ :		
David ANDRIOT	No flux	With flux (GCG)	
Introduction GCG and type II SUSY conditions $T \oplus T^*$ bundle Twist Dual solutions Heterotic Conclusions	$D_m \eta^1 = 0$ $D_m \eta^2 = 0$ $\eta^{1,2} \text{ spinors on } T$ $CY$ • First define different $J_{\mu\nu} = -i\eta_+^{\dagger}\gamma_{\mu}$ SUSY conditions can	$\begin{aligned} d(e^{3A} \ \Psi_1) &= 0\\ d(e^{2A} \ \mathrm{Re}(\Psi_2)) &= 0\\ d(e^{4A} \ \mathrm{Im}(\Psi_2)) &= e^{4A} e^{-B} \ast \lambda(\sum_p F_p)\\ \Psi_{1,2} \text{ spinors on } T \oplus T^*\\ \mathrm{GCY} \end{aligned}$ ial forms (with $\eta = (\eta_+)^*$ ) for $SU(3)$ : $_{\mu\nu}\eta_+ \ , \qquad \Omega_{\mu\nu\rho} &= -i\eta^{\dagger}\gamma_{\mu\nu\rho}\eta_+ \ . \end{aligned}$ n be expressed in terms of forms !	
	• Further define the b $\Phi_+ = \pi$ $\Phi_\pm = \text{polyforms, for } \Phi_+, \Phi, \Phi, \Phi, \Phi, \Phi, \Phi, \Phi, \Phi$	i-spinors: $\eta_{+}^{1} \otimes \eta_{+}^{2\dagger},  \Phi_{-} = \eta_{+}^{1} \otimes \eta_{-}^{2\dagger}$ or $SU(3)$ : $N_{+} = e^{i\theta} \qquad e^{-iJ}$ $N_{-} = -i \qquad \Omega_{3}$	

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$ \begin{array}{c} \begin{array}{c} \text{Introduction} \\ \text{GCG and type II} \\ \text{SUSY conditions} \\ T \oplus T^* \text{ bundle} \end{array} \end{array} \begin{array}{c} D_m \eta^1 = 0 \\ D_m \eta^2 = 0 \\ \text{d}(e^{2A} \operatorname{Re}(\Psi_2)) = 0 \\ \text{d}(e^{4A} \operatorname{Im}(\Psi_2)) = e^{4A} e^{-B} * \lambda(\sum_{i=1}^{n} e^{-iA} e^{-iA}) \end{array} \end{array} $	$F_n$ )
Twist Dual solutions Heterotic Conclusions $\eta^{1,2} \text{ spinors on } T \qquad \Psi_{1,2} \text{ spinors on } T \oplus T^*$ GCY <b>o</b> First define differential forms (with $\eta = (\eta_+)^*$ ) for $SU(\eta_+)^*$	(3):
SUSY  conditions can be expressed in terms of forms !	

• Further define the bi-spinors:  $\Phi_+ = \eta^1_+ \otimes \eta^{2\dagger}_+, \quad \Phi_- = \eta^1_+ \otimes \eta^{2\dagger}_-$ 

$$\begin{split} \Phi_{\pm} &= \text{polyforms, for } SU(3): \\ \Psi_{+} &= e^{-\phi} e^{-B} \; \Phi_{+}/N_{+} = e^{i\theta} e^{-\phi} e^{-B} e^{-iJ} \\ \Psi_{-} &= e^{-\phi} e^{-B} \; \Phi_{-}/N_{-} = -i e^{-\phi} e^{-B} \Omega_{3} \end{split}$$

	• SUSY conditions on	$\mathcal{M}_{\mathrm{internal}}$ :
David ANDRIOT	No flux	With flux (GCG)
Introduction GCG and type II SUSY conditions $T \oplus T^*$ bundle Twist Dual solutions Heterotic Conclusions	$D_m \eta^1 = 0$ $D_m \eta^2 = 0$ $\eta^{1,2} \text{ spinors on } T$ $CY$ • First define different $J_{\mu\nu} = -i\eta_+^{\dagger} \gamma_{\mu\nu}$	$\begin{aligned} & \operatorname{d}(e^{3A} \ \Psi_1) = 0 \\ & \operatorname{d}(e^{2A} \operatorname{Re}(\Psi_2)) = 0 \\ & \operatorname{d}(e^{4A} \ \operatorname{Im}(\Psi_2)) = e^{4A} e^{-B} \ast \lambda(\sum_p F_p) \\ & \Psi_{1,2} \text{ spinors on } T \oplus T^* \\ & \operatorname{GCY} \end{aligned}$ tial forms (with $\eta = (\eta_+)^*$ ) for $SU(3)$ : $_{\mu\nu}\eta_+ \ , \qquad \Omega_{\mu\nu\rho} = -i\eta^{\dagger}\gamma_{\mu\nu\rho}\eta_+ \ . \end{aligned}$
	SUSY conditions can • Further define the b $\Phi_+ = r$	in be expressed in terms of forms ! $\eta_{+}^{1} \otimes \eta_{+}^{2\dagger},  \Phi_{-} = \eta_{+}^{1} \otimes \eta_{-}^{2\dagger}$ $\Phi_{-} = \eta_{+}^{1} \otimes \eta_{-}^{2\dagger}$

$$\begin{split} \Psi_{\pm} &= p \, {\rm ory}\, {\rm forms, \, for \, b \, c \, (5),} \\ \Psi_{\pm} &= e^{-\phi} e^{-B} \, \Phi_{\pm}/N_{\pm} = e^{i\theta} e^{-\phi} e^{-B} e^{-iJ} \\ \Psi_{\pm} &= e^{-\phi} e^{-B} \, \Phi_{\pm}/N_{\pm} = -i e^{-\phi} e^{-B} \Omega_{3} \end{split}$$

SUSY conditions can be expressed in terms of  $\Psi_{\pm}$  !

	• SUSY conditions on	$\mathcal{M}_{\mathrm{internal}}$ :
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Introduction GCG and type II SUSY conditions $T \oplus T^*$ bundle Twist Dual solutions Heterotic Conclusions	$D_m \eta^1 = 0$ $D_m \eta^2 = 0$ $\eta^{1,2} \text{ spinors on } T$ $CY$ • First define different $J_{\mu\nu} = -i\eta_+^{\dagger} \gamma_{\mu}$ SUSY conditions ca	$d(e^{3A} \Psi_1) = 0$ $d(e^{2A} \operatorname{Re}(\Psi_2)) = 0$ $d(e^{4A} \operatorname{Im}(\Psi_2)) = e^{4A}e^{-B} * \lambda(\sum_p F_p)$ $\Psi_{1,2} \text{ spinors on } T \oplus T^*$ GCY dial forms (with $\eta = (\eta_+)^*$ ) for $SU(3)$ : $u_{\mu\nu}\eta_+ , \qquad \Omega_{\mu\nu\rho} = -i\eta^{\dagger}\gamma_{\mu\nu\rho}\eta_+ .$ In be expressed in terms of forms !
	• Further define the b $\Phi_{\pm} = r$ $\Phi_{\pm} = \text{polyforms, fo}$	$egin{aligned} & ext{i-spinors:} \ \eta_+^1\otimes\eta_+^{2\dagger}, & \Phi=\eta_+^1\otimes\eta^{2\dagger} \ & ext{ir} \ SU(3): \end{aligned}$

$$\begin{split} \Psi_{+} &= e^{-\phi} e^{-B} \ \Phi_{+}/N_{+} = e^{i\theta} e^{-\phi} e^{-B} e^{-iJ} \\ \Psi_{-} &= e^{-\phi} e^{-B} \ \Phi_{-}/N_{-} = -i e^{-\phi} e^{-B} \Omega_{3} \end{split}$$

SUSY conditions can be expressed in terms of  $\Psi_{\pm}$  ! Spinors on  $T \oplus T^*$ ,

	• SUSY conditions on $\mathcal{M}_{internal}$ :	
David ANDRIOT	No flux	With flux (GCG)
Introduction GCG and type II SUSY conditions $T \oplus T^*$ bundle Twist Dual solutions Heterotic Conclusions	$D_m \eta^1 = 0$ $D_m \eta^2 = 0$ $\eta^{1,2} \text{ spinors on } T$ $CY$ • First define different	$d(e^{3A} \Psi_1) = 0$ $d(e^{2A} \operatorname{Re}(\Psi_2)) = 0$ $d(e^{4A} \operatorname{Im}(\Psi_2)) = e^{4A}e^{-B} * \lambda(\sum_p F_p)$ $\Psi_{1,2} \text{ spinors on } T \oplus T^*$ $GCY$ tial forms (with $\eta = (\eta_+)^*$ ) for $SU(3)$ :
	$J_{\mu\nu} = -i\eta_{+}^{*}\gamma_{\mu}$ SUSY conditions ca • Further define the b $\Phi_{+} = \eta_{-}$ $\Phi_{+} = polyforms for$	$M_{\mu\nu\rho} = -i\eta_{-}^{\dagger}\gamma_{\mu\nu\rho}\eta_{+} .$ In be expressed in terms of forms ! i-spinors: $\eta_{+}^{1} \otimes \eta_{+}^{2\dagger},  \Phi_{-} = \eta_{+}^{1} \otimes \eta_{-}^{2\dagger}$ or $SU(3)$ .

$$\begin{split} \Psi_{\pm} &= \text{polytorms, for } SU(3). \\ \Psi_{+} &= e^{-\phi}e^{-B} \; \Phi_{+}/N_{+} = e^{i\theta}e^{-\phi}e^{-B}e^{-iJ} \\ \Psi_{-} &= e^{-\phi}e^{-B} \; \Phi_{-}/N_{-} = -ie^{-\phi}e^{-B}\Omega_{3} \end{split}$$

SUSY conditions can be expressed in terms of  $\Psi_{\pm}$  ! Spinors on  $T \oplus T^*$ , use of  $\Psi_{\pm} \Rightarrow$  GCG interpretations !

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## The generalized tangent bundle

GCG considers the fibration:

 $T^* \quad \hookrightarrow \quad E$ 

E: the generalized tangent bundle.

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# The generalized tangent bundle

GCG considers the fibration:

 $\begin{array}{cccc} \Gamma^* & \hookrightarrow & E \\ & & \downarrow & E \text{: the ge} \\ & & - \end{array}$ 

 $E\colon$  the generalized tangent bundle.

Locally:  $T \oplus T^*$ .

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## The generalized tangent bundle

GCG considers the fibration:

 $\begin{array}{cccc} T^* & \hookrightarrow & E \\ & & \downarrow & E \text{: the generalized tangent bundle.} \\ & & T \end{array}$ 

Locally:  $T \oplus T^*$ . Sections: generalized vectors:

$$X = v + \xi = \begin{pmatrix} v \\ \xi \end{pmatrix} , \ v \in T \ , \ \xi \in T^*$$

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## The generalized tangent bundle

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This non-trivial fibration: given by the 2-form B. Geometrical interpretation of B: "connective structure" of a gerbe.

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 $\hookrightarrow \omega = d\Lambda$  is "a gauge transformation" from  $U_{\alpha}$  to  $U_{\beta}$ .

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• E is equipped with a natural metric  $\eta$  :

$$\eta(X,X) = i_v \xi = v^{\mu} \xi_{\mu} \quad \Leftrightarrow \quad X^T \eta X = \frac{1}{2} \begin{pmatrix} v & \xi \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} v \\ \xi \end{pmatrix}$$

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$$P = e^B \begin{pmatrix} a & 0 \\ 0 & a^{-T} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ B & 1 \end{pmatrix} \begin{pmatrix} a & 0 \\ 0 & a^{-T} \end{pmatrix} = \begin{pmatrix} a & 0 \\ Ba & a^{-T} \end{pmatrix}$$

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Twist duality  $\subset G_{\text{geom}}$ .

•  $\mathcal{M}_{\text{internal}}$ : a bundle with base  $\mathcal{B}$  along  $dx^{\mu}$ , fiber  $\mathcal{F}$  along  $dy^{m}$ , and connection  $A_{\mu}^{m}$ :

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 $ds^{2} = g_{\mu\nu} dx^{\mu} dx^{\nu} + g_{mn} (dy^{m} + A^{m}_{\ \rho} dx^{\rho}) (dy^{n} + A^{n}_{\ \sigma} dx^{\sigma}) .$ 

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In the  $(dx^{\mu}, dy^{m})$  basis: vielbeins and metric:

$$e = \begin{pmatrix} e_{\mathcal{B}} & 0\\ e_{\mathcal{F}}A & e_{\mathcal{F}} \end{pmatrix} \qquad g = e^{T}e = \begin{pmatrix} g_{\mathcal{B}} + A^{T}g_{\mathcal{F}}A & A^{T}g_{\mathcal{F}}\\ g_{\mathcal{F}}A & g_{\mathcal{F}} \end{pmatrix}$$
$$g_{\mathcal{B}} = e_{\mathcal{B}}^{T}e_{\mathcal{B}} , \qquad g_{\mathcal{F}} = e_{\mathcal{F}}^{T}e_{\mathcal{F}} .$$

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### • On E, B plays the role of a connection

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On E, B plays the role of a connection

 → in GCG, we introduce generalized vielbeins/metric:

$$\mathcal{E} = \begin{pmatrix} e & 0 \\ -e^{-T}B & e^{-T} \end{pmatrix}$$
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Transformation under O(d, d):

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$$\mathcal{E} \mapsto \mathcal{E}' = \mathcal{E}O , \ \mathcal{H} \mapsto \mathcal{H}' = O^T \mathcal{H}O , \ e^{\phi} \mapsto e^{\phi'} = e^{\phi} \left( \frac{\det(g')}{\det(g)} \right)^T$$

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 $\mathcal{H}$  looks like an object in T-duality on torus.

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• Pure spinors: 
$$\Psi_{\pm} = \frac{1}{N_{\pm}} e^{-\phi} e^{-B} \eta_{\pm}^{1} \otimes \eta_{\pm}^{2\dagger}$$

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$$= \begin{pmatrix} e & 0 \\ -e^{-T}B & e^{-T} \end{pmatrix} , \qquad \mathcal{H} = \mathcal{E}^{T}\mathcal{E} = \begin{pmatrix} g - Bg^{-1}B & Bg^{-1} \\ -g^{-1}B & g^{-1} \end{pmatrix}$$
$$\eta = \mathcal{E}^{T} \begin{pmatrix} 0 & \mathbb{I} \\ \mathbb{I} & 0 \end{pmatrix} \mathcal{E} , \qquad \mathcal{H} = \mathcal{E}^{T} \begin{pmatrix} \mathbb{I} & 0 \\ 0 & \mathbb{I} \end{pmatrix} \mathcal{E} .$$

Transformation under O(d, d):

 $\mathcal{E} \mapsto \mathcal{E}' = \mathcal{E}O, \ \mathcal{H} \mapsto \mathcal{H}' = O^T \mathcal{H}O, \ e^{\phi} \mapsto e^{\phi'} = e^{\phi} \left(\frac{\det(g')}{\det(g)}\right)^{\frac{1}{4}}.$ 

 $\mathcal{H}$  looks like an object in T-duality on torus. Same transformation under T-duality group  $O(d_{\mathcal{F}}, d_{\mathcal{F}})$ . Here  $d = d_{\mathcal{M}} > d_{\mathcal{F}} \Rightarrow$  T-duality natural in GCG; we will consider a "generalization of T-duality".

• Pure spinors: 
$$\Psi_{\pm} = \frac{1}{N_{\pm}} e^{-\phi} e^{-B} \eta_{\pm}^1 \otimes \eta_{\pm}^{2\dagger} .$$

Majorana-Weyl Spin(d, d) spinors on E (locally  $T \oplus T^*$ )

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Geometrical info. contained in J,  $\Omega_3$ ,  $(B, \phi)$  (SU(3)).

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## On generalized vielbeins ${\mathcal E}$

• We consider 
$$O \in G_{geom} \subset O(d, d)$$
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with separation of  $\mathcal{B}$  and  $\mathcal{F}$ .

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$$\hookrightarrow e' = eA , B' = A^T B A - A^T C , e^{\phi'} = e^{\phi} |\det(A)|^{\frac{1}{2}} .$$

### Concrete example:

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$$e = \begin{pmatrix} e_{\mathcal{B}} & 0\\ 0 & e_{\mathcal{F}} \end{pmatrix} , \ \mathrm{d}s^2 = g_{\mu\nu} \mathrm{d}x^{\mu} \mathrm{d}x^{\nu} + g_{mn} (\mathrm{d}y^m + 0) (\mathrm{d}y^n + 0) ,$$

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The twist transformation gives:

$$\begin{aligned} e' &= \begin{pmatrix} e_{\mathcal{B}}A_{\mathcal{B}} & 0\\ (e_{\mathcal{F}}A_{\mathcal{F}})(A_{\mathcal{F}}^{-1}A_{\mathcal{C}}) & e_{\mathcal{F}}A_{\mathcal{F}} \end{pmatrix}\\ \mathrm{d}s^2 &= g'_{\mu\nu}\mathrm{d}x^{\mu}\mathrm{d}x^{\nu} + g'_{mn}(\mathrm{d}y^m + A'^m_{\rho}\mathrm{d}x^{\rho})(\mathrm{d}y^n + A'^n_{\sigma}\mathrm{d}x^{\sigma}) \ ,\\ A' &= A_{\mathcal{F}}^{-1}A_{\mathcal{C}} \ , \ g'_{\mu\nu} = (A_{\mathcal{B}}^T g_{\mathcal{B}} A_{\mathcal{B}})_{\mu\nu} \ , \ g'_{mn} = (A_{\mathcal{F}}^T g_{\mathcal{F}} A_{\mathcal{F}})_{mn} \ . \end{aligned}$$

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$$B' = -\begin{pmatrix} \tilde{C}_{\mathcal{B}} & -\tilde{C}_{\mathcal{C}}^T \\ \tilde{C}_{\mathcal{C}} & \tilde{C}_{\mathcal{F}} \end{pmatrix}$$

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## On pure spinors $\Psi_\pm$

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## On pure spinors $\Psi_{\pm}$

Metric  $\eta \Rightarrow$  Clifford algebra Cliff(d, d) on E:  $\{\Gamma^m, \Gamma^n\} = \{\Gamma_m, \Gamma_n\} = 0$ ,  $\{\Gamma^m, \Gamma_n\} = \delta_n^m \quad m, \ n = 1 \dots d$ .

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### On pure spinors $\Psi_{\pm}$

Metric  $\eta \Rightarrow$  Clifford algebra Cliff(d, d) on E:  $\{\Gamma^m, \Gamma^n\} = \{\Gamma_m, \Gamma_n\} = 0$ ,  $\{\Gamma^m, \Gamma_n\} = \delta_n^m \quad m, \ n = 1 \dots d$ .

Action on  $\Psi_{\pm}$ : wedges, contractions:  $\Gamma^n = dx^n \wedge , \ \Gamma_m = \iota_{\partial_m}$ .

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 $O = e^{-\frac{1}{4}\Theta_{MN}\sigma^{MN}}$ ,  $M, N = 1 \dots d + d$ .

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$$\sigma^{MN} = [\Gamma^M, \Gamma^N] \quad ,$$

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GL(d) transformation (previous A):

$$O_a = e^{-\frac{1}{4}(a^m{}_n[\Gamma_m,\Gamma^n] - a_m{}^n[\Gamma^m,\Gamma_n])} = \dots = \frac{1}{\sqrt{\det A}} e^{a^m{}_n \mathrm{d}x^n \wedge \iota_{\partial_m}}$$

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B-transform (previous C):

$$O_B = e^{-\frac{1}{2}B_{mn}\Gamma^{mn}} = e^{-\frac{1}{2}B_{mn}\mathrm{d}x^m \wedge \mathrm{d}x^n}$$

Previous twist transformation:

David ANDRIOT  $\mathcal{E} \mapsto \mathcal{E}' = \mathcal{E}O \ , \ O = \begin{pmatrix} A & 0 \\ C & A^{-T} \end{pmatrix}$ 

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Previous twist transformation:

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$$\mathcal{E} \mapsto \mathcal{E}' = \mathcal{E}O \ , \ O = \begin{pmatrix} A & 0 \\ C & A^{-T} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ B & 1 \end{pmatrix} \begin{pmatrix} A & 0 \\ 0 & A^{-T} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -B' & 1 \end{pmatrix}$$

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$$\mathcal{E} \mapsto \mathcal{E}' = \mathcal{E}O , \ O = \begin{pmatrix} A & 0 \\ C & A^{-T} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ B & 1 \end{pmatrix} \begin{pmatrix} A & 0 \\ 0 & A^{-T} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -B' & 1 \end{pmatrix}$$
  
becomes:

 $\Psi \mapsto O_f \Psi' \ , \ O_f = \frac{1}{\sqrt{\det A}} e^{-B'} e^{a^m{}_n \mathrm{d} x^n \wedge \iota_{\partial_m}} e^B \ .$ 

$$\mathcal{E} \mapsto \mathcal{E}' = \mathcal{E}O , \ O = \begin{pmatrix} A & 0 \\ C & A^{-T} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ B & 1 \end{pmatrix} \begin{pmatrix} A & 0 \\ 0 & A^{-T} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -B' & 1 \end{pmatrix}$$
  
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 $\Psi \mapsto O_f \Psi'$ ,  $O_f = \frac{1}{\sqrt{\det A}} e^{-B'} e^{a^m {}_n \mathrm{d} x^n \wedge \iota_{\partial m}} e^B$ .

We consider a further phase transformation:

$$O_c^{\pm} = e^{i\theta_c^{\pm}} O_f \; .$$

 $\mathcal{E} \mapsto \overline{\mathcal{E}'} = \mathcal{E}O \ , \ O = \begin{pmatrix} A & 0 \\ C & A^{-T} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ B & 1 \end{pmatrix} \begin{pmatrix} A & 0 \\ 0 & A^{-T} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -B' & 1 \end{pmatrix}$ becomes:

 $\Psi \mapsto O_f \Psi'$ ,  $O_f = \frac{1}{\sqrt{\det A}} e^{-B'} e^{a^m} e^{Ax^n \wedge \iota_{\partial_m}} e^B$ .

We consider a further phase transformation:

$$O_c^{\pm} = e^{i\theta_c^{\pm}} O_f$$
 .

Check the twist on  $\Psi_+$ : do we have:

$$\Psi_{+} = e^{i\theta} e^{-\phi} e^{-B} e^{-iJ} \longrightarrow \Psi'_{+} = e^{i(\theta+\theta_{c}^{+})} e^{-\phi'} e^{-B'} e^{-iJ'} ,$$
  
$$\Psi_{-} = -ie^{-\phi} e^{-B} \Omega_{3} \longrightarrow \Psi'_{-} = -ie^{i\theta_{c}^{-}} e^{-\phi'} e^{-B'} \Omega'_{3} .$$

On pure spinors

 $\mathcal{E} \mapsto \mathcal{E}' = \mathcal{E}O \ , \ O = \begin{pmatrix} A & 0 \\ C & A^{-T} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ B & 1 \end{pmatrix} \begin{pmatrix} A & 0 \\ 0 & A^{-T} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -B' & 1 \end{pmatrix}$ becomes:

On pure spinors

$$\Psi \mapsto O_f \ \Psi' \ , \ O_f = rac{1}{\sqrt{\det A}} e^{-B'} \ e^{a^m{}_n \mathrm{d} x^n \wedge \, \iota_{\partial_m}} \ e^B \ .$$

We consider a further phase transformation:

$$O_c^{\pm} = e^{i\theta_c^{\pm}} O_f \; .$$

Check the twist on  $\Psi_{\pm}$ : do we have:

$$\begin{split} \Psi_{+} &= e^{i\theta} e^{-\phi} e^{-B} e^{-iJ} \quad \longrightarrow \quad \Psi'_{+} = e^{i(\theta+\theta_{c}^{+})} e^{-\phi'} e^{-B'} e^{-iJ'} ,\\ \Psi_{-} &= -i e^{-\phi} e^{-B} \Omega_{3} \quad \longrightarrow \quad \Psi'_{-} = -i e^{i\theta_{c}^{-}} e^{-\phi'} e^{-B'} \Omega'_{3} .\\ \text{ransform: } \checkmark, \end{split}$$

David

 $\mathcal{E} \mapsto \mathcal{E}' = \mathcal{E}O, \quad O = \begin{pmatrix} A & 0 \\ C & A^{-T} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ B & 1 \end{pmatrix} \begin{pmatrix} A & 0 \\ 0 & A^{-T} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -B' & 1 \end{pmatrix}$ becomes:

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We consider a further phase transformation:

$$O_c^{\pm} = e^{i\theta_c^{\pm}} O_f \; .$$

Check the twist on  $\Psi_+$ : do we have:

 $\Psi_{+} = e^{i\theta} e^{-\phi} e^{-B} e^{-iJ} \longrightarrow \Psi'_{+} = e^{i(\theta + \theta_{c}^{+})} e^{-\phi'} e^{-B'} e^{-iJ'} ,$  $\Psi_{-} = -ie^{-\phi}e^{-B}\Omega_{3} \quad \longrightarrow \quad \Psi' = -ie^{i\theta_{c}^{-}}e^{-\phi'}e^{-B'}\Omega'_{2} \; .$ *B*-transform:  $\checkmark$ , dilaton transform:  $\checkmark$ ,

On pure spinors

David ANDRIOT  $\mathcal{E} \mapsto \mathcal{E}' = \mathcal{E}O , \ O = \begin{pmatrix} A & 0 \\ C & A^{-T} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ B & 1 \end{pmatrix} \begin{pmatrix} A & 0 \\ 0 & A^{-T} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -B' & 1 \end{pmatrix}$ becomes:

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David ANDRIOT  $\mathcal{E} \mapsto \mathcal{E}' = \mathcal{E}O , \quad O = \begin{pmatrix} A & 0 \\ C & A^{-T} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ B & 1 \end{pmatrix} \begin{pmatrix} A & 0 \\ 0 & A^{-T} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -B' & 1 \end{pmatrix}$ becomes:

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$$\Psi \mapsto O_f \ \Psi' \ , \ O_f = \frac{1}{\sqrt{\det A}} e^{-B'} \ e^{a^m \, {}_n \mathrm{d} x^n \wedge \, \iota_{\partial_m}} \ e^B \ .$$

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left with the "A-transform"  $O_{a}$ : action on  $J, \Omega_{3}$ .  
should transform the metric and provide a connection

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On pure spinors

 $\mathcal{E} \mapsto \mathcal{E}' = \mathcal{E}O, \quad O = \begin{pmatrix} A & 0 \\ C & A^{-T} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ B & 1 \end{pmatrix} \begin{pmatrix} A & 0 \\ 0 & A^{-T} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -B' & 1 \end{pmatrix}$ becomes:

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We consider a further phase transformation:

$$O_c^{\pm} = e^{i\theta_c^{\pm}} O_f \; .$$

Check the twist on  $\Psi_{\pm}$ : do we have:

$$\begin{split} \Psi_{+} &= e^{i\theta} \, e^{-\phi} e^{-B} \, e^{-iJ} \quad \longrightarrow \quad \Psi'_{+} = e^{i(\theta+\theta_{c}^{+})} \, e^{-\phi'} \, e^{-B'} \, e^{-iJ'} \ , \\ \Psi_{-} &= -i e^{-\phi} e^{-B} \Omega_{3} \quad \longrightarrow \quad \Psi'_{-} = -i e^{i\theta_{c}^{-}} \, e^{-\phi'} \, e^{-B'} \Omega'_{3} \ . \\ B\text{-transform: } \checkmark, \text{ dilaton transform: } \checkmark, \text{ phases: } \checkmark. \\ &\hookrightarrow \text{ left with the "A-transform"} \, O_{a} : \text{ action on } J, \, \Omega_{3}. \\ \text{It should transform the metric and provide a connection.} \\ \text{Particular case: } \mathcal{F} = T^{2}, \text{ only } g_{\mathcal{F}} \text{ transformation, and provide} \\ \text{a holomorphic connection } \alpha: \end{split}$$

David

 $\mathcal{E} \mapsto \mathcal{E}' = \mathcal{E}O, \quad O = \begin{pmatrix} A & 0 \\ C & A^{-T} \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ B & 1 \end{pmatrix} \begin{pmatrix} A & 0 \\ 0 & A^{-T} \end{pmatrix} \begin{pmatrix} 1 & 0 \\ -B' & 1 \end{pmatrix}$ becomes:

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$$J = J_{\mathcal{B}} + \frac{i}{2} g_{z\overline{z}} \, dz \wedge d\overline{z} \, \mapsto \, J' = J_{\mathcal{B}} + \frac{i}{2} g'_{z\overline{z}} \, (dz + \alpha) \wedge (d\overline{z} + \overline{\alpha})$$
  
$$\Omega_3 = \sqrt{g_{z\overline{z}}} \, \omega_{\mathcal{B}} \wedge dz \, \mapsto \, \Omega'_3 = \sqrt{g'_{z\overline{z}}} \, \omega_{\mathcal{B}} \wedge (dz + \alpha) \, .$$

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$$\Omega_{3} = \sqrt{g_{z\overline{z}}} \, \omega_{\mathcal{B}} \wedge dz \, \mapsto \, \Omega'_{3} = \sqrt{g'_{z\overline{z}}} \, \omega_{\mathcal{B}} \wedge (dz + \alpha) \, . \, \checkmark$$

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$$\Omega_3 = \sqrt{g_{z\overline{z}}} \, \omega_{\mathcal{B}} \wedge dz \, \mapsto \, \Omega'_3 = \sqrt{g'_{z\overline{z}}} \, \omega_{\mathcal{B}} \wedge (dz+\alpha) \, . \, \checkmark$$

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Remarks (work in progress):

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Remarks (work in progress):

• Twist does not change  $T\mathcal{M}_{internal}$  structure group (SU(3))

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### Remarks (work in progress):

• Twist does not change  $T\mathcal{M}_{internal}$  structure group (SU(3))T-duality can change it...

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Remarks (work in progress):

• Twist does not change  $T\mathcal{M}_{internal}$  structure group (SU(3))T-duality can change it...

• RR flux transform?

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Remarks (work in progress):

- Twist does not change  $T\mathcal{M}_{\text{internal}}$  structure group (SU(3))T-duality can change it...
- RR flux transform? In GCG, defined through SUSY from  $\Psi_{\pm}...$

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Remarks (work in progress):

- Twist does not change  $T\mathcal{M}_{internal}$  structure group (SU(3))T-duality can change it...
- RR flux transform? In GCG, defined through SUSY from Ψ<sub>±</sub>...
   Definition of the transformed RR as new solutions of SUSY.

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# Mapping type IIB solutions

Constraints to generate solutions

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# Mapping type IIB solutions

Constraints to generate solutions

Consider a solution of the SUSY conditions.

 $\begin{array}{l} {\rm d}(e^{3A}\Psi_1)=0\\ {\rm d}(e^{2A}\operatorname{Re}\Psi_2)=0\\ {\rm d}(e^{4A}\operatorname{Im}\Psi_2)=e^{4A}e^{-B}*\lambda(F)=R \;. \end{array}$ 

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# Mapping type IIB solutions

## Constraints to generate solutions

Consider a solution of the SUSY conditions. Perform a twist:  $\Psi'_{\pm} = e^{i\theta_c^{\pm}} O_f \Psi_{\pm}$ 

```
\begin{split} & \mathrm{d}(e^{3A}\Psi_1) = 0 \\ & \mathrm{d}(e^{2A}\operatorname{Re}\Psi_2) = 0 \\ & \mathrm{d}(e^{4A}\operatorname{Im}\Psi_2) = e^{4A}e^{-B}*\lambda(F) = R \; . \end{split}
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# Mapping type IIB solutions

### Constraints to generate solutions

Consider a solution of the SUSY conditions. Perform a twist:  $\Psi'_{\pm} = e^{i\theta_c^{\pm}} O_f \Psi_{\pm} \Rightarrow \text{Get a new solution}?$ 

$$\begin{array}{ll} \mathrm{d}(e^{3A}\Psi_1) = 0 & \qquad \mathrm{d}(e^{3A}\Psi_1') = 0 \\ \mathrm{d}(e^{2A}\operatorname{Re}\Psi_2) = 0 & \Rightarrow & \mathrm{d}(e^{2A}\operatorname{Re}\Psi_2') = 0 \\ \mathrm{d}(e^{4A}\operatorname{Im}\Psi_2) = e^{4A}e^{-B}*\lambda(F) = R \ . & \qquad \mathrm{d}(e^{4A}\operatorname{Im}\Psi_2') = R' \end{array}$$

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Develop in terms of  $\Psi_{1,2}, R, O_f, \theta_c^{\pm}$ 

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 $O_f, \theta_c^+$  constrained with respect to the first solution  $\Psi_{1,2}, R$ .

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## Mapping type IIB solutions

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 $O_f$ ,  $\theta_c^+$  constrained with respect to the first solution  $\Psi_{1,2}$ , R. Note that automatically satisfied for ordinary T-duality.

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$$s_{\theta_c^+} d(e^{2A} O_f) \ e^{2A} \operatorname{Re} \Psi_2 + c_{\theta_c^+} d(O_f) \ e^{4A} \operatorname{Im} \Psi_2 = R' - c_{\theta_c^+} O_f \ R \ .$$

 $O_f$ ,  $\theta_c^+$  constrained with respect to the first solution  $\Psi_{1,2}$ , R. Note that automatically satisfied for ordinary T-duality. Last equation: only definition of R':

$$R' = c_{\theta_c^+} O_f R + s_{\theta_c^+} d(e^{2A} O_f) e^{2A} \operatorname{Re} \Psi_2 + c_{\theta_c^+} d(O_f) e^{4A} \operatorname{Im} \Psi_2 .$$

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 $O_f$ ,  $\theta_c^+$  constrained with respect to the first solution  $\Psi_{1,2}$ , R. Note that automatically satisfied for ordinary T-duality. Last equation: only definition of R':

$$\begin{split} R' &= c_{\theta_c^+} O_f R + s_{\theta_c^+} \mathrm{d}(e^{2A} O_f) e^{2A} \operatorname{Re} \Psi_2 + c_{\theta_c^+} \mathrm{d}(O_f) e^{4A} \operatorname{Im} \Psi_2 \ . \\ \theta_c^+ &\neq 0, \text{ coordinate dependent } O_f \end{split}$$

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## Examples of type IIB dual solutions

$\mathcal{M}_{ ext{internal}}$	$T^6 = T^2 \times T^4$	

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## Examples of type IIB dual solutions

$\mathcal{M}_{ ext{internal}}$	$T^6 = T^2 \times T^4$	
$ds_6^2$	$e^{-2A}\mathrm{d}x_{\mathcal{B}}^2 + e^{-2A}\mathrm{d}x_{\mathcal{F}}^2$	

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Sources, $\theta$	$O3$ , $\theta = \frac{\pi}{2}$	$O5 \ // \ \mathcal{F} \ , \ \theta = 0$
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Sources, $\theta$	$O3$ , $\theta = \frac{\pi}{2}$	$O5 \ // \ \mathcal{F} \ , \ \theta = 0$
RR	$g_s F_5 = e^{4A} * d(e^{-4A}) , \ (F_3)$	$g_s F_3 = -e^{-4A} * \mathrm{d}(e^{2A}J)$
NSNS	$(H = g_s * F_3)$	
$e^{\phi}$	$g_s$	

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NSNS	$(H = g_s * F_3)$	0
$e^{\phi}$	$g_s$	$g_s e^{2A}$

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## Examples of type IIB dual solutions

### Solutions:

$\mathcal{M}_{ ext{internal}}$	$T^6 = T^2 \times T^4$	$\begin{array}{cccc} T^2 & \hookrightarrow & \mathcal{M} \\ & \downarrow \\ & T^4 \end{array}$
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NSNS	$(H = g_s * F_3)$	0
$e^{\phi}$	$g_s$	$g_s e^{2A}$

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## Examples of type IIB dual solutions

### Solutions:

$\mathcal{M}_{ ext{internal}}$	$T^6 = T^2 \times T^4$	$\begin{array}{cccc} T^2 & \hookrightarrow & \mathcal{M} \\ & & \downarrow \\ & & T^4 \end{array}$
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NSNS	$(H = g_s * F_3)$	0
$e^{\phi}$	$g_s$	$g_s e^{2A}$

Twist duality map?

 $A_{\mathcal{B}} = 1_4$ 

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### Solutions:

$\mathcal{M}_{ ext{internal}}$	$T^6 = T^2 \times T^4$	$\begin{array}{cccc} T^2 & \hookrightarrow & \mathcal{M} \\ & \downarrow \\ & T^4 \end{array}$
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Sources, $\theta$	$O3$ , $\theta = \frac{\pi}{2}$	$O5 \; // \; \mathcal{F} \;, \; \theta = 0$
RR	$g_s F_5 = e^{4A} * d(e^{-4A}) , \ (F_3)$	$g_s F_3 = -e^{-4A} * \mathbf{d}(e^{2A}J)$
NSNS	$(H = g_s * F_3)$	0
$e^{\phi}$	$g_s$	$g_s e^{2A}$

Twist duality map?

 $A_{\mathcal{B}} = 1_4 , \ A_{\mathcal{F}} = 1_2 \times e^{2A}$ 

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## Examples of type $\overline{\text{IIB dual solutions}}$

### Solutions:

$\mathcal{M}_{ ext{internal}}$	$T^6 = T^2 \times T^4$	$\begin{array}{cccc} T^2 & \hookrightarrow & \mathcal{M} \\ & & \downarrow \\ & & T^4 \end{array}$
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NSNS	$(H = g_s * F_3)$	0
$e^{\phi}$	$g_s$	$g_s e^{2A}$

Twist duality map?

$$A_{\mathcal{B}} = 1_4 \ , \ A_{\mathcal{F}} = 1_2 \times e^{2A} \ , \ A_{\mathcal{C}}{}^{I}{}_{\mu} = e^{2A} A^{I}_{\mu}$$

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### Solutions:

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Sources, $\theta$	$O3$ , $ heta=rac{\pi}{2}$	$O5 \; // \; \mathcal{F} \;, \;  heta = 0$
RR	$g_s F_5 = e^{4A} * d(e^{-4A}) , \ (F_3)$	$g_s F_3 = -e^{-4A} * \mathbf{d}(e^{2A}J)$
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Twist duality map?

$$A_{\mathcal{B}} = 1_4 , \ A_{\mathcal{F}} = 1_2 \times e^{2A} , \ A_{\mathcal{C}}{}^{I}{}_{\mu} = e^{2A} A^{I}_{\mu} , \ \theta^+_c = -\frac{\pi}{2}$$

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## Examples of type IIB dual solutions

### Solutions:

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 $\hookrightarrow \mathrm{Twist}\ \mathrm{duals}\ \checkmark$ 

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Explicit non-trivial fibration solutions?

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Explicit non-trivial fibration solutions? Among nilmanifolds: twisted tori, GCY...

### Among 34 nilmanifolds, only 5 non-trivial $T^2$ bundles.

David		
ANDRIOT		
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	Among 54 minimannoids, only 5 non-trivial $1^-$ bundles.
David	They have different topologies:
ANDRIOT	
$\rm GCG$ and type $\rm II$	
Examples	

	Among 34 nilm	anifolds, only	5 non	-trivia	TT bi	indles.	
David	They have diffe	rent topologie	es:				
ANDRIOT			$T^2$	$\hookrightarrow$	$\mathcal{M}$		
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$n \ 4.4 \ , \ n \ 4.7$			$T^2$	$\hookrightarrow$	$\mathcal{M} \ \downarrow \ T^4$		
n 4.5	$S^1$	$\hookrightarrow$	$\mathcal{M}_1 \ \downarrow \ T^2$	×	$S^1$	$\hookrightarrow$	$\mathcal{M}_2 \ \downarrow \ T^2$

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			$T^2$	$\hookrightarrow$	$\mathcal{M}$		
$n \ 4.4 \ , \ n \ 4.7$					$\downarrow \\ T^4$		
	$S^1$	$\hookrightarrow$	$\mathcal{M}_1$		$\frac{1}{S^1}$	$\hookrightarrow$	$\mathcal{M}_2$
$n \ 4.5$			$\downarrow$	×			$\downarrow$
		$T^2$	$T^2$	11			12
$n \ 4.6$		T	,	$\downarrow$ $T^3$	×	$S^1$	
				1			

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	$S^1$	$\hookrightarrow$	$\mathcal{M}_1$		$\frac{T^{4}}{S^{1}}$	$\hookrightarrow$	$\mathcal{M}_2$
n  4.5			$\downarrow \\ T^2$	×			$\stackrel{\downarrow}{T^2}$
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		~1		$T^3$		~	
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1 1				$T^4$			

They have precise curvatures  $F = d\alpha$ .

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Another solution on  $n \ 3.14 : S^1 \hookrightarrow \mathcal{M} \to \mathcal{M}_1$ .

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SUSY conditions in the presence of  $H \neq 0$  and Bianchi id.: Nucl. Phys. B 274 (1986) 253 by A. Strominger

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Non-trivial solutions were found later : hep-th/9908088 by K. Dasgupta, G. Rajesh, S. Sethi

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B-field	0	

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Transform the gauge field? Included in H, but...

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Transform the gauge field? Included in H, but... Extend  $T \oplus T^*$  with gauge bundle, apply an O(d+16, d+16)...

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Twist  $\Psi_{\pm}$ : only connection transformation, no *B*-transform

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For heterotic: no GCG construction  $\Rightarrow$  first try to do it !  $\hookrightarrow$  construct pure spinors, twist them; SUSY conditions.  $\mathcal{N}_{10D} = 1 : \epsilon \Rightarrow$  decomposition on 4D + 6D for  $\mathcal{N}_{4D} = 1$ :

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Only one internal spinor  $\eta_+$  (SU(3) structure)  $\Rightarrow$  Pure spinors:

$$\begin{array}{rcl} \Psi_+ &=& 8 \, e^{-\phi} \eta_+ \otimes \eta_+^\dagger = e^{-\phi} \, e^{-iJ} \ , \\ \\ \Psi_- &=& 8 \, e^{-\phi} \eta_+ \otimes \eta_-^\dagger = -i e^{-\phi} \Omega_3 \ . \end{array}$$

As before, except that no B-field: B plays a very different role in heterotic string...

 $T \oplus T^*$  and not E...

 $\hookrightarrow$  The generalized vielbein  ${\mathcal E}$  point of view: more difficult.

Twist  $\Psi_{\pm}$ : only connection transformation, no *B*-transform  $\hookrightarrow$  previous solutions mapped !

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 $\hookrightarrow$  constraints...

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