# Twist duality for flux backgrounds of type II and heterotic String Theory from Generalized Complex Geometry 

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arXiv:0903.0633 by D. A., R. Minasian, M. Petrini

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Geometric change: essentially provides a connection.
Transform metric and dilaton, and B-transform.
It relates backgrounds not related (simply) before.

Flux backgrounds for compactifications.

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math.DG/0209099 by N. Hitchin, math.DG/0401221 by M. Gualtieri
$\mathcal{M}_{\text {internal }}$ preserving at least $\mathcal{N}=1$ are Generalized CY (GCY)


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 GCG and type II TwistDual solutions
Heterotic Conclusions

- More generally: GCG is a natural set-up to study type II SUSY flux backgrounds.


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$\hookrightarrow$ Introduce some GCG objects, and then perform Twist duality to map two solutions.


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- Review appearance of GCG in type II SUSY flux backgrounds.


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- Heterotic backgrounds treatment.


## Generalized Complex Geometry and type II flux compactifications

The SUSY conditions for a SUGRA vacuum

- Type II SUGRA: $\mathcal{N}_{10 D}=2$ Spectrum: $g, \phi, H=d B, F_{p}, \psi_{\mu}^{1,2}, \lambda^{1,2}$


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## SUSY conditions $+\mathrm{BI} \Rightarrow$ e.o.m.

Main focus: solve the SUSY conditions

## Introduction

GCG and type II SUSY conditions $T \oplus T^{*}$ bundle

## Twist

Dual solutions
Heterotic
Conclusions

- SUSY conditions :

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\begin{aligned}
& 0=\delta \psi_{\mu}=D_{\mu} \epsilon+\frac{1}{4} H_{\mu} \mathcal{P} \epsilon+\frac{1}{16} e^{\phi} \sum_{n} H_{2 n} \gamma_{\mu} \mathcal{P}_{n} \epsilon \\
& 0=\delta \lambda=\left(\not \partial \phi+\frac{1}{2} H \mathcal{P}\right) \epsilon+\frac{1}{8} e^{\phi} \sum_{n}(-1)^{2 n}(5-2 n) \not \mathscr{F}_{2 n} \mathcal{P}_{n} \epsilon
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- SUSY conditions : CY condition

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Fluxes in the SUSY conditions $\Rightarrow$ GCG rewriting.

## Introduction

 GCG and type II SUSY conditions $T \oplus T^{*}$ bundle
## Twist

Dual solutions
Heterotic
Conclusions

- SUSY conditions :

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& 0=\delta \psi_{\mu}=D_{\mu} \epsilon+\frac{1}{4} H_{\mu} \mathcal{P} \epsilon+\frac{1}{16} e^{\phi} \sum_{n} H_{2 n} \gamma_{\mu} \mathcal{P}_{n} \epsilon \\
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Introduction GCG and type II SUSY conditions $T \oplus T^{*}$ bundle Twist

Dual solutions Heterotic Conclusions

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Dual solutions Heterotic Conclusions

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## Two globally defined non-vanishing spinors on $\mathcal{M}_{\text {internal }}: \eta^{1}, \eta^{2}$

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## SUSY conditions

 $T \oplus T^{*}$ bundleTwist
Dual solutions
Heterotic
Conclusions

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## Two globally defined non-vanishing spinors on $\mathcal{M}_{\text {internal }}: \eta^{1}, \eta^{2}$ Consider in the following $\eta^{1}=\eta^{2}(S U(3)$ structure).

- SUSY conditions on $\mathcal{M}_{\text {internal }}$ :

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## Introduction

GCG and type II SUSY conditions $T \oplus T^{*}$ bundle

## Twist

Dual solutions
Heterotic

David ANDRIOT

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GCG and type II SUSY conditions $T \oplus T^{*}$ bundle

## Twist

Dual solutions

## Heterotic

- SUSY conditions on $\mathcal{M}_{\text {internal }}$ : No flux
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$T \oplus T^{*}$ bundle

## Twist

Dual solutions

## Heterotic

Conclusions

No flux

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With flux (GCG)

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## Introduction

## SUSY conditions

$T \oplus T^{*}$ bundle

## Twist

Dual solutions

## Heterotic

Conclusions

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Introduction

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& \mathrm{d}\left(e^{3 A} \Psi_{1}\right)=0 \\
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GCY

- First define differential forms (with $\eta_{-}=\left(\eta_{+}\right)^{*}$ ) for $S U(3)$ :

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\Phi_{+} / N_{+}=e^{i \theta} & e^{-i J} \\
\Phi_{-} / N_{-}=-i & \Omega_{3}
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Introduction
GCG and type II SUSY conditions $T \oplus T^{*}$ bundle Twist

Dual solutions

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Introduction
GCG and type II SUSY conditions $T \oplus T^{*}$ bundle Twist

Dual solutions

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Introduction
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SUSY conditions can be expressed in terms of $\Psi_{ \pm}$! Spinors on $T \oplus T^{*}$, use of $\Psi_{ \pm} \Rightarrow$ GCG interpretations !

The generalized tangent bundle

GCG considers the fibration:
$T^{*} \hookrightarrow E$
$\downarrow \quad E$ : the generalized tangent bundle.
$\stackrel{\downarrow}{T}$

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Transition functions between two patches $U_{\alpha}$ and $U_{\beta}$ are

$$
\binom{v}{\xi}_{(\alpha)}=\left(\begin{array}{cc}
a & 0 \\
\omega a & a^{-T}
\end{array}\right)_{(\alpha \beta)}\binom{v}{\xi}_{(\beta)}=\binom{a v}{a^{-T} \xi-i_{a v} \omega} .
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$a \in G L(d, \mathbb{R}): a, a^{-T}$ usual patching of vectors and 1-forms.

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Additional shift of the 1 -form, given by the 2 -form $\omega_{(\alpha \beta)}=\mathrm{d} \Lambda_{(\alpha \beta)}$ :

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Additional shift of the 1 -form, given by the 2 -form $\omega_{(\alpha \beta)}=\mathrm{d} \Lambda_{(\alpha \beta)}$ : due to the non-trivial fibration of $T^{*}$ over $T$.

This non-trivial fibration: given by the 2 -form $B$. Geometrical interpretation of $B$ : "connective structure" of a gerbe.

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Additional shift of the 1 -form, given by the 2 -form $\omega_{(\alpha \beta)}=\mathrm{d} \Lambda_{(\alpha \beta)}$ : due to the non-trivial fibration of $T^{*}$ over $T$.

This non-trivial fibration: given by the 2 -form $B$.
Geometrical interpretation of $B$ : "connective structure" of a gerbe.
$\hookrightarrow \omega=\mathrm{d} \Lambda$ is "a gauge transformation" from $U_{\alpha}$ to $U_{\beta}$.

## Introduction

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Dual solutions

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Conclusions

- $E$ is equipped with a natural metric $\eta$ :

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In the $\left(\mathrm{d} x^{\mu}, \mathrm{d} y^{m}\right)$ basis: vielbeins and metric:

$$
\begin{aligned}
e=\left(\begin{array}{cc}
e_{\mathcal{B}} & 0 \\
e_{\mathcal{F}} A & e_{\mathcal{F}}
\end{array}\right) & g=e^{T} e=\left(\begin{array}{cc}
g_{\mathcal{B}}+A^{T} g_{\mathcal{F}} A & A^{T} g_{\mathcal{F}} \\
g_{\mathcal{F}} A & g_{\mathcal{F}}
\end{array}\right) \\
g_{\mathcal{B}}=e_{\mathcal{B}}^{T} e_{\mathcal{B}}, & g_{\mathcal{F}}=e_{\mathcal{F}}^{T} e_{\mathcal{F}}
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- On $E, B$ plays the role of a connection

David ANDRIOT

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Majorana-Weyl $\operatorname{Spin}(d, d)$ spinors on $E$ (locally $T \oplus T^{*}$ )

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Geometrical info. contained in $J, \Omega_{3},(B, \phi)(S U(3))$.

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A_{\mathcal{C}} & A_{\mathcal{F}} & 0 & 0 \\
\hline C_{\mathcal{B}} & C_{\mathcal{C}} & A_{\mathcal{B}}^{-T} & -A_{\mathcal{B}}^{-T} A_{\mathcal{C}}^{T} A_{\mathcal{F}}^{-T} \\
C_{\mathcal{C}^{\prime}} & C_{\mathcal{F}} & 0 & A_{\mathcal{F}}^{-T}
\end{array}\right)
\end{aligned}
$$

with separation of $\mathcal{B}$ and $\mathcal{F}$.

## The Twist transformation

On generalized vielbeins $\mathcal{E}$

- We consider $O \in G_{\text {geom }} \subset O(d, d)$ :

$$
\begin{aligned}
& A^{T} C+C^{T} A=0, \\
& O=\left(\begin{array}{cc}
A & 0 \\
C & A^{-T}
\end{array}\right)=\left(\begin{array}{cc|cc}
A_{\mathcal{B}} & 0 & 0 & 0 \\
A_{\mathcal{C}} & A_{\mathcal{F}} & 0 & 0 \\
\hline C_{\mathcal{B}} & C_{\mathcal{C}} & A_{\mathcal{B}}^{-T} & -A_{\mathcal{B}}^{-T} A_{\mathcal{C}}^{T} A_{\mathcal{F}}^{-T} \\
C_{\mathcal{C}^{\prime}} & C_{\mathcal{F}} & 0 & A_{\mathcal{F}}^{-T}
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with separation of $\mathcal{B}$ and $\mathcal{F}$.
The $O(d, d)$ constraint parameterizes $C$ as

$$
C=\left(\begin{array}{cc}
A_{\mathcal{B}}^{-T}\left(\tilde{C}_{\mathcal{B}}-A_{\mathcal{C}}^{T} A_{\mathcal{F}}^{-T} \tilde{C}_{\mathcal{C}}\right) & -A_{\mathcal{B}}^{-T}\left(\tilde{C}_{\mathcal{C}}^{T}+A_{\mathcal{C}}^{T} A_{\mathcal{F}}^{-T} \tilde{C}_{\mathcal{F}}\right) \\
A_{\mathcal{F}}^{-T} \tilde{\mathcal{C}}_{\mathcal{C}} & \tilde{C}_{\mathcal{F}}
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A_{\mathcal{F}}^{-T} \tilde{\mathcal{C}}_{\mathcal{C}}^{-T} \tilde{C}_{\mathcal{F}}
\end{array}\right),
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with $\tilde{C}_{B}$ and $\tilde{C}_{\mathcal{F}}$ anti-symmetric, $\tilde{C}_{C}$ unconstrained.

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\end{array}\right),
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with $\tilde{C}_{\mathcal{B}}$ and $\tilde{C}_{\mathcal{F}}$ anti-symmetric, $\tilde{C}_{\mathcal{C}}$ unconstrained.
We perform the transformation:

$$
\phi, \mathcal{E}=\left(\begin{array}{cc}
e & 0 \\
-e^{-T} B & e^{-T}
\end{array}\right) \mapsto \phi^{\prime}, \mathcal{E}^{\prime}=\mathcal{E} O .
$$

## The Twist transformation

On generalized vielbeins $\mathcal{E}$

- We consider $O \in G_{\text {geom }} \subset O(d, d)$ :

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A_{\mathcal{C}} & A_{\mathcal{F}} & 0 & 0 \\
\hline C_{\mathcal{B}} & C_{\mathcal{C}} & A_{\mathcal{B}}^{-T} & -A_{\mathcal{B}}^{-T} A_{C^{T}}^{T} A_{\mathcal{F}}^{-T} \\
C_{\mathcal{C}^{\prime}} & C_{\mathcal{F}} & 0 & A_{\mathcal{F}}
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A_{\mathcal{F}}^{-T} \tilde{C}_{\mathcal{C}} & A_{\mathcal{F}}^{-} \tilde{C}_{\mathcal{F}}
\end{array}\right),
$$

with $\tilde{C}_{\mathcal{B}}$ and $\tilde{C}_{\mathcal{F}}$ anti-symmetric, $\tilde{C}_{C}$ unconstrained.
We perform the transformation:

$$
\begin{gathered}
\phi, \mathcal{E}=\left(\begin{array}{cc}
e & 0 \\
-e^{-T} B & e^{-T}
\end{array}\right) \mapsto \varnothing^{\prime}, \mathcal{E}^{\prime}=\mathcal{E} O . \\
\leftrightarrow e^{\prime}=e A, B^{\prime}=A^{T} B A-A^{T} C, e^{\phi^{\prime}}=e^{\phi}|\operatorname{det}(A)|^{\frac{1}{2}} .
\end{gathered}
$$

## Concrete example:

 NDRIOT
## Introduction

## GCG and type II

## Twist

On generalized vielbeins
On pure spinors
Dual solutions

## Heterotic

## Conclusions

Concrete example: $A=0$ :

David ANDRIOT

$$
e=\left(\begin{array}{cc}
e_{\mathcal{B}} & 0 \\
0 & e_{\mathcal{F}}
\end{array}\right), \mathrm{d} s^{2}=g_{\mu \nu} \mathrm{d} x^{\mu} \mathrm{d} x^{\nu}+g_{m n}\left(\mathrm{~d} y^{m}+0\right)\left(\mathrm{d} y^{n}+0\right)
$$

On generalized vielbeins
On pure spinors
Dual solutions

## Heterotic

## Conclusions

Concrete example: $A=0$ :

David ANDRIOT

## Introduction

GCG and type II
Twist
On generalized vielbeins
On pure spinors
Dual solutions

## Heterotic

Conclusions

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## Introduction

 GCG and type II Twist
## On generalized

 vielbeinsOn pure spinors
Dual solutions
Heterotic
Conclusions

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\begin{aligned}
& e=\left(\begin{array}{cc}
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\end{aligned}
$$

The twist transformation gives:

$$
\begin{aligned}
& e^{\prime}=\left(\begin{array}{cc}
e_{\mathcal{B}} A_{\mathcal{B}} & 0 \\
\left(e_{\mathcal{F}} A_{\mathcal{F}}\right)\left(A_{\mathcal{F}}^{-1} A_{\mathcal{C}}\right) & e_{\mathcal{F}} A_{\mathcal{F}}
\end{array}\right) \\
& \mathrm{d} s^{2}=g_{\mu \nu}^{\prime} \mathrm{d} x^{\mu} \mathrm{d} x^{\nu}+g_{m n}^{\prime}\left(\mathrm{d} y^{m}+A_{\rho}^{\prime m} \mathrm{~d} x^{\rho}\right)\left(\mathrm{d} y^{n}+A_{\sigma}^{\prime n} \mathrm{~d} x^{\sigma}\right), \\
& A^{\prime}=A_{\mathcal{F}}^{-1} A_{\mathcal{C}}, g_{\mu \nu}^{\prime}=\left(A_{\mathcal{B}}^{T} g_{\mathcal{B}} A_{\mathcal{B}}\right)_{\mu \nu}, g_{m n}^{\prime}=\left(A_{\mathcal{F}}^{T} g_{\mathcal{F}} A_{\mathcal{F}}\right)_{m n} .
\end{aligned}
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\begin{aligned}
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\end{aligned}
$$

$\hookrightarrow A_{\mathcal{C}}$ generates a connection $A^{\prime}$

Concrete example: $A=0$ :

$$
\begin{aligned}
& e=\left(\begin{array}{cc}
e_{\mathcal{B}} & 0 \\
0 & e_{\mathcal{F}}
\end{array}\right), \mathrm{d} s^{2}=g_{\mu \nu} \mathrm{d} x^{\mu} \mathrm{d} x^{\nu}+g_{m n}\left(\mathrm{~d} y^{m}+0\right)\left(\mathrm{d} y^{n}+0\right), \\
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e_{\mathcal{B}} A_{\mathcal{B}} & 0 \\
\left(e_{\mathcal{F}} A_{\mathcal{F}}\right)\left(A_{\mathcal{F}}^{-1} A_{\mathcal{C}}\right) & e_{\mathcal{F}} A_{\mathcal{F}}
\end{array}\right) \\
& \mathrm{d} s^{2}=g_{\mu \nu}^{\prime} \mathrm{d} x^{\mu} \mathrm{d} x^{\nu}+g_{m n}^{\prime}\left(\mathrm{d} y^{m}+A_{\rho}^{\prime m} \mathrm{~d} x^{\rho}\right)\left(\mathrm{d} y^{n}+A_{\sigma}^{\prime n} \mathrm{~d} x^{\sigma}\right), \\
& A^{\prime}=A_{\mathcal{F}}^{-1} A_{\mathcal{C}}, g_{\mu \nu}^{\prime}=\left(A_{\mathcal{B}}^{T} g_{\mathcal{B}} A_{\mathcal{B}}\right)_{\mu \nu}, g_{m n}^{\prime}=\left(A_{\mathcal{F}}^{T} g_{\mathcal{F}} A_{\mathcal{F}}\right)_{m n} .
\end{aligned}
$$

$\hookrightarrow A_{\mathcal{C}}$ generates a connection $A^{\prime}\left(A_{\mathcal{C}}\right.$ has to be coordinate dependent... Important difference with T-duality),

Concrete example: $A=0$ :

$$
\begin{aligned}
& e=\left(\begin{array}{cc}
e_{\mathcal{B}} & 0 \\
0 & e_{\mathcal{F}}
\end{array}\right), \mathrm{d} s^{2}=g_{\mu \nu} \mathrm{d} x^{\mu} \mathrm{d} x^{\nu}+g_{m n}\left(\mathrm{~d} y^{m}+0\right)\left(\mathrm{d} y^{n}+0\right), \\
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\end{array}\right) \\
& \mathrm{d} s^{2}=g_{\mu \nu}^{\prime} \mathrm{d} x^{\mu} \mathrm{d} x^{\nu}+g_{m n}^{\prime}\left(\mathrm{d} y^{m}+A_{\rho}^{\prime m} \mathrm{~d} x^{\rho}\right)\left(\mathrm{d} y^{n}+A_{\sigma}^{\prime n} \mathrm{~d} x^{\sigma}\right), \\
& A^{\prime}=A_{\mathcal{F}}^{-1} A_{\mathcal{C}}, g_{\mu \nu}^{\prime}=\left(A_{\mathcal{B}}^{T} g_{\mathcal{B}} A_{\mathcal{B}}\right)_{\mu \nu}, g_{m n}^{\prime}=\left(A_{\mathcal{F}}^{T} g_{\mathcal{F}} A_{\mathcal{F}}\right)_{m n} .
\end{aligned}
$$

$\hookrightarrow A_{\mathcal{C}}$ generates a connection $A^{\prime}\left(A_{\mathcal{C}}\right.$ has to be coordinate dependent... Important difference with T-duality), $\hookrightarrow A_{\mathcal{B}}, A_{\mathcal{F}}$ transform the metric .

Concrete example: $A=0$ :

$$
\begin{aligned}
& e=\left(\begin{array}{cc}
e_{\mathcal{B}} & 0 \\
0 & e_{\mathcal{F}}
\end{array}\right), \mathrm{d} s^{2}=g_{\mu \nu} \mathrm{d} x^{\mu} \mathrm{d} x^{\nu}+g_{m n}\left(\mathrm{~d} y^{m}+0\right)\left(\mathrm{d} y^{n}+0\right), \\
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$\hookrightarrow$ The dilaton transforms accordingly:

$$
e^{\phi^{\prime}}=e^{\phi}\left|\operatorname{det}\left(A_{\mathcal{B}}\right) \operatorname{det}\left(A_{\mathcal{F}}\right)\right|^{\frac{1}{2}} .
$$

Concrete example: $A=0$ :

## Twist

## On generalized

 vielbeinsOn pure spinors
Dual solutions
Heterotic
Conclusions

$$
\begin{aligned}
& e=\left(\begin{array}{cc}
e_{\mathcal{B}} & 0 \\
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\end{array}\right), \mathrm{d} s^{2}=g_{\mu \nu} \mathrm{d} x^{\mu} \mathrm{d} x^{\nu}+g_{m n}\left(\mathrm{~d} y^{m}+0\right)\left(\mathrm{d} y^{n}+0\right), \\
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$$
B^{\prime}=-\left(\begin{array}{cc}
\tilde{C}_{\mathcal{B}} & -\tilde{\tilde{C}}_{\mathcal{C}}^{T} \\
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Concrete example: $A=0$ :

## Twist

## On generalized

 vielbeinsOn pure spinors
Dual solutions
Heterotic
Conclusions

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\begin{aligned}
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& A^{\prime}=A_{\mathcal{F}}^{-1} A_{\mathcal{C}}, g_{\mu \nu}^{\prime}=\left(A_{\mathcal{B}}^{T} g_{\mathcal{B}} A_{\mathcal{B}}\right)_{\mu \nu}, g_{m n}^{\prime}=\left(A_{\mathcal{F}}^{T} g_{\mathcal{F}} A_{\mathcal{F}}\right)_{m n}
\end{aligned}
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\end{array}\right)
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## On pure spinors $\Psi_{ \pm}$

David ANDRIOT

## Introduction

## GCG and type II

## Twist

On generalized vielbeins

On pure spinors
Dual solutions

## Heterotic

## Conclusions

## On pure spinors $\Psi_{ \pm}$

$$
\left\{\Gamma^{m}, \Gamma^{n}\right\}=\left\{\Gamma_{m}, \Gamma_{n}\right\}=0, \quad\left\{\Gamma^{m}, \Gamma_{n}\right\}=\delta_{n}^{m} \quad m, n=1 \ldots d
$$

## On pure spinors $\Psi_{ \pm}$

Metric $\eta \Rightarrow \operatorname{Clifford}$ algebra $\operatorname{Cliff}(d, d)$ on $E$ :

$$
\left\{\Gamma^{m}, \Gamma^{n}\right\}=\left\{\Gamma_{m}, \Gamma_{n}\right\}=0, \quad\left\{\Gamma^{m}, \Gamma_{n}\right\}=\delta_{n}^{m} \quad m, n=1 \ldots d .
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Action on $\Psi_{ \pm}$: wedges, contractions: $\Gamma^{n}=\mathrm{d} x^{n} \wedge, \Gamma_{m}=\iota_{\partial_{m}}$.

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Action on $\Psi_{ \pm}$: wedges, contractions: $\Gamma^{n}=\mathrm{d} x^{n} \wedge, \Gamma_{m}=\iota_{\partial_{m}}$. $\hookrightarrow O(d, d)$ in the spinorial representation:

$$
O=e^{-\frac{1}{4} \Theta_{M N} \sigma^{M N}} \quad, \quad M, N=1 \ldots d+d .
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\begin{aligned}
& O=e^{-\frac{1}{4} \Theta_{M N} \sigma^{M N}} \quad, \quad M, N=1 \ldots d+d . \\
& \sigma^{M N}=\left[\Gamma^{M}, \Gamma^{N}\right],
\end{aligned}
$$

## On pure spinors $\Psi_{ \pm}$

On pure spinors

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\sigma^{M N}=\left[\Gamma^{M}, \Gamma^{N}\right], & \Theta_{M N}=\left(\begin{array}{cc}
a^{m}{ }_{n} & \beta^{m n} \\
B_{m n} & -a_{m}{ }^{n}
\end{array}\right) .
\end{array}
$$

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a^{m}{ }_{n} & \beta^{m n} \\
B_{m n} & -a_{m}{ }^{n}
\end{array}\right) .
\end{array}
$$

$G L(d)$ transformation (previous $A$ ):
$O_{a}=e^{-\frac{1}{4}\left(a^{m}{ }_{n}\left[\Gamma_{m}, \Gamma^{n}\right]-a_{m}{ }^{n}\left[\Gamma^{m}, \Gamma_{n}\right]\right)}=\cdots=\frac{1}{\sqrt{\operatorname{det} A}} e^{a^{m}{ }_{n} \mathrm{~d} x^{n} \wedge \iota \partial_{m}}$.

## On pure spinors $\Psi_{ \pm}$

Metric $\eta \Rightarrow \operatorname{Clifford}$ algebra $\operatorname{Cliff}(d, d)$ on $E$ :

$$
\left\{\Gamma^{m}, \Gamma^{n}\right\}=\left\{\Gamma_{m}, \Gamma_{n}\right\}=0, \quad\left\{\Gamma^{m}, \Gamma_{n}\right\}=\delta_{n}^{m} \quad m, n=1 \ldots d .
$$

Action on $\Psi_{ \pm}$: wedges, contractions: $\Gamma^{n}=\mathrm{d} x^{n} \wedge, \Gamma_{m}=\iota_{\partial_{m}}$.
$\hookrightarrow O(d, d)$ in the spinorial representation:

$$
\begin{array}{cl}
O=e^{-\frac{1}{4} \Theta_{M N} \sigma^{M N}} & , \quad M, N=1 \ldots d+d . \\
\sigma^{M N}=\left[\Gamma^{M}, \Gamma^{N}\right], & \Theta_{M N}=\left(\begin{array}{cc}
a^{m}{ }_{n} & \beta^{m n} \\
B_{m n} & -a_{m}{ }^{n}
\end{array}\right) .
\end{array}
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$B$-transform (previous $C$ ):

$$
O_{B}=e^{-\frac{1}{2} B_{m n} \Gamma^{m n}}=e^{-\frac{1}{2} B_{m n} \mathrm{~d} x^{m} \wedge \mathrm{~d} x^{n}}
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## Previous twist transformation:

## David ANDRIOT <br> Introduction

$$
\mathcal{E} \mapsto \mathcal{E}^{\prime}=\mathcal{E} O, O=\left(\begin{array}{cc}
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On generalized vielbeins

On pure spinors
Dual solutions

## Heterotic

## Conclusions

## Previous twist transformation:

David ANDRIOT
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## Previous twist transformation:

## Introduction

GCG and type II

## Twist

On generalized vielbeins
On pure spinors
Dual solutions

## Heterotic

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Previous twist transformation:

Introduction GCG and type II Twist
On generalized vielbeins
On pure spinors
Dual solutions

## Heterotic

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Introduction GCG and type II Twist
On generalized vielbeins
On pure spinors
Dual solutions
Heterotic
Conclusions
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Introduction GCG and type II Twist

On pure spinors
Dual solutions
Heterotic
Conclusions
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Previous twist transformation:

Introduction GCG and type II Twist

On pure spinors
Dual solutions
Heterotic
Conclusions
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Previous twist transformation:

Introduction
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Previous twist transformation:

## Twist

## On generalized

 vielbeinsOn pure spinors
Dual solutions
Heterotic
Conclusions
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Previous twist transformation:

## Twist

## On generalized

 vielbeinsOn pure spinors
Dual solutions
Heterotic
Conclusions
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Previous twist transformation:

## Twist

## On generalized

 vielbeinsOn pure spinors
Dual solutions
Heterotic
Conclusions
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David ANDRIOT

## Introduction

GCG and type II Twist

On generalized vielbeins
On pure spinors
Dual solutions

## Heterotic

Conclusions

Remarks (work in progress):

David ANDRIOT

## Introduction

On pure spinors
Dual solutions

## Heterotic

Conclusions

Remarks (work in progress):

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## Introduction

On pure spinors

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## Introduction

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Definition of the transformed RR as new solutions of SUSY.


## Mapping type IIB solutions

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## Constraints to generate solutions

Consider a solution of the SUSY conditions.

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\begin{aligned}
& \mathrm{d}\left(e^{3 A} \Psi_{1}\right)=0 \\
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Develop in terms of $\Psi_{1,2}, R, O_{f}, \theta_{c}^{ \pm}$

## Mapping type IIB solutions

## Constraints to generate solutions

Consider a solution of the SUSY conditions.
Perform a twist: $\Psi_{ \pm}^{\prime}=e^{i \theta_{c}^{ \pm}} O_{f} \Psi_{ \pm} \Rightarrow$ Get a new solution?

$$
\begin{array}{ll}
\mathrm{d}\left(e^{3 A} \Psi_{1}\right)=0 & \mathrm{~d}\left(e^{3 A} \Psi_{1}^{\prime}\right)=0 \\
\mathrm{~d}\left(e^{2 A} \operatorname{Re} \Psi_{2}\right)=0 \\
\mathrm{~d}\left(e^{4 A} \operatorname{Im} \Psi_{2}\right)=e^{4 A} e^{-B} * \lambda(F)=R . & \Rightarrow \\
\mathrm{d}\left(e^{2 A} \operatorname{Re} \Psi_{2}^{\prime}\right)=0 \\
\mathrm{~d}\left(e^{4 A} \operatorname{Im} \Psi_{2}^{\prime}\right)=R^{\prime}
\end{array}
$$

Develop in terms of $\Psi_{1,2}, R, O_{f}, \theta_{c}^{ \pm} \Rightarrow$ constraints on the twist:
$\mathrm{d}\left(O_{f}\right) \Psi_{1}=0$
$c_{\theta_{c}^{+}} \mathrm{d}\left(O_{f}\right) e^{2 A} \operatorname{Re} \Psi_{2}-s_{\theta_{c}^{+}} \mathrm{d}\left(e^{-2 A} O_{f}\right) e^{4 A} \operatorname{Im} \Psi_{2}=e^{-2 A} s_{\theta_{c}^{+}} O_{f} R$
$s_{\theta_{c}^{+}} \mathrm{d}\left(e^{2 A} O_{f}\right) e^{2 A} \operatorname{Re} \Psi_{2}+c_{\theta_{c}^{+}} \mathrm{d}\left(O_{f}\right) e^{4 A} \operatorname{Im} \Psi_{2}=R^{\prime}-c_{\theta_{c}^{+}} O_{f} R$.

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## Mapping type IIB solutions

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## Mapping type IIB solutions

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$$

$\theta_{c}^{+} \neq 0$, coordinate dependent $O_{f}$

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$$

$\theta_{c}^{+} \neq 0$, coordinate dependent $O_{f} \Rightarrow$ mix NSNS and RR sectors.

## Examples of type IIB dual solutions

## Introduction

Solutions:

| $\mathcal{M}_{\text {internal }}$ | $T^{6}=T^{2} \times T^{4}$ |  |
| :--- | :--- | :--- |
|  |  |  |
|  |  |  |
|  |  |  |

## Examples of type IIB dual solutions

## Introduction

Solutions:

|  | $T^{6}=T^{2} \times T^{4}$ |  |
| :---: | :---: | :---: |
| $d s_{6}^{2}$ | $e^{-2 A} \mathrm{~d} x_{\mathcal{B}}^{2}+e^{-2 A} \mathrm{~d} x_{\mathcal{F}}^{2}$ |  |
|  |  |  |
|  |  |  |
|  |  |  |
|  |  |  |

Examples of type IIB dual solutions

## Introduction

## Constraints

Examples

## Heterotic

 Conclusions
## Solutions:

| $\mathcal{M}_{\text {internal }}$ | $T^{6}=T^{2} \times T^{4}$ |  |
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| Sources, $\theta$ | $O 3, \theta=\frac{\pi}{2}$ |  |
|  |  |  |
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|  |  |  |

Examples of type IIB dual solutions

## Introduction

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| RR | $g_{s} F_{5}=e^{4 A} * \mathrm{~d}\left(e^{-4 A}\right),\left(F_{3}\right)$ |  |  |  |
|  |  |  |  |  |
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Examples of type IIB dual solutions

## Introduction

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|  |  |  |  |  |

Examples of type IIB dual solutions

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| $e^{\phi}$ | $g_{s}$ |  |

Examples of type IIB dual solutions

## Solutions:

| $\mathcal{M}$ internal | $T^{6}=T^{2} \times T^{4}$ | $T^{2}$ | $\hookrightarrow$ |
| :---: | :---: | :---: | :---: |
|  |  | $\mathcal{M}$ |  |
|  |  | $\downarrow$ |  |
| $d s_{6}^{2}$ | $e^{-2 A} \mathrm{~d} x_{\mathcal{B}}^{2}+e^{-2 A} \mathrm{~d} x_{\mathcal{F}}^{2}$ |  |  |
| Sources, $\theta$ | $O 3, \theta=\frac{\pi}{2}$ |  |  |
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Examples of type IIB dual solutions

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| :---: | :---: | :---: | :---: |
|  |  |  |  |
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| Sources, $\theta$ | $O 3, \theta=\frac{\pi}{2}$ |  |  |
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Examples of type IIB dual solutions

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| Sources, $\theta$ | $O 3, \theta=\frac{\pi}{2}$ | $O 5 / / \mathcal{F}, \theta=0$ |
| RR | $g_{s} F_{5}=e^{4 A} * \mathrm{~d}\left(e^{-4 A}\right),\left(F_{3}\right)$ | $g_{s} F_{3}=-e^{-4 A} * \mathrm{~d}\left(e^{2 A} J\right)$ |
| NSNS | $\left(H=g_{s} * F_{3}\right)$ | 0 |
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Examples of type IIB dual solutions

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Examples of type IIB dual solutions

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| $e^{\phi}$ | $g_{s}$ | $g_{s} e^{2 A}$ |

Twist duality map?

Examples of type IIB dual solutions

## Solutions:

| internal | $T^{6}=T^{2} \times T^{4}$ | $T^{2} \hookrightarrow$ |
| :---: | :---: | :---: |
|  |  |  |
| $d s_{6}^{2}$ | $e^{-2 A} \mathrm{~d} x_{\mathcal{B}}^{2}+e^{-2 A} \mathrm{~d} x_{\mathcal{F}}^{2}$ | $e^{-2 A} \mathrm{~d} x_{\mathcal{B}}^{2}+e^{2 A}\left(\mathrm{~d} x_{\mathcal{F}}+A\right)^{2}$ |
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| $e^{\phi}$ | $g_{s}$ | $g_{s} e^{2 A}$ |

Twist duality map?

$$
A_{\mathcal{B}}=1_{4}
$$

Examples of type IIB dual solutions

## Solutions:

| $\mathcal{M}$ internal | $T^{6}=T^{2} \times T^{4}$ | $T^{2} \hookrightarrow$ |
| :---: | :---: | :---: |
|  |  |  |
| $d s_{6}^{2}$ | $e^{-2 A} \mathrm{~d} x_{\mathcal{B}}^{2}+e^{-2 A} \mathrm{~d} x_{\mathcal{F}}^{2}$ | $e^{-2 A} \mathrm{~d} x_{\mathcal{B}}^{2}+e^{2 A}\left(\mathrm{~d} x_{\mathcal{F}}+A\right)^{2}$ |
| $T^{4}$ |  |  |
| Sources, $\theta$ | $O 3, \theta=\frac{\pi}{2}$ | $O 5 / / \mathcal{F}, \theta=0$ |
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Twist duality map?

$$
A_{\mathcal{B}}=1_{4}, A_{\mathcal{F}}=1_{2} \times e^{2 A}
$$

Examples of type IIB dual solutions
Solutions:

| internal | $T^{6}=T^{2} \times T^{4}$ | $T^{2} \hookrightarrow$ |
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$$
A_{\mathcal{B}}=1_{4}, A_{\mathcal{F}}=1_{2} \times e^{2 A}, A_{C}{ }^{I}{ }_{\mu}=e^{2 A} A_{\mu}^{I}
$$

Examples of type IIB dual solutions
Solutions:

| internal | $T^{6}=T^{2} \times T^{4}$ | $T^{2} \hookrightarrow$ |
| :---: | :---: | :---: |
|  |  |  |
|  | $e^{-2 A} \mathrm{~d} x_{\mathcal{B}}^{2}+e^{-2 A} \mathrm{~d} x_{\mathcal{F}}^{2}$ | $e^{-2 A} \mathrm{~d} x_{\mathcal{B}}^{2}+e^{2 A}\left(\mathrm{~d} x_{\mathcal{F}}+A\right)^{2}$ |
| $T^{4}$ |  |  |
| Sources, $\theta$ | $O 3, \theta=\frac{\pi}{2}$ | $O 5 / / \mathcal{F}, \theta=0$ |
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Twist duality map?

$$
A_{\mathcal{B}}=1_{4}, A_{\mathcal{F}}=1_{2} \times e^{2 A}, A_{\mathcal{C}}{ }^{I}{ }_{\mu}=e^{2 A} A_{\mu}^{I}, \theta_{c}^{+}=-\frac{\pi}{2} .
$$

Examples of type IIB dual solutions
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| internal | $T^{6}=T^{2} \times T^{4}$ | $T^{2} \hookrightarrow$ |
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|  |  |  |
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| $T^{4}$ |  |  |
| Sources, $\theta$ | $O 3, \theta=\frac{\pi}{2}$ | $O 5 / / \mathcal{F}, \theta=0$ |
| RR | $g_{s} F_{5}=e^{4 A} * \mathrm{~d}\left(e^{-4 A}\right),\left(F_{3}\right)$ | $g_{s} F_{3}=-e^{-4 A} * \mathrm{~d}\left(e^{2 A} J\right)$ |
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Twist duality map?

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$B$-transform if needed...

Examples of type IIB dual solutions
Solutions:

| internal | $T^{6}=T^{2} \times T^{4}$ | $T^{2} \hookrightarrow$ |
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|  |  |  |
| $d s_{6}^{2}$ | $e^{-2 A} \mathrm{~d} x_{\mathcal{B}}^{2}+e^{-2 A} \mathrm{~d} x_{\mathcal{F}}^{2}$ | $e^{-2 A} \mathrm{~d} x_{\mathcal{B}}^{2}+e^{2 A}\left(\mathrm{~d} x_{\mathcal{F}}+A\right)^{2}$ |
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Explicit non-trivial fibration solutions? Among nilmanifolds: twisted tori, GCY...

Among 34 nilmanifolds, only 5 non-trivial $T^{2}$ bundles.

David ANDRIOT

## Introduction

GCG and type II
Twist
Dual solutions
Constraints
Examples
Heterotic
Conclusions


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David ANDRIOT

## Introduction

## Twist

Dual solutions

## Heterotic

Solutions
Pure spinors, SUSY Conclusions

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$\hookrightarrow$ decompose on $4 D+6 D$, and work out conditions for $\Psi_{ \pm}$ Same as "type A" solutions of type IIB with $F=0, A=0$ $\hookrightarrow$ worked out in

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Rewrite the result as a wedge product:
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