

# Time-resolved Diffuse Scattering: phonon spectroscopy with ultrafast x rays

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SLAC National Accelerator Laboratory*

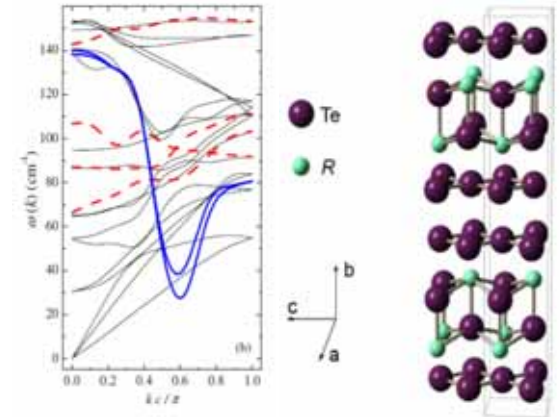
# phonons play defining role in materials properties



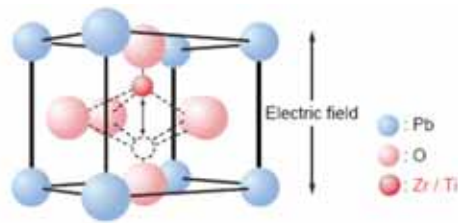
## Electrical and thermal conductivity



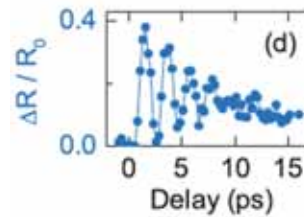
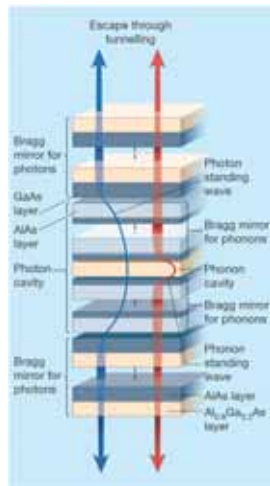
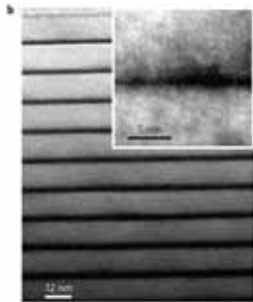
## materials with coupled degrees of freedom



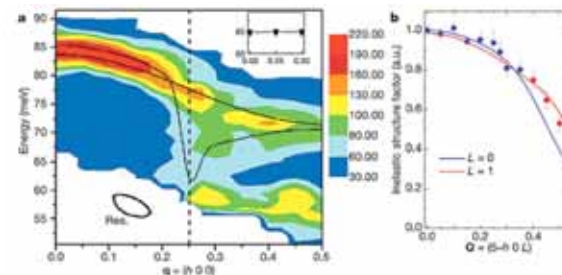
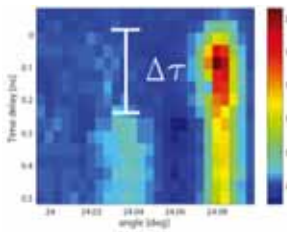
## Ferroelectricity



## Acoustic phonon "engineering"

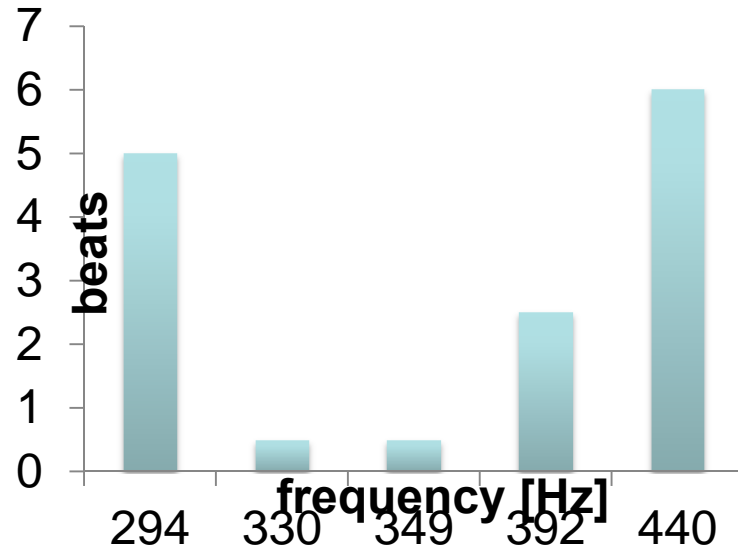


## Superconductivity



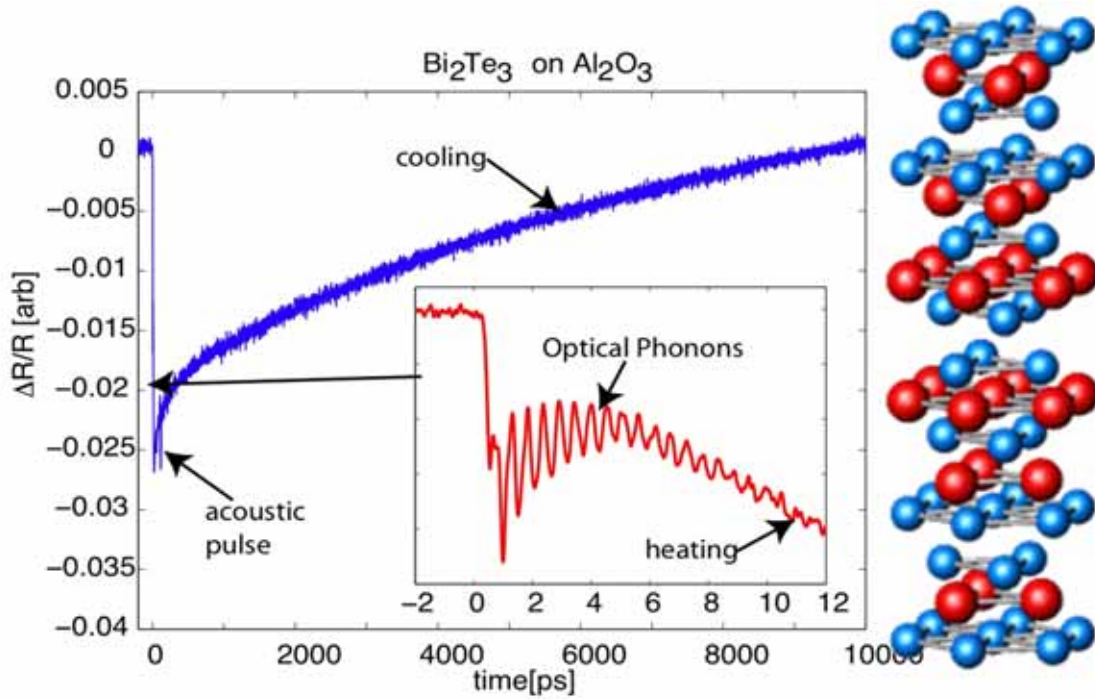
## Ultrafast Tickle and Probe

# Why do physics in the time domain?

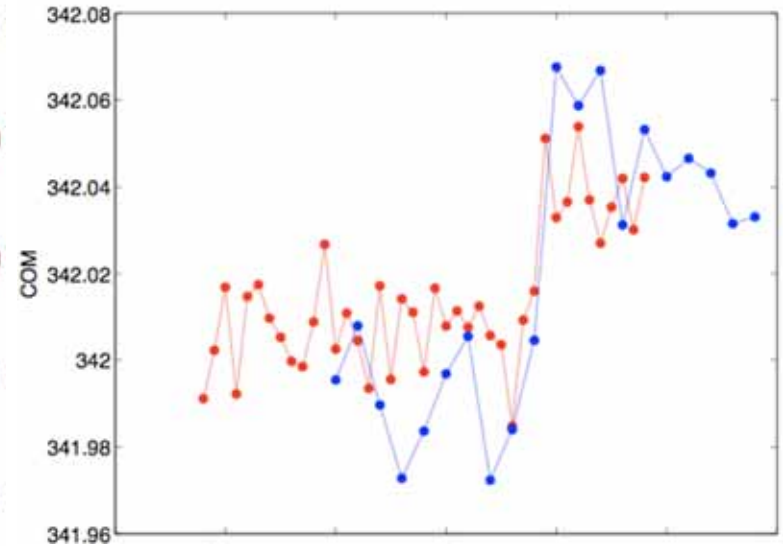


**Spectra alone are  
sometimes  
inadequate**

# Separation of time-scales

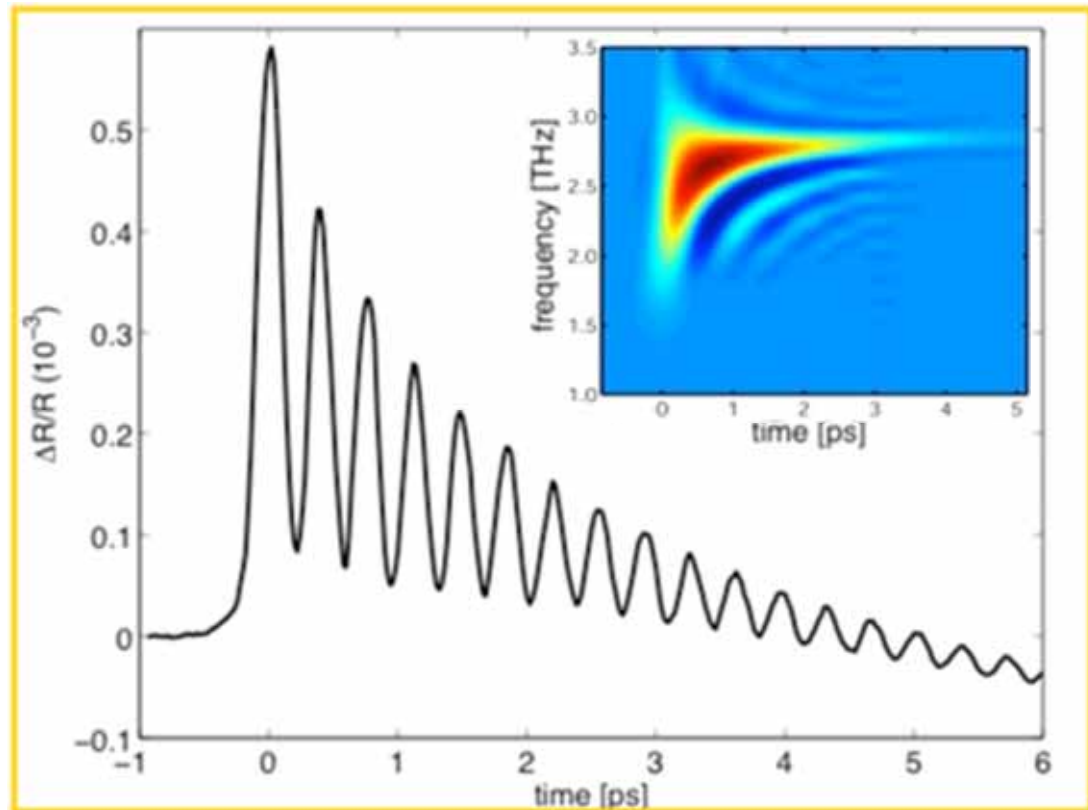
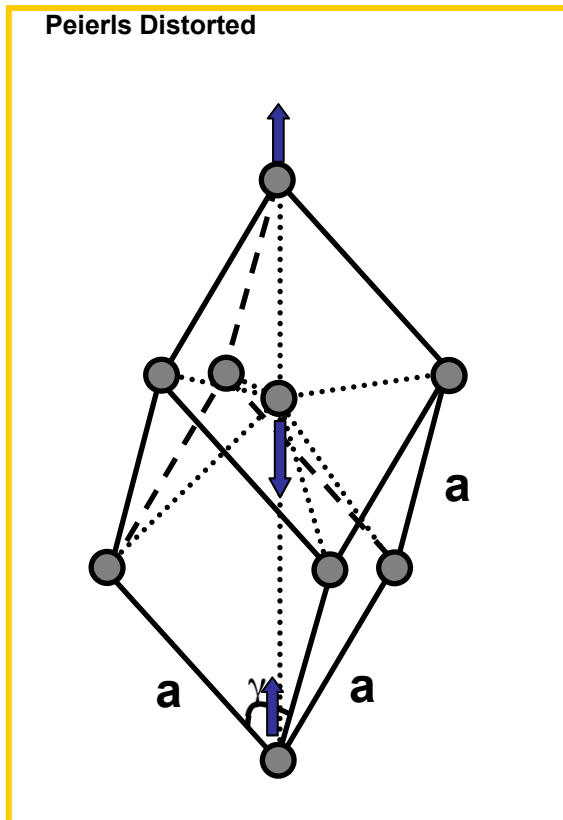


Sheu *et al.* unpublished



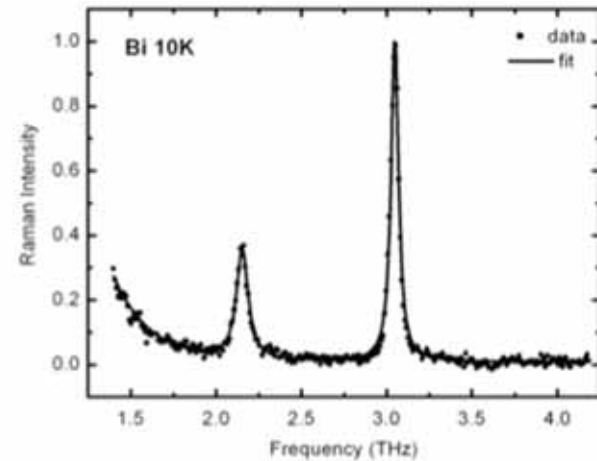
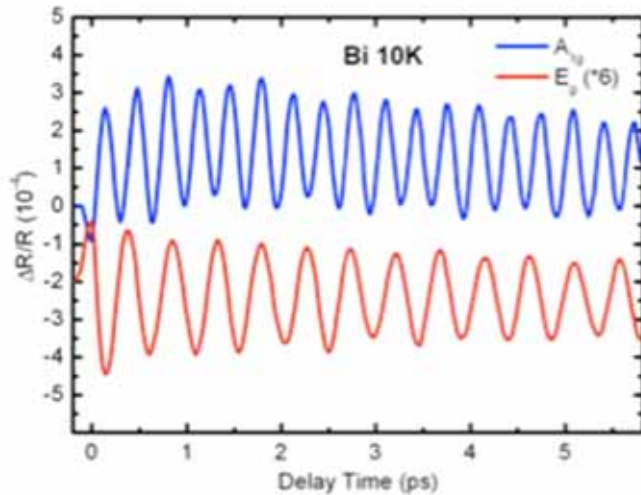
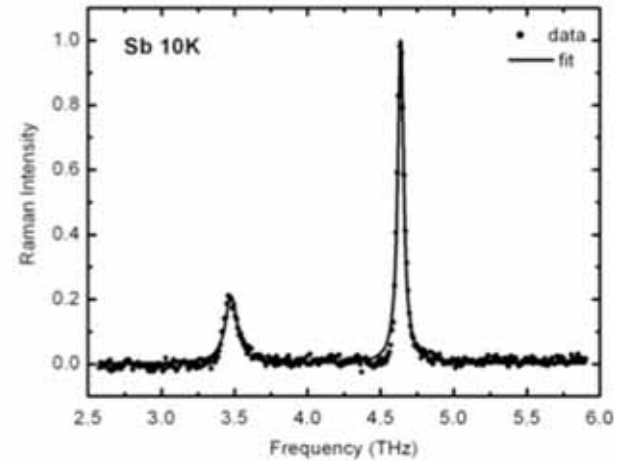
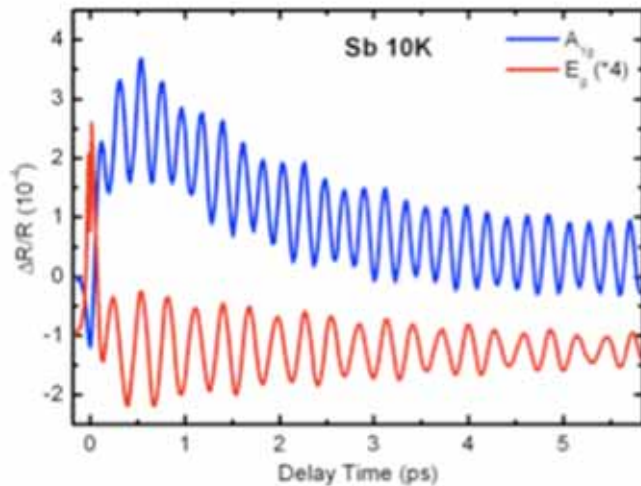
Trigo *et al.* unpublished

Example: Photoexcited bismuth (all optical experiments)



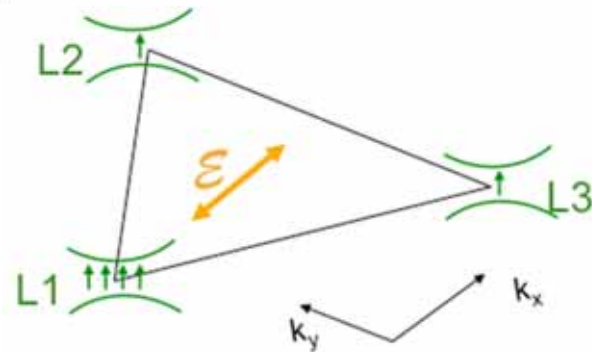
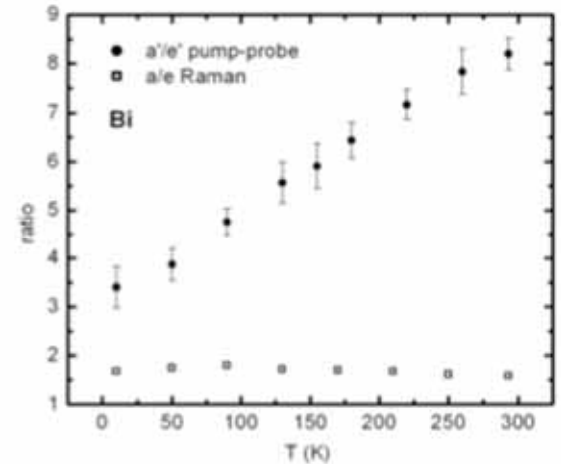
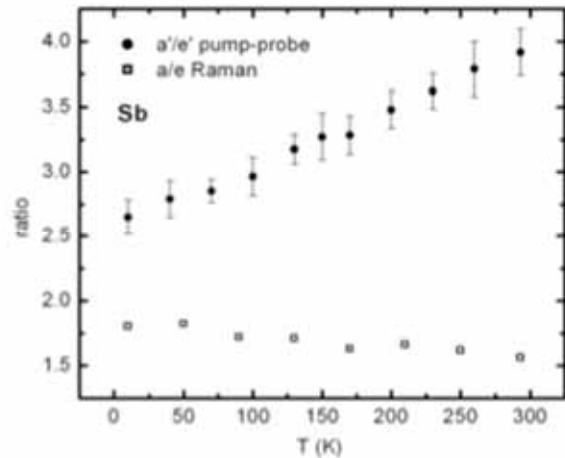
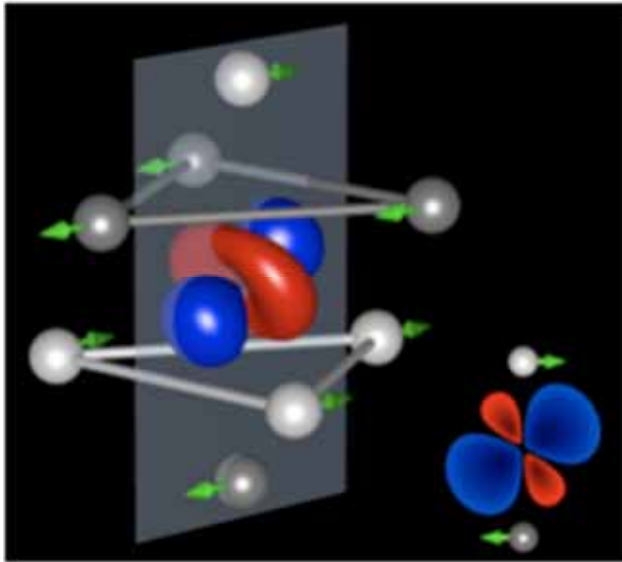
Coherent  $A_{1g}$  Mode is strongly softened and chirped.

# Low excitation (tickle regime) Pump-probe vs. Spontaneous Raman



Collaboration with Roberto Merlin, UM

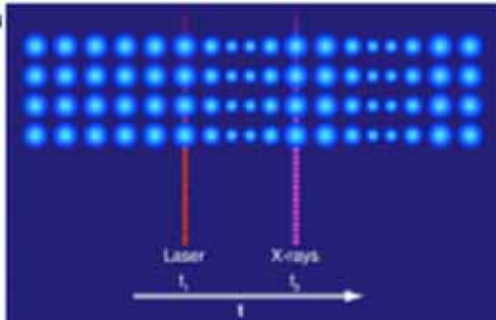
# Subtle difference between the two...



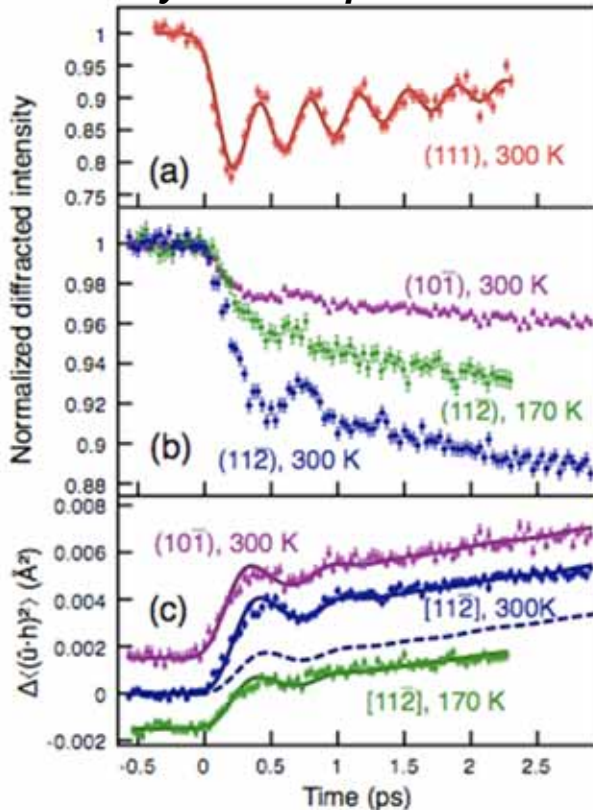
$$F(t) = \begin{cases} F_0 e^{-\Gamma t} & t \geq 0 \\ 0 & t < 0 \end{cases}$$

$$Q(t) = \int_{-\infty}^t \frac{\sin[\Omega(t-\tau)]}{\Omega} e^{-\gamma(t-\tau)} F(\tau) d\tau = \frac{F_0 e^{-\Gamma t}}{\Omega^2 + (\Gamma - \gamma)^2} - \frac{F_0 e^{-\gamma t} \cos(\Omega t + \varphi)}{\Omega \sqrt{\Omega^2 + (\Gamma - \gamma)^2}}$$

# Electronic Softening in Bi by femtosecond X-ray Diffraction

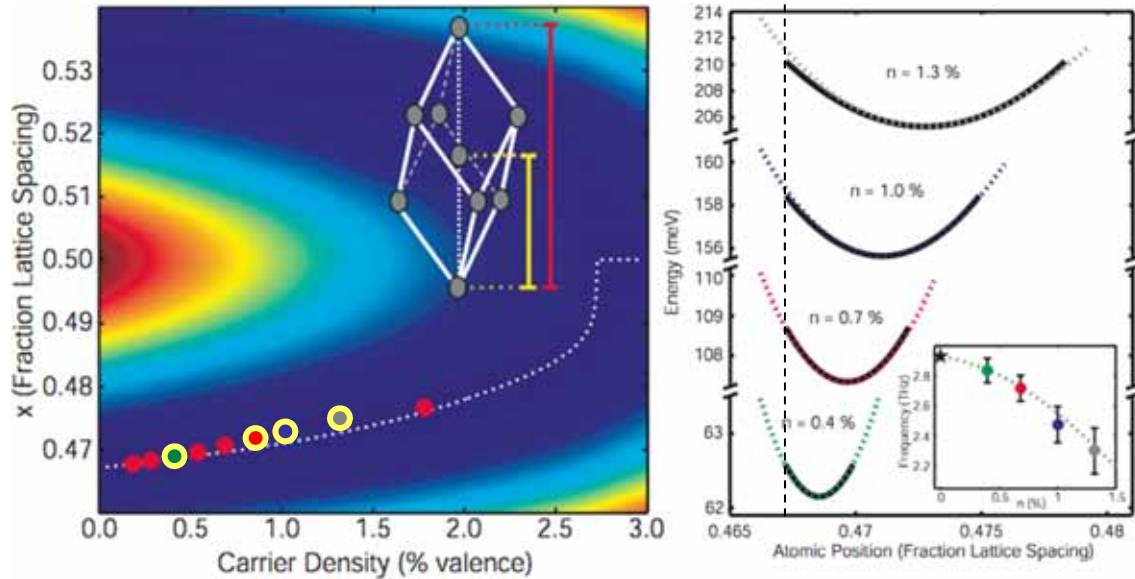


Reis *Physics Viewpoint* 2009



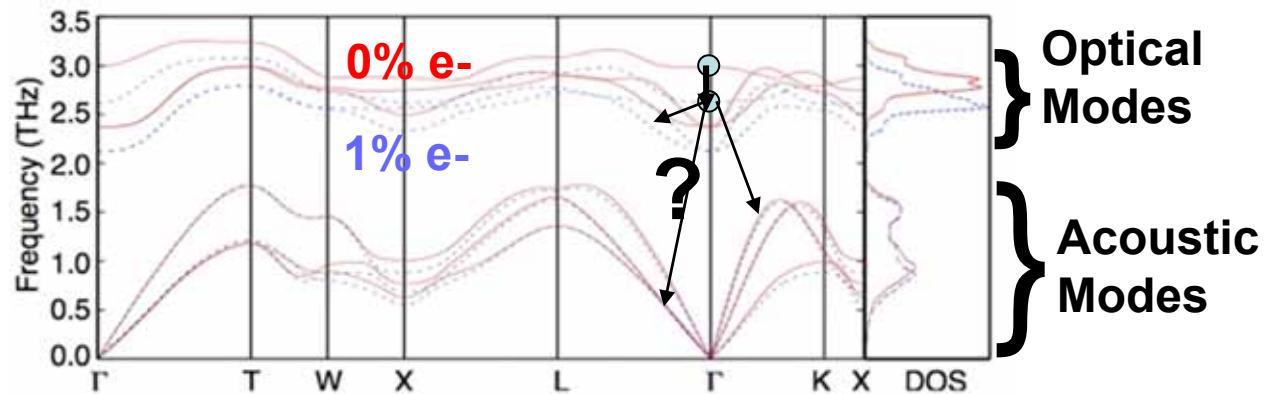
Johnson *et al.* PRL 2009.

Ultrafast Tickle and Probe



D. M. Fritz *et al.* Science 315, 2007.

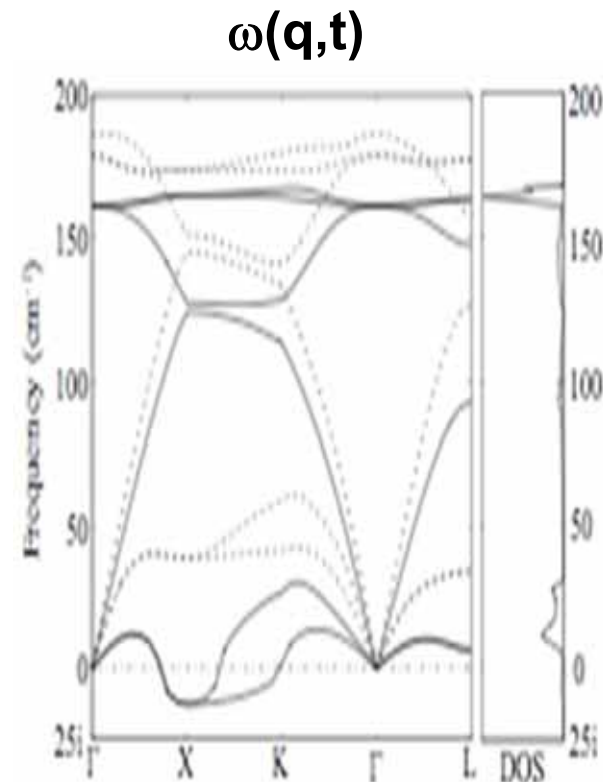
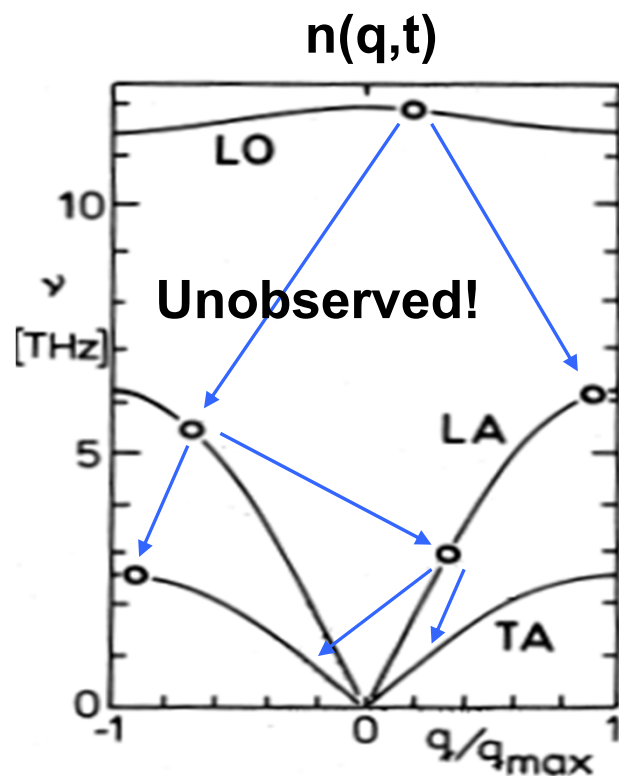
## DFPT calculation



Murray *et al.* PRB 75 2007.

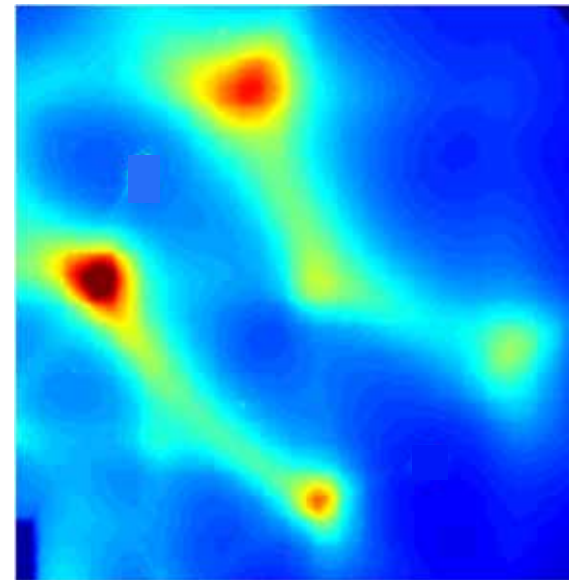
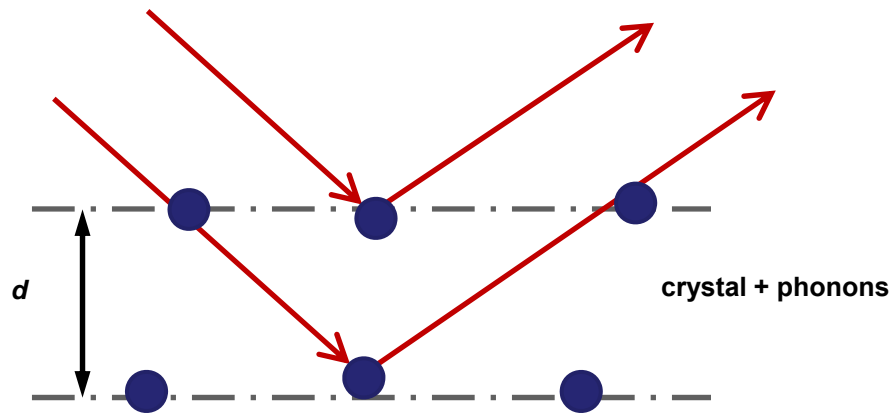


# Momentum and Time-resolved Phonon Spectroscopies virtually nonexistent



phonon-phonon and electron-phonon coupling, interatomic forces.

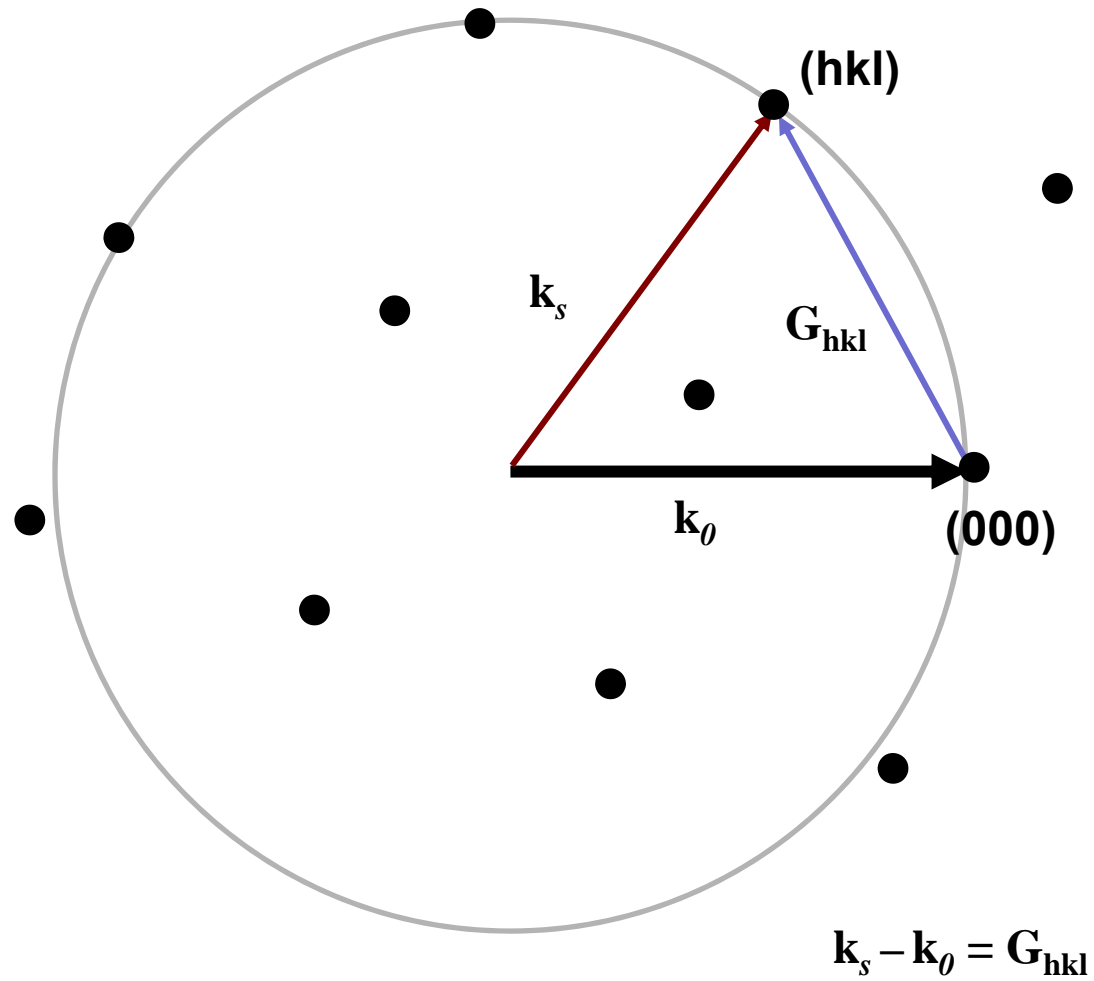
# x-ray diffuse scattering: measure deviations from average structure



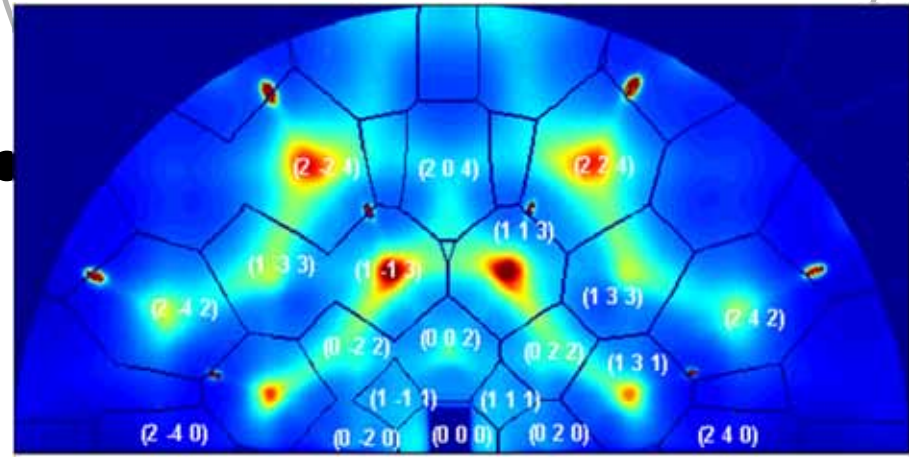
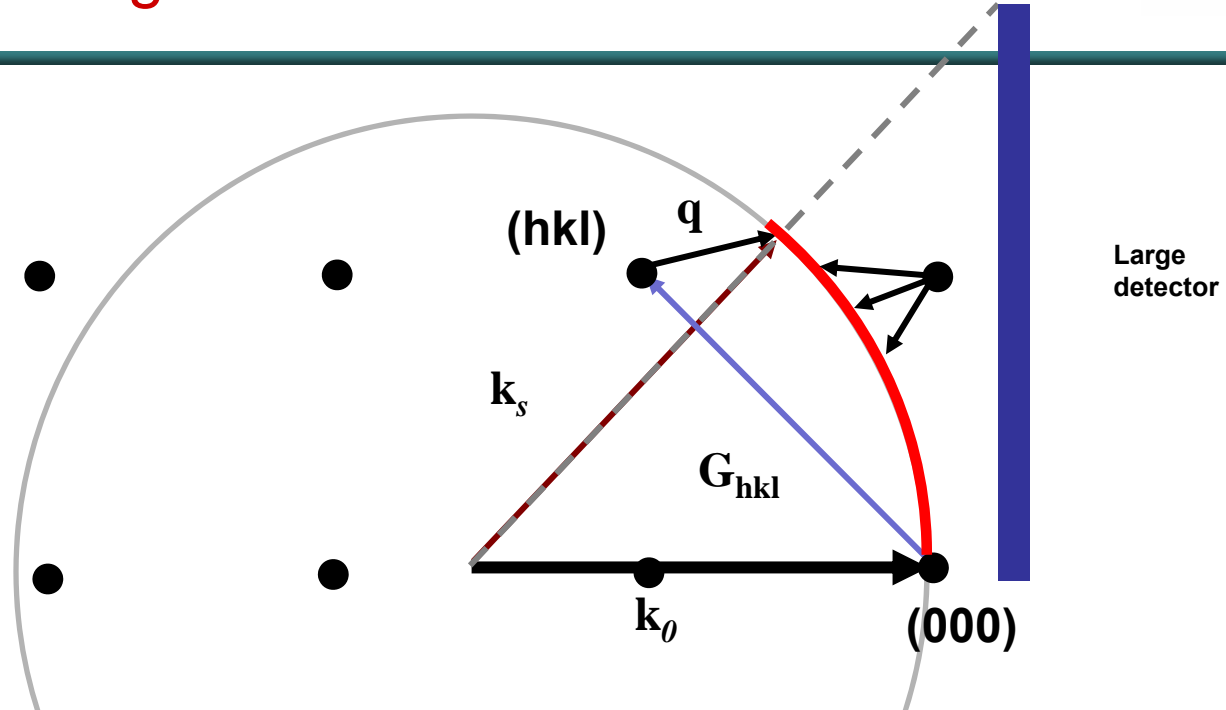
Very weak, we need bright x-ray pulses!

Example, diffuse scattering by thermal phonons

# Bragg scattering

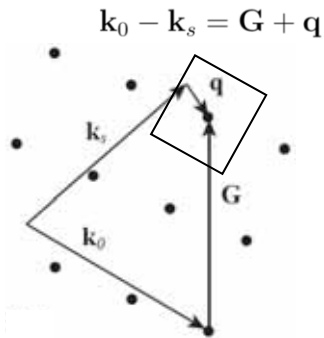
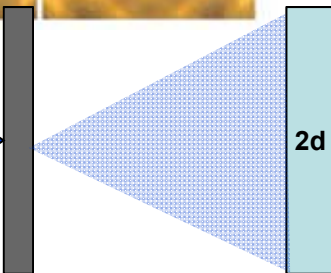
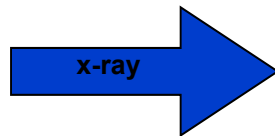
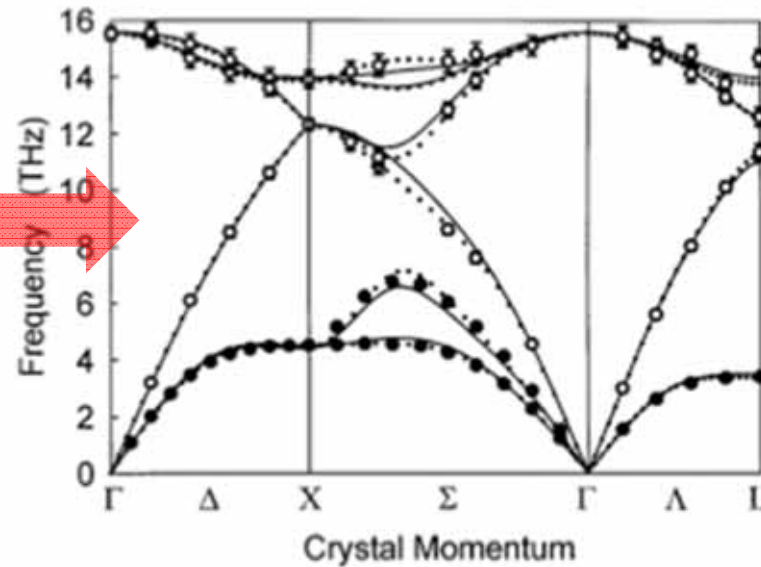
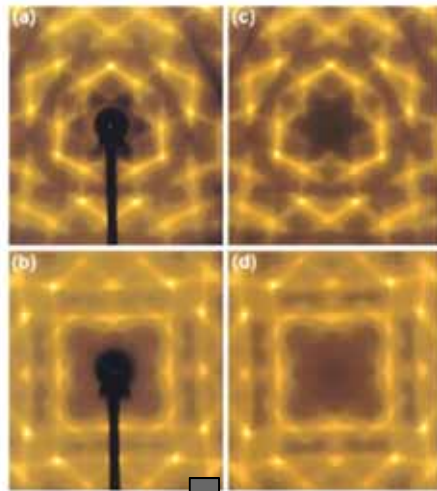


# Diffuse scattering



$$I(\mathbf{Q}, t) \propto \sum_j (\mathbf{Q} \cdot \epsilon_{\mathbf{q}j})^2 \frac{2n_{\mathbf{q}j} + 1}{\omega_{\mathbf{q}j}}$$

# Phonon Dispersion from TDS and limitations



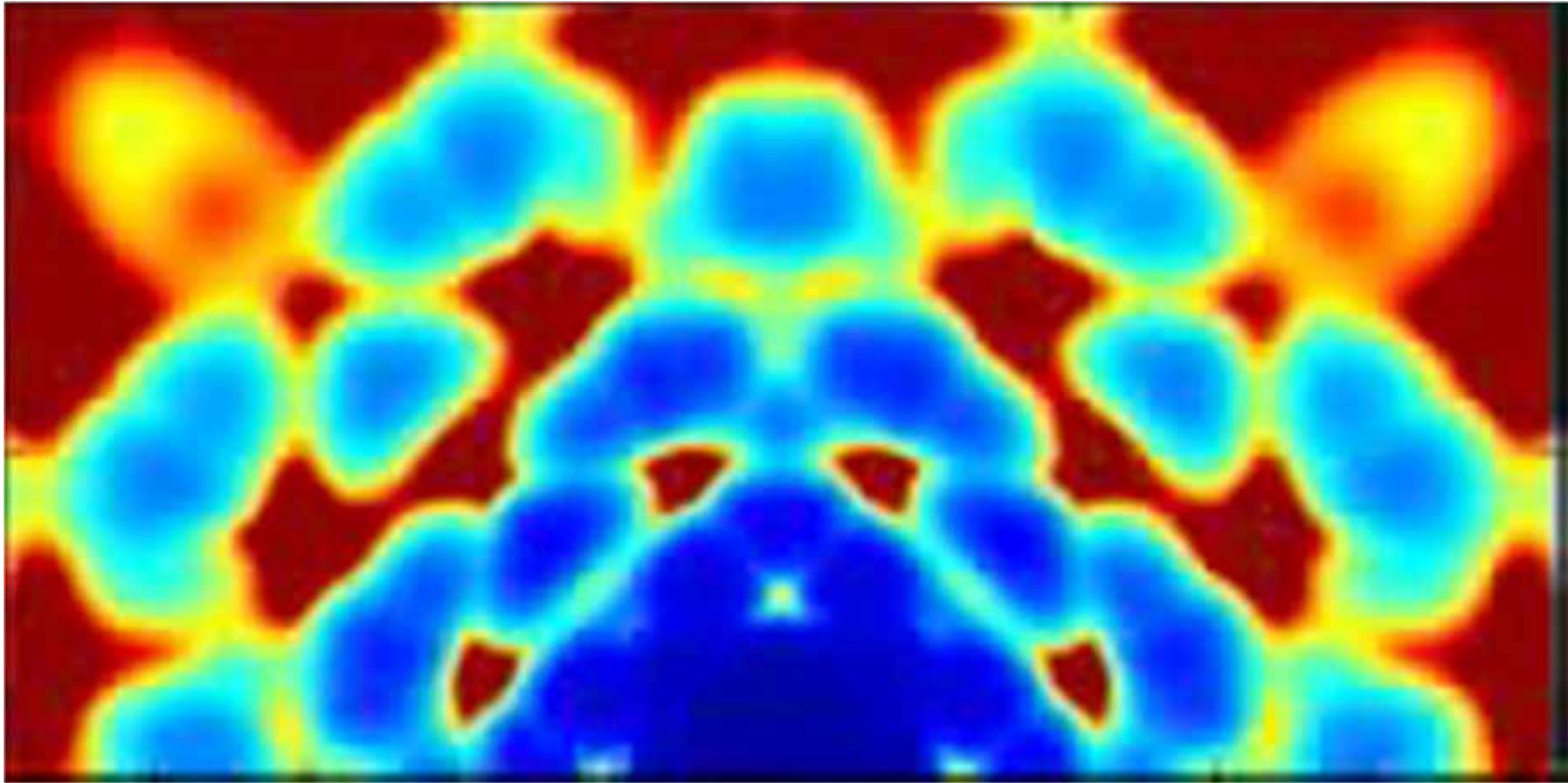
Joynson, Phys. Rev. 94, 851 (1954)...  
 ...M. Holt et al., PRL 83, 1999.

$$I(\vec{Q}, t) \propto (\vec{Q} \cdot \vec{\epsilon}_{\lambda, \vec{q}})^2 \frac{(2n_{\lambda, \vec{q}}(t) + 1)}{\hbar\omega_{\lambda, \vec{q}}}$$

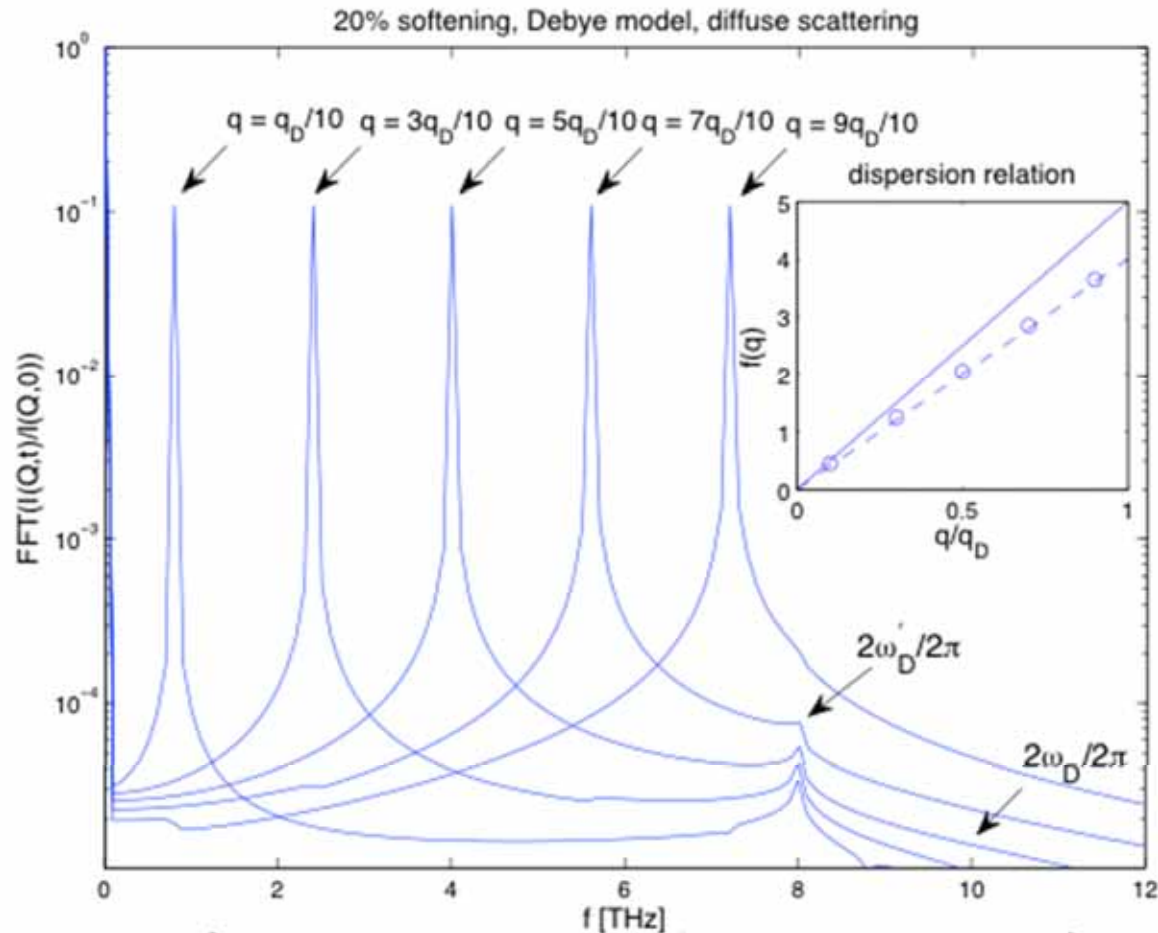
**TDS: Limited to simple cases (# fit parameters low) and have a constraint (assumes Bose-Einstein distribution)**

# Simulation of InP impulse softening of TA by 20%

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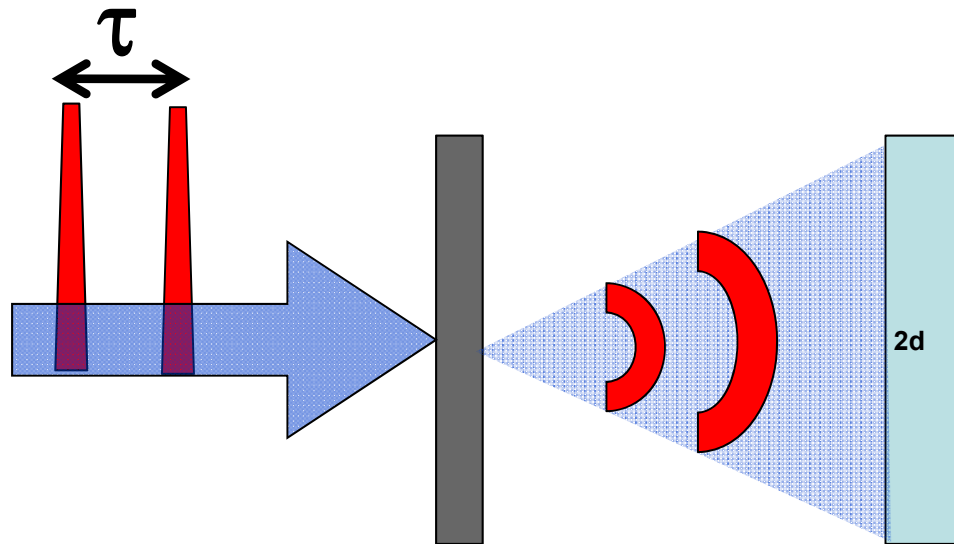
# Fourier transform of $I(q,t)$ yields phonon dispersion (excited state)



$$\langle u_q(t)^2 \rangle = \frac{1}{2} \langle u_q(0)^2 \rangle \left( \alpha (1 + \cos(2\alpha\omega(q)t)) + \frac{1}{\alpha} (1 - \cos(2\alpha\omega(q)t)) \right),$$

$$\langle u^2(t) \rangle = \frac{9k_B T}{2M\omega_D^2} \left( \left( 1 + \frac{1}{\alpha^2} \right) + \frac{\sin(2\alpha\omega_D t)}{2\alpha\omega_D t} \left( 1 - \frac{1}{\alpha^2} \right) \right).$$

Alternatively, can use multiple coherent pulses

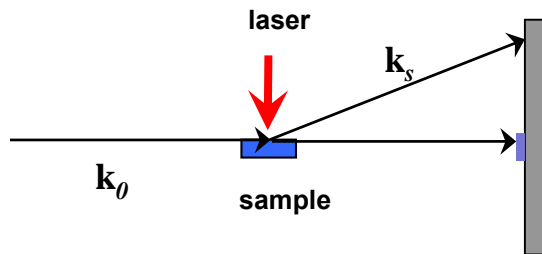


If you want to go really crazy, multiple colors, incidence angle, stimulated x-ray (electronic) Raman selected  $w$  &  $q$ ...

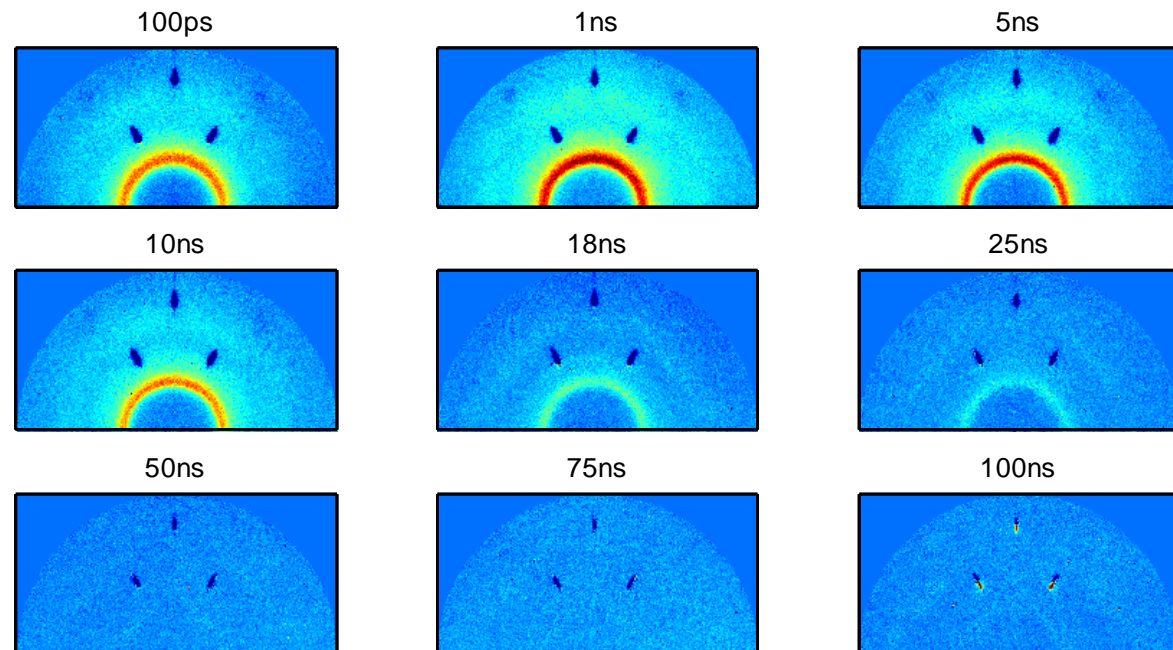
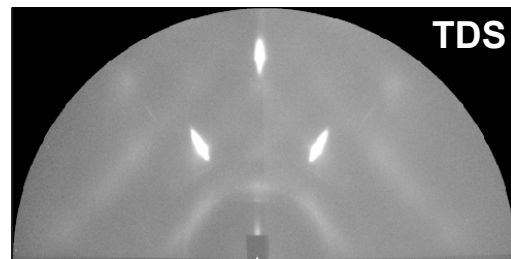
Can we do x-ray 4-wave mixing, ala K. Nelson or S. Mukamel?



# Synchrotron data limited by time-resolution



## Laser-melt and epitaxial regrowth of Bi on Sapphire

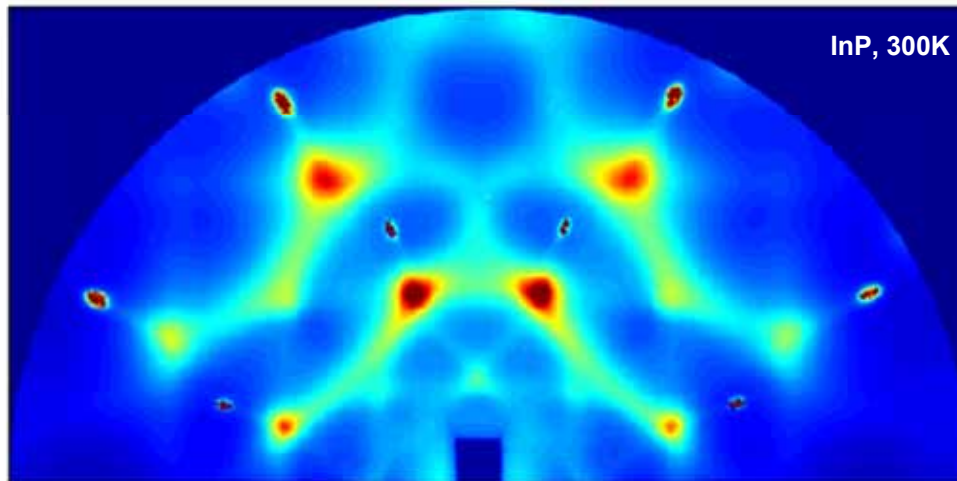
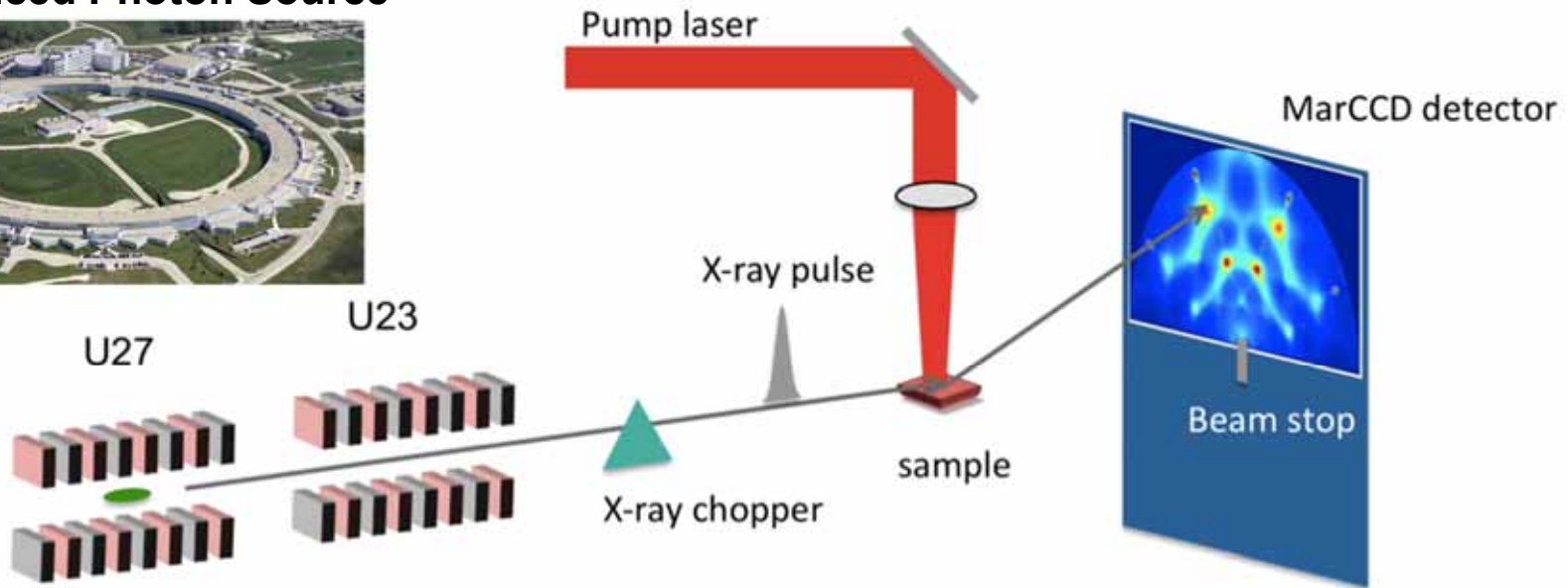


Average of 5 shots

BioCARS beamline at APS  
~1% of LCLS photons/pulse but 100ps

# Time-resolved x-ray diffuse scattering with 100ps resolution

## Advanced Photon Source



15 keV x-rays

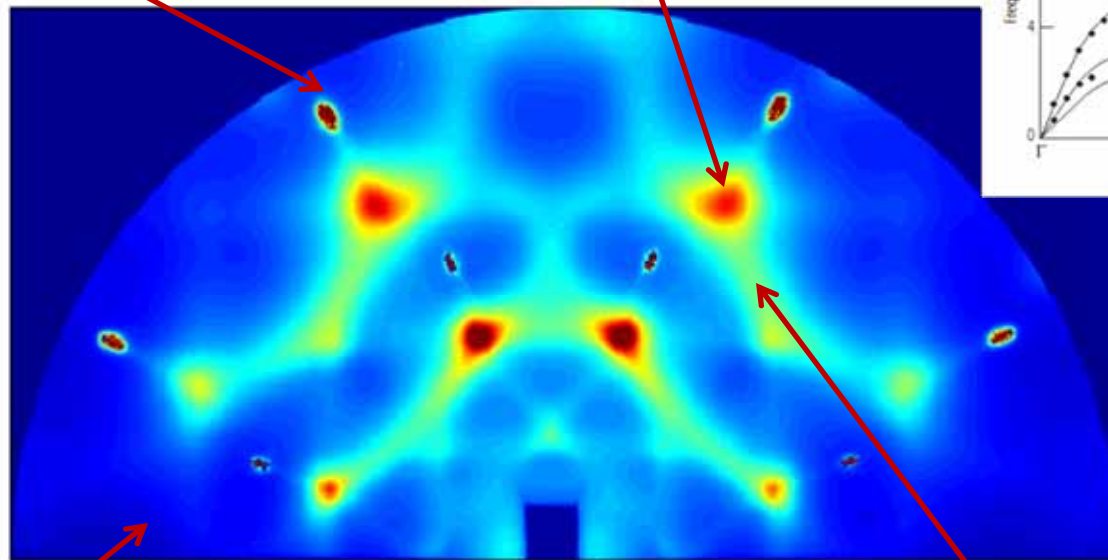
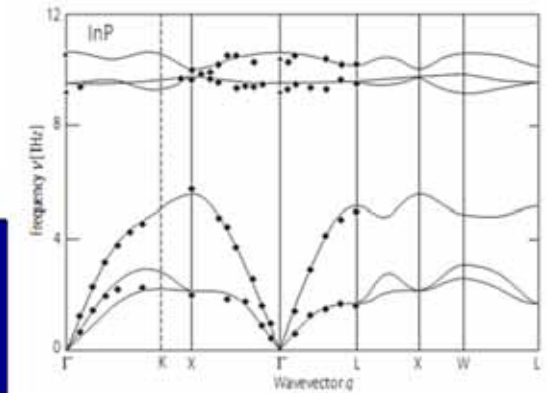
~ 100 single x-ray Pulses

Equivalent to a single LCLS shot!

# Primarily TDS

“Bragg Rods” from surface scattering

bright spots due to acoustic phonons near **zone center**



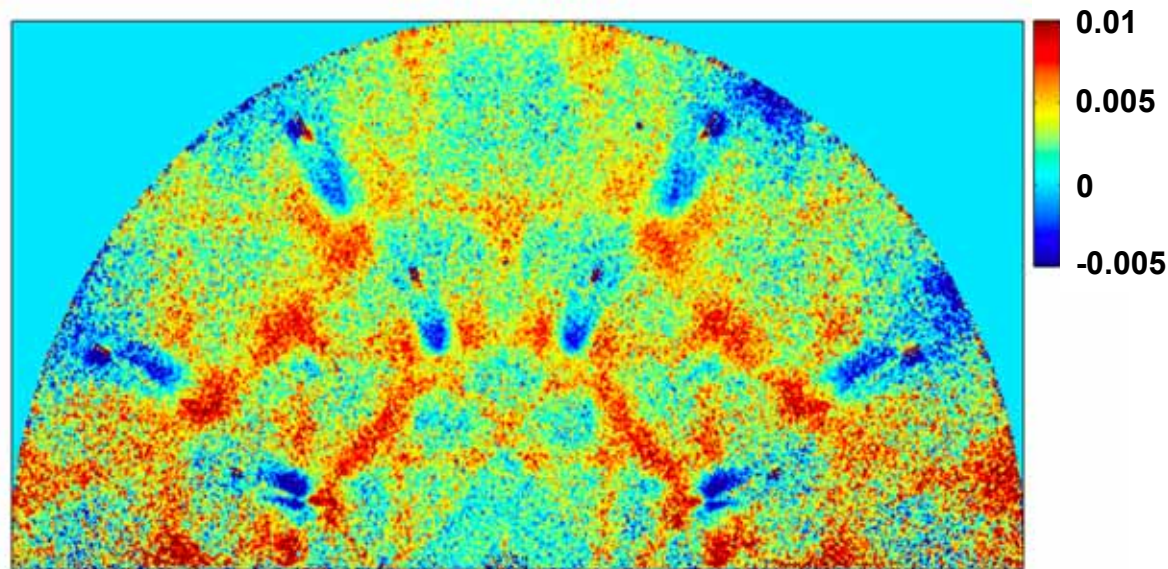
p-Polarization forbids scattering at  $90^\circ$

Beam stop shadow

High-symmetry directions, where acoustic branches are “soft”.

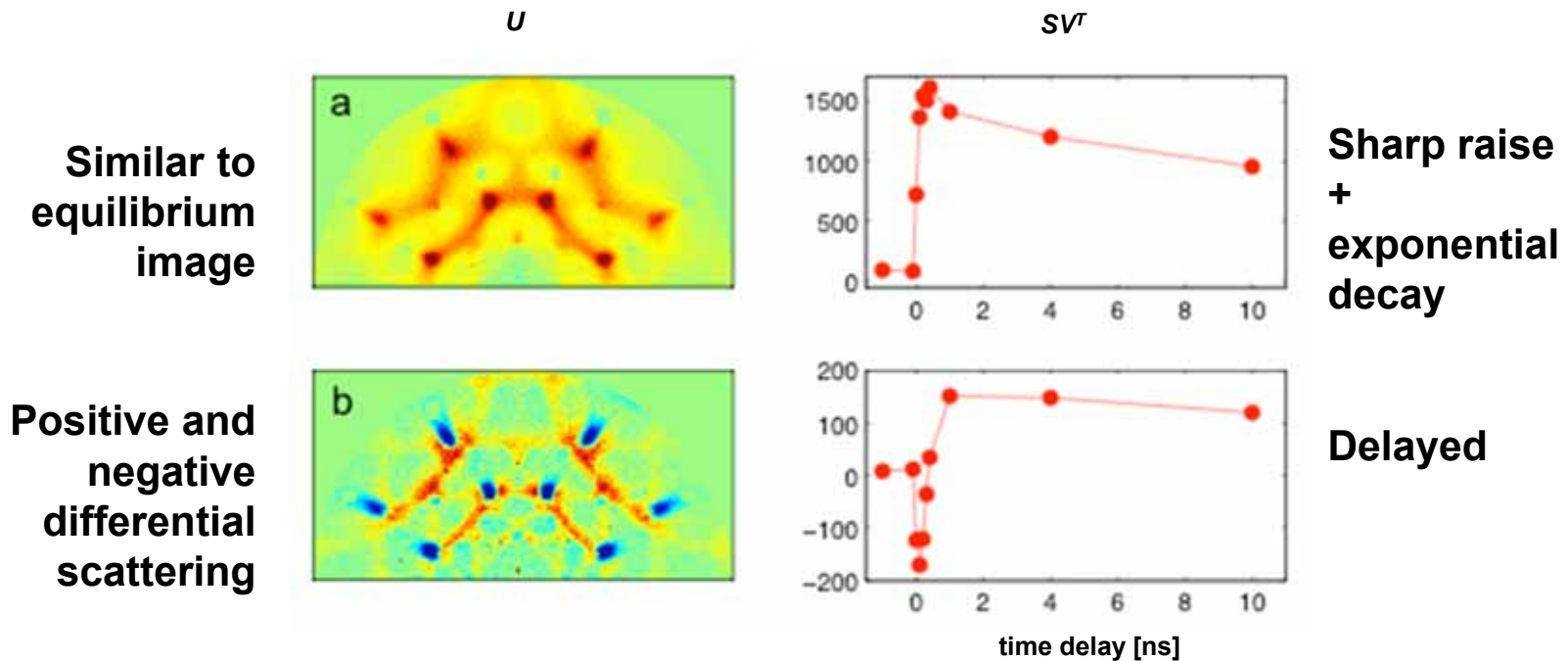
# But, more than heating

$$[ I(400\text{ps}) - I(100\text{ps}) ] / I(\text{off})$$



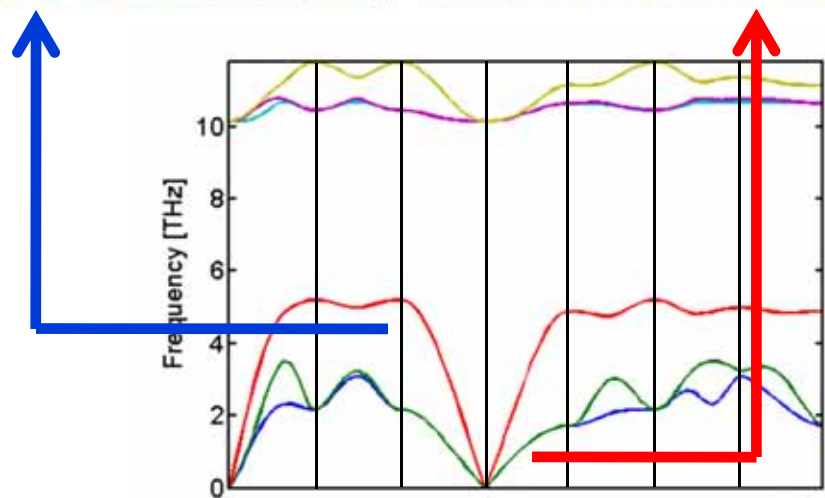
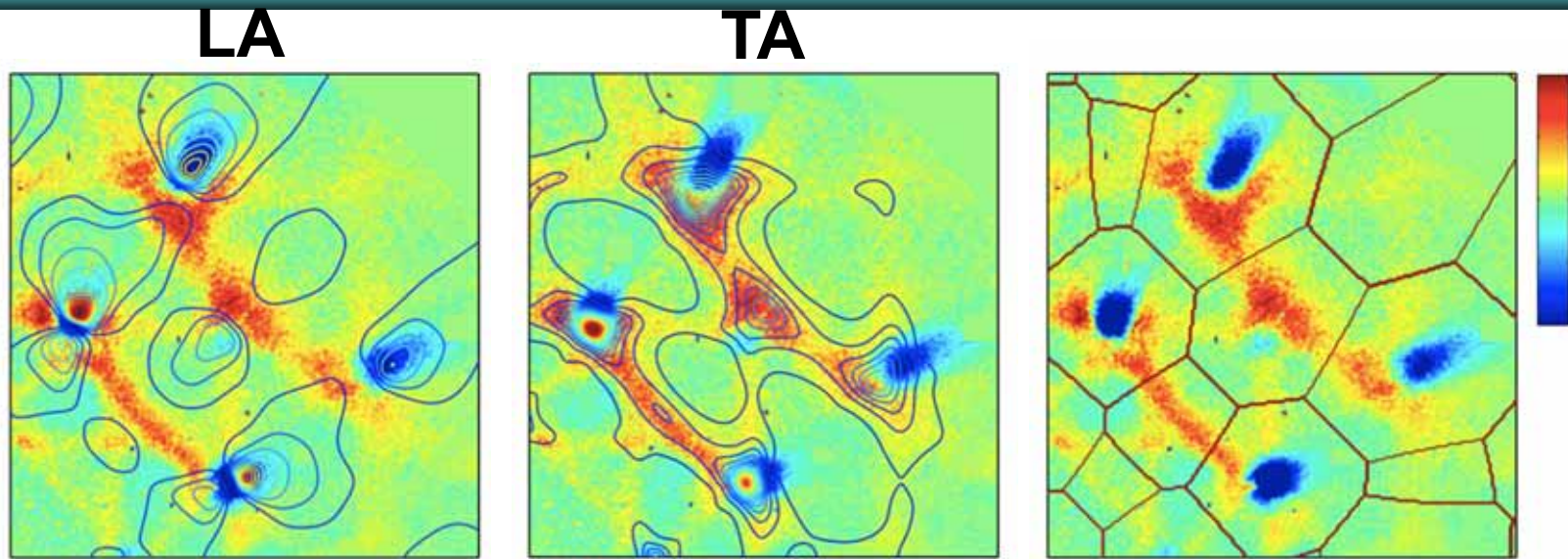
If processes were only thermal,  $\frac{\partial I(t)}{\partial t} \propto \Delta n \rightarrow \frac{1}{I_0} \frac{\partial I(t)}{\partial t} \propto \Delta T = \text{const}$

# Interpretation by Singular Value Decomposition

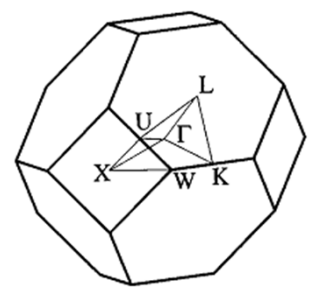


Complex dynamics in the phonon populations due to the anharmonic coupling between modes

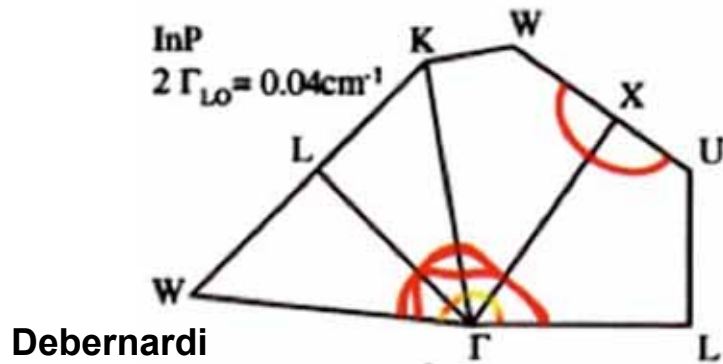
# Contribution from acoustic phonon branches



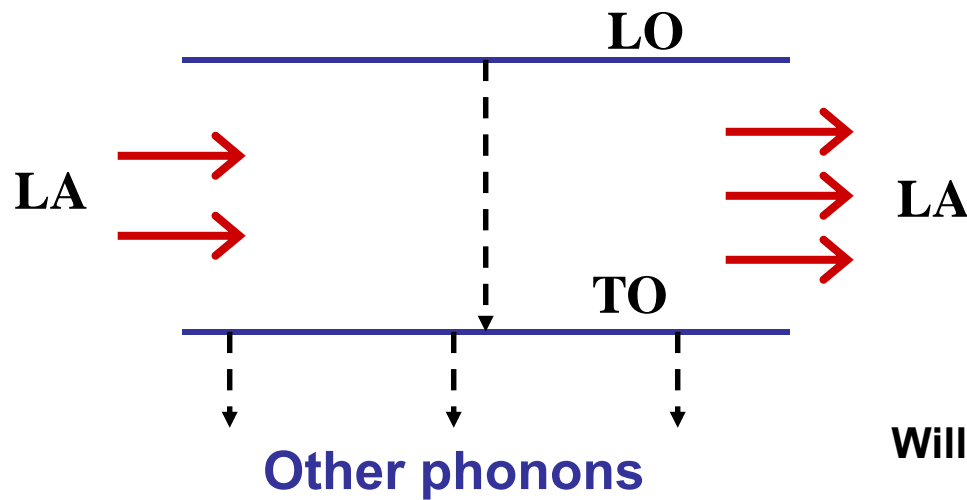
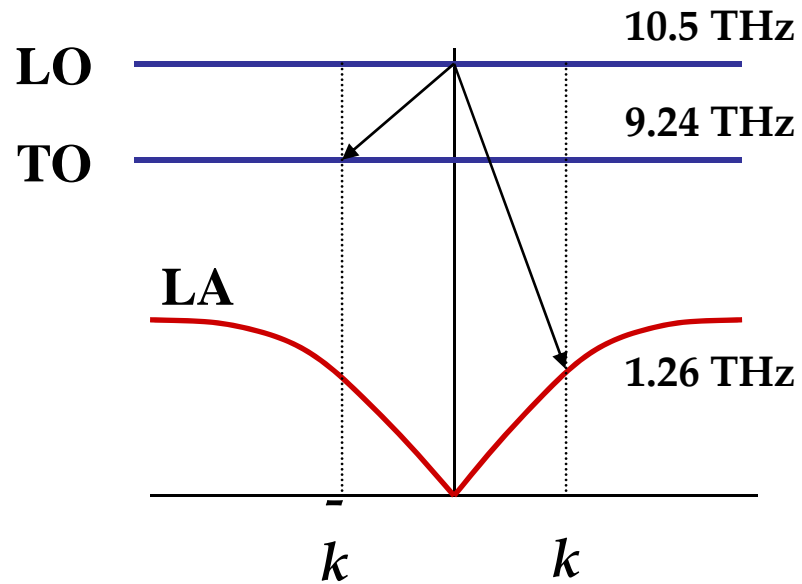
Brillouin zones



# Can we modify optical phonon lifetime?



**Decay channel:  
 LO  $\rightarrow$  TO + LA**



**Stimulated decay  
 and emission of  
 TO + LA**

Chen, Khurgin, Merlin, APL 80 2901, 2002.

Will study anharmonic decay at LCLS this fall!

# Acknowledgements

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Shambhu Ghimire, Vinayak Vishwanath, PULSE, SLAC, Stanford

*Stanford PULSE Institute, SLAC National Accelerator Laboratory*

Yu-Miin Sheu,  
*Los Alamos*

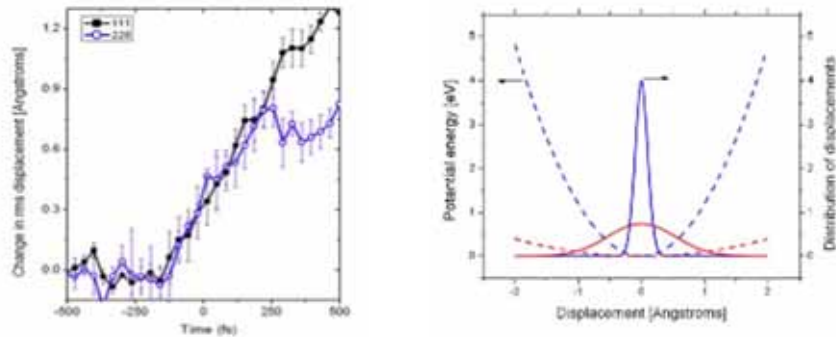
Tim Graber, Robert Henning,  
*CARS, U. Chicago*

Supported by the U.S. Department of Energy, Office of Basic Energy Science

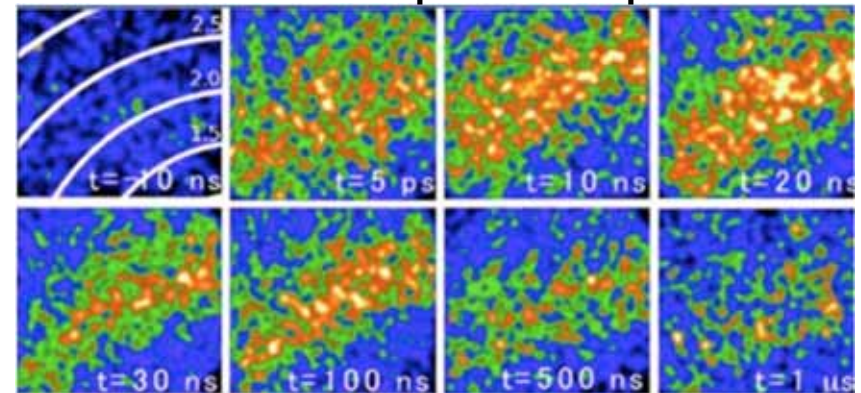


# Lattice instabilities

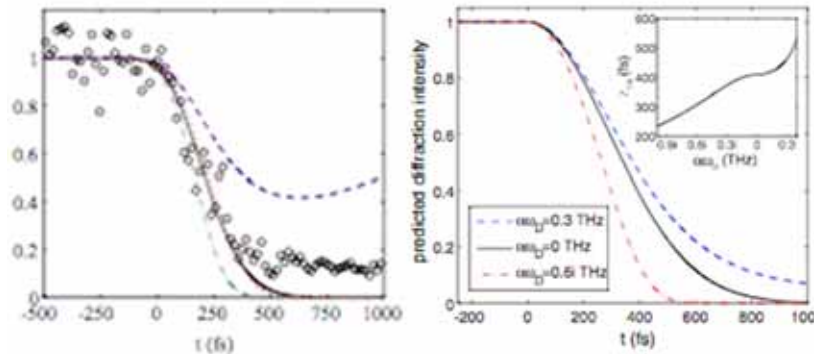
Diffraction data consistent with complete softening



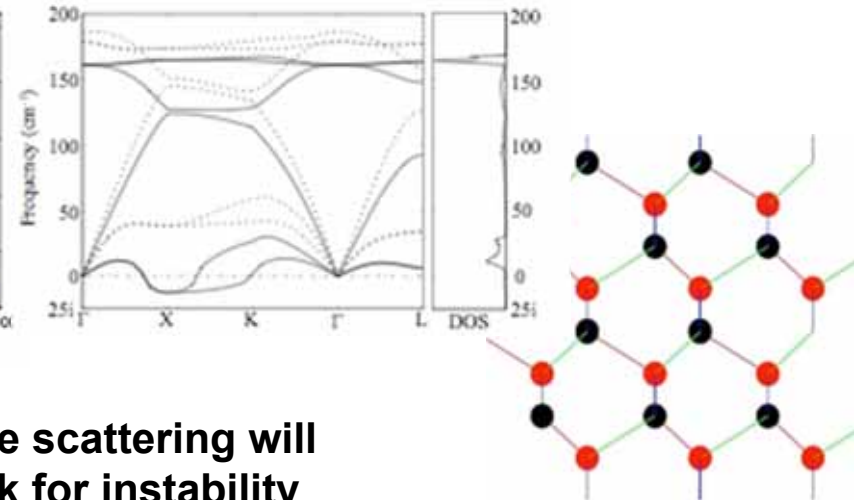
formation of a nonequilibrium liquid



Model assuming uniform softening  
Gives similar results to inertial dynamics



DFPT predicts instability first develops at X point



LCLS beamtime this fall to study Ultrafast diffuse scattering will measure evolution of interatomic forces and look for instability