

Can pulse coherence reveal the low energy charge modes in condensed matter?

Peter Abbamonte

University of Illinois at Urbana-Champaign

Acknowledgements: Discussions with William Graves, Bates Lab/MIT; Michael Rübhausen, University of Hamburg; David Moncton, MIT; Sunny Sinha, UCSD; Phil Platzman, Bell Labs

Why do we care about charge modes?

21st century grand challenge (apologies for DOE speak):

Understand the physics of interacting electron systems, in particular the doped-Mott insulator

We “understand” an interacting electron system when ...

1. We can deduce its ground state (preferably a Slater determinant)
2. We can describe its elementary excitations

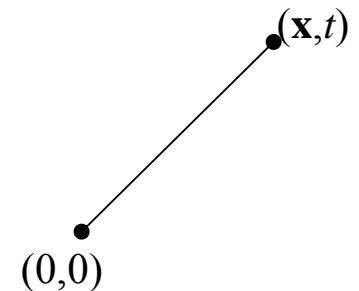
“Charge modes” \equiv those excitations that modulate the electron density

Characterized by dynamic structure factor,

$$S(\mathbf{k}, \omega) = \sum_n |\langle n | \rho(\mathbf{k}) | 0 \rangle|^2 \delta(\omega - \omega_n) = -1/\pi \text{Im}[\chi(\mathbf{k}, \omega)]$$

Measured directly by inelastic x-ray or electron scattering

$$\chi(\mathbf{x}, t) = -i \langle 0 | \hat{n}(\mathbf{x}, t) \hat{n}(0, 0) | 0 \rangle \theta(t)$$

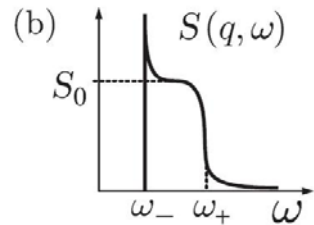


Why do we care about charge modes?

A Classic Example:

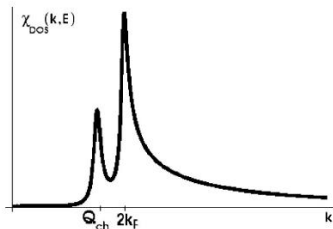
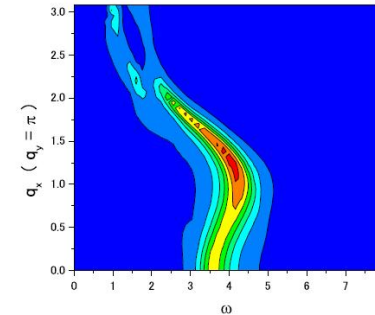
Interacting electron gas in RPA (Pines & Nozieres, 1952) $S(\mathbf{k}, \omega)$

Some Recent Examples:



Luttinger liquid with nonlinear dispersion
Pustilnik, *PRL*, **96**, 196405 (2006)

Amplitude modes in a d-wave superconductor,
Lee & Nagaosa, *PRB*, **68**, 24516 (2003)



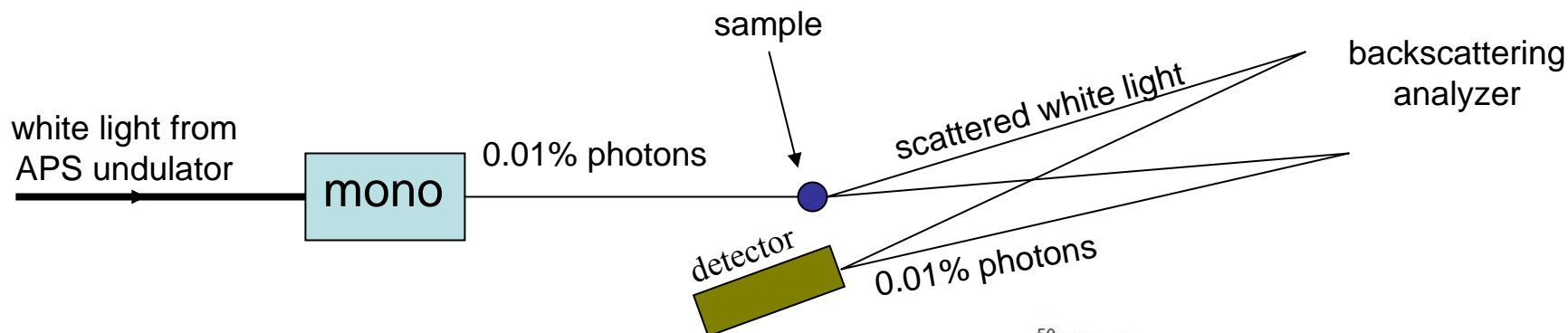
Stripe theory of HTSC

Kivelson, Bindloss, Fradkin, et. al., *RMP*, **75**,
1201 (2004)

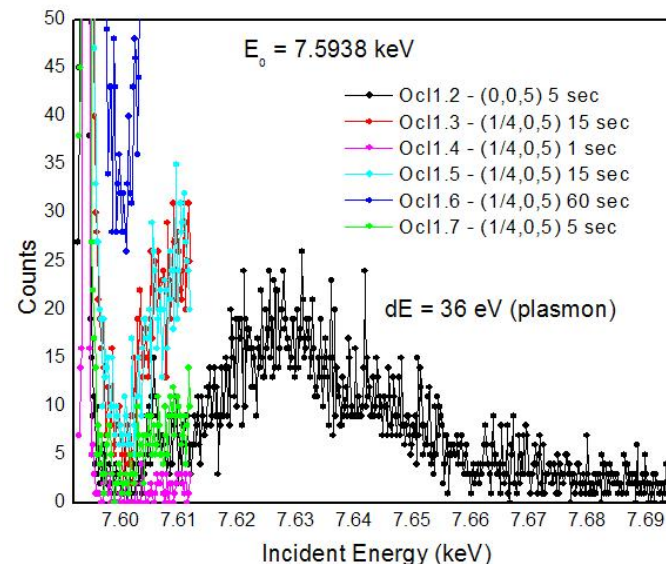
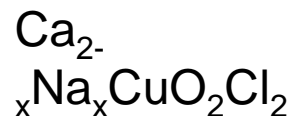
Inelastic x-ray scattering – why doesn't it work?!

IXS measures $S(\mathbf{k}, \omega)$, i.e. the charge response. Why doesn't it work?

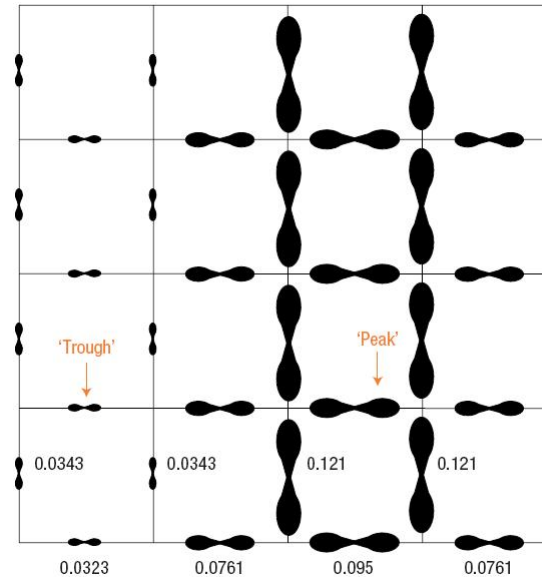
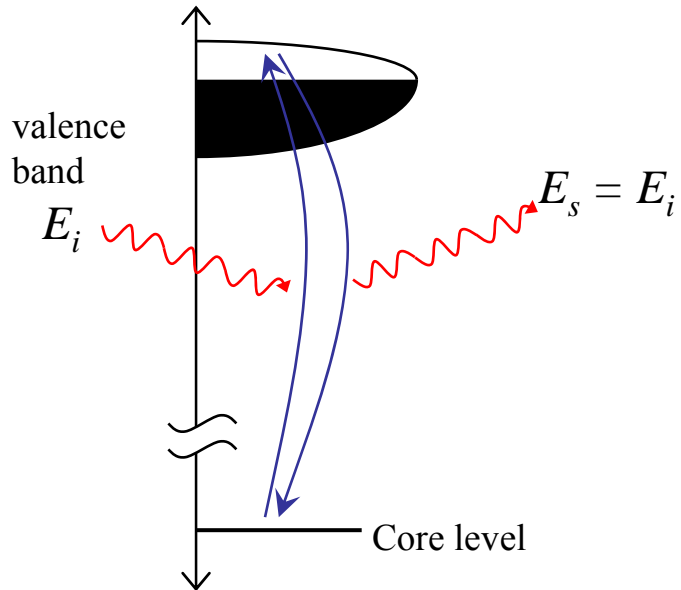
1. XRD is coherent, scales like N^2
2. IXS incoherent, scales like $N_{\text{valence}} \sim 0.1 e \ll N$
3. Mono and analyzer both throw away 99.9999% of the photons



Upshot: Signal is weak and sitting on tail of elastic, which is stronger by N^2/N_{valence}



Resonance techniques - diffraction



Charge amplitude
of static stripes
Abbamonte,
Rusydi, et. al.,
Nature Physics, **1**,
155 (2005)

$\Delta v = 0.063$ holes
 $A = 0.59$ holes

Wigner crystal melting, Rusydi, et. al, *PRL* (2006)

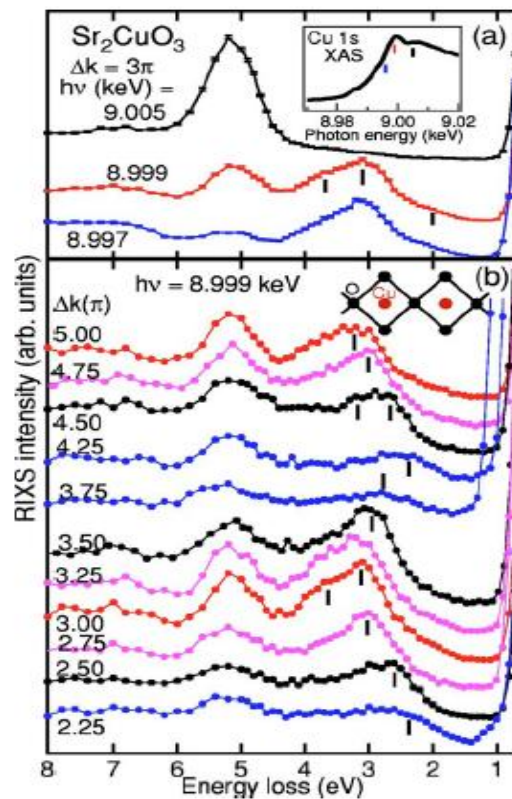
Wigner crystallization in ladders, Abbamonte, et. al., *Nature*, **431**, 1078 (2004)

Charge disproportionation in PCMO, Thomas, et. al. *PRL*, **92**, 237204 (2004)

Depletion zone at an interface, Abbamonte, et. al., *Science*, **297**, 581 (2002)

Are these correlations present dynamically in SC state?

Resonance techniques – inelastic scattering (RIXS)

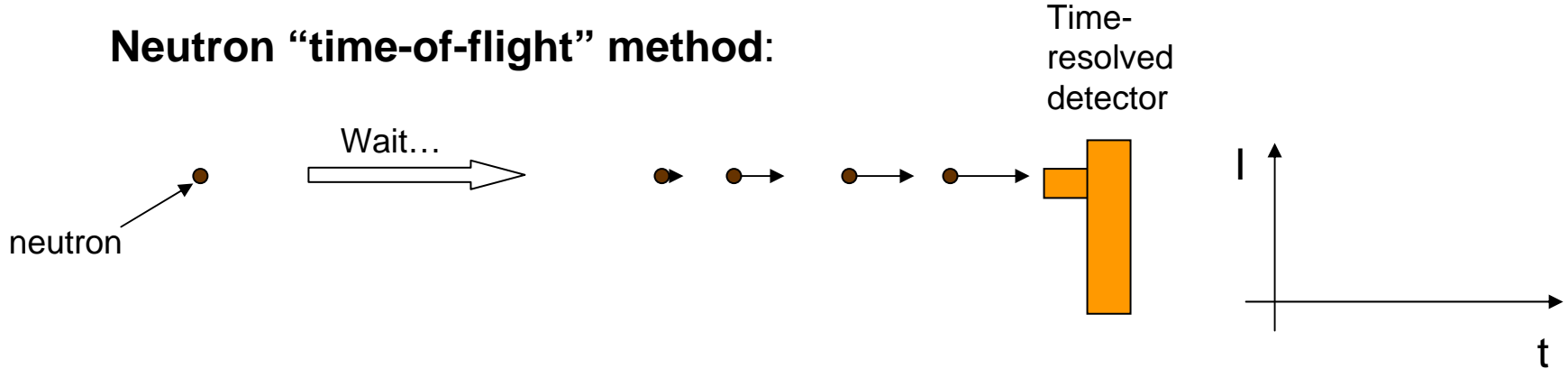


S. Suga, et. al., *PRB*, **72**, 81101(R) (2005)

We're doing chemistry. Useful, but not good enough.

Work in time domain?

Neutron “time-of-flight” method:



- Uses all neutrons in the bunch (no monochromator)
- No a-priori limit on time resolution
- Does not work for photons because they all travel at c . (except maybe near an edge)

Q: How can it be that the frequency structure changes but not the time structure?

$$|\mathbf{k}\rangle \rightarrow \exp(i\omega_n t) |\mathbf{k}\rangle$$

No change in time trace, $E^2(t)$.

Phase space reformulation

System in its ground state: $|0\rangle$

System in its ground state, plus a photon: $a^\dagger_{\mathbf{k}} |0\rangle$

Single photon wave packet: $|\psi\rangle = \sum_{\mathbf{k}} A(\mathbf{k}) a^\dagger_{\mathbf{k}} |0\rangle$

Satisfies $\Delta E \Delta t = \hbar/2$

Many photons: $|\psi\rangle = [\sum_{\mathbf{k}_1} A(\mathbf{k}_1) a^\dagger_{\mathbf{k}_1}] [\sum_{\mathbf{k}_2} A(\mathbf{k}_2) a^\dagger_{\mathbf{k}_2}] \dots |0\rangle$

Still satisfies $\Delta E \Delta t = \hbar/2$

Inelastic x-ray scattering:

System undergoes transition: $|0\rangle \rightarrow |n\rangle$

Probability of doing this: $|\langle n | \rho(\mathbf{k}) |0\rangle|^2$

Photon undergoes frequency shift: $|\psi\rangle \rightarrow \exp(i\omega_n t) |\psi\rangle$

No change in time trace, $E^2(t)$

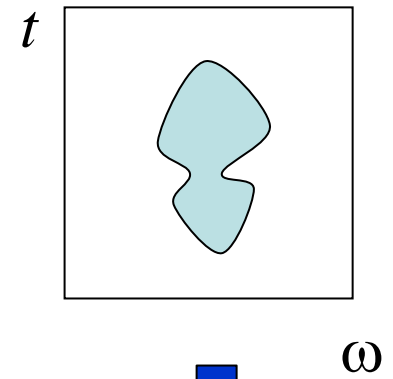
BUT ... incoherent sum \Rightarrow increase in phase space area

Scattered pulse is only partially coherent.
Contains noise.

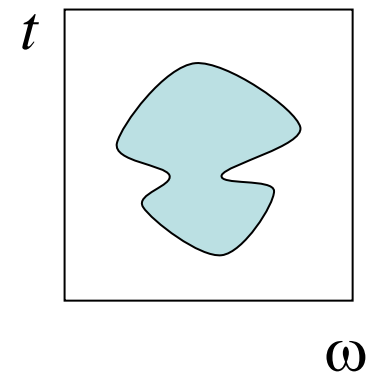
Phase space:

$$E(\omega) \cdot E(t)$$

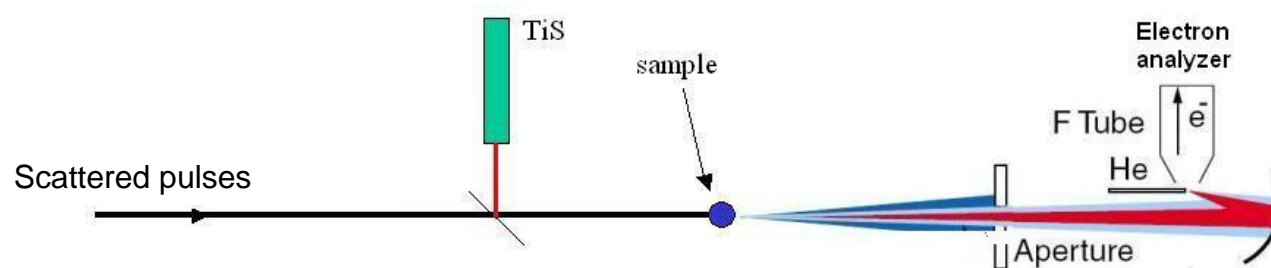
$$\Delta E \Delta t = \hbar/2$$



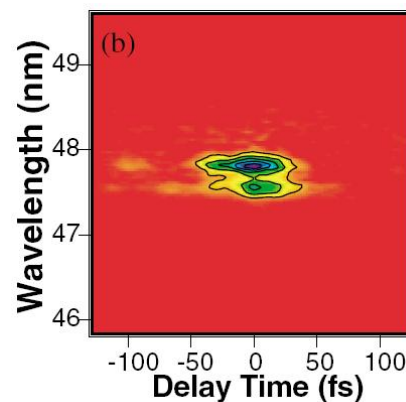
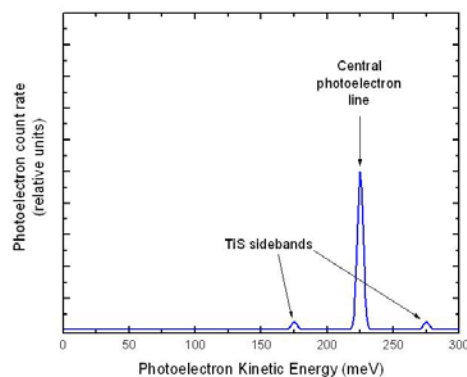
$$\Delta E \Delta t = \hbar/2$$



Possible route: X-FROG



Adapted from
Sekikawa, *PRL*, **91**,
103902 (2003)



X-FROG at
25 eV

Tune time delay. Complete phase space reconstruction via simple matrix inversion, provided TiS pulse shape is known

- Entire scattered pulse used at once. More efficient than monochromator.
- Resolution determined by delay line. No intrinsic limit. $\Delta E = 2\pi/\text{rep rate} \sim 3.3 \text{ peV} @ 10\text{Hz}$
- Needs transform-limited pulses
- X-ray analogue to neutron time-of-flight or spin echo

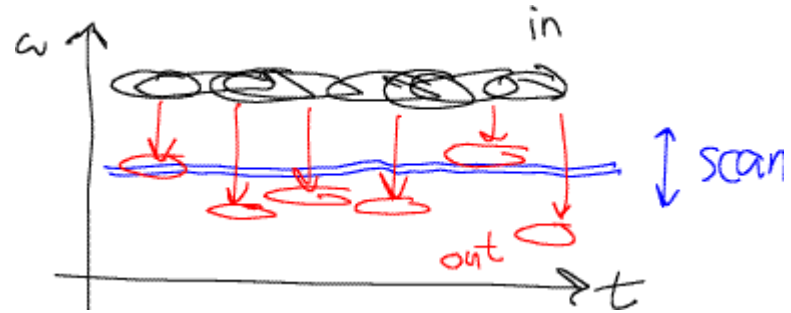
Conclusions

- IXS is really measurement of violation of Liouville's theorem
- Coherent pulses \Rightarrow more ways to do this than with monochromators
- X-FROG is one possible alternative

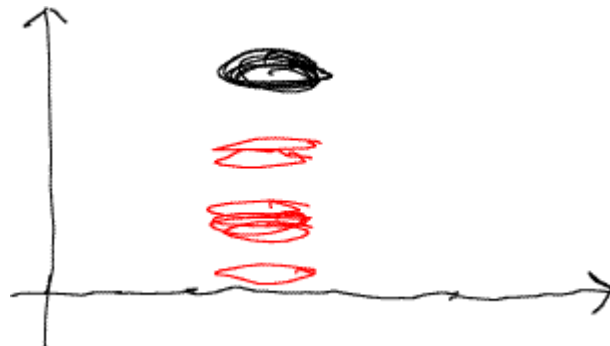
Detectors: Windowless APDs

How to probe decoherence?

Incoherent pulses - "regular IXS"



Coherent pulses



Can detect instead through noise pattern in $E(t)$