SYNCHROTRON RADIATION ANALYSIS OF THE SUPERKEKB POSITRON STORAGE RING

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Abstract

We report on modeling results for synchrotron radiation absorption in the SuperKEKB storage ring vacuum chamber including the effects of photon scattering on the interior walls. A detailed model of the geometry of the inner vacuum chamber profile, including roughness parameters, has been developed and used as input to a photon tracking code. Particular emphasis is placed on the photon absorption rates in the electron-positron interaction region.

INTRODUCTION

The SuperKEKB e+e- collider will provide highprecision measurements of production and decay of bound states of B quarks in the interactions of 4-GeV positrons with 7 GeV electrons. Commissioning is scheduled to begin in 2016. The design of the vacuum system [1] incorporates a variety of countermeasures to limit electron cloud buildup in the positron ring. The Cornell Electron Storage Ring Test Accelerator (CESRTA) program [2] in operation since 2008 has provided a wealth of information on the efficacy of such mitigation techniques via extensive measurements and model development which has also been used in the design of the International Linear Collider positron damping ring. [3,4]. Here we report on results from a detailed modeling study of synchrotron radiation photon scattering and absorption on the interior surfaces of the vacuum chamber in the positron ring.

VACUUM CHAMBER MODEL

Standard beam pipes in the positron storage ring have a circular beam channel and two rectangular antechambers on either side, as shown in Fig. 1. Synchrotron radiation photons irradiate the outer wall of the antechamber on the outside of the ring. The diameter of the beam channel and the half width at the horizontal axis are 90 mm and 110 mm, respectively. Strip-type non-evaporable getters (NEG) pumps are installed in the antechamber on the inside of the ring. The pump channel is connected to the beam channel through a screen. A cross section of the standard beam pipe in the vacuum chamber model for the present calculations is shown in Fig. 2. The screens for the NEG pumps are defined as photon-absorbing walls in the model. In the wiggler sections, the photons are strike both sides

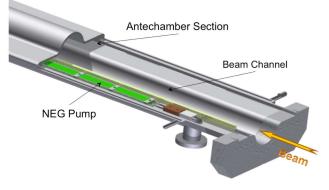


Figure 1: Schematic drawing of a beam pipe in arc sections of the positron storage ring

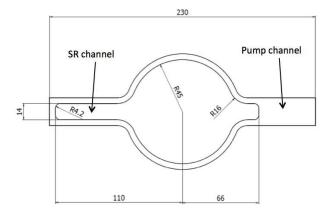


Figure 2: Cross section of the vacuum chamber model in the arc sections

of the beam pipe, so both antechambers are used as channels for photon absorption. Circular beam pipes have been adopted in the interaction region (IR). The diameters vary in a staircase pattern from 80 mm at the entrance of the final focusing magnets (QCs) to 20 mm at the interaction point (IP). Beam pipes in the QCs and IP are shown in Figs. 3 and 4, respectively. The beam pipes nearest the IP are tapered from 20 mm to 10 mm in order to cut off the photons to the Belle-II detector. The beam pipes also have ridge structures on the inner surface to diffuse the photons, but these structures are not included in this model.

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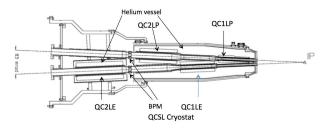


Figure 3: Schematic view from above of the beam pipes in the cryostat on one side of the IP. The positron beam exits on the left in the upper of the two pipes shown.

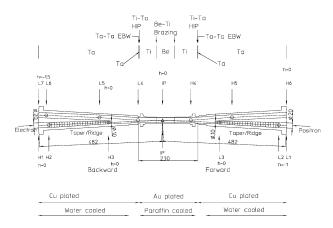


Figure 4: Beam pipes around the interaction point

SYNCHROTRON RADIATION PATTERN

Figure 5 shows a calculation of the synchrotron radiation pattern on the outside wall of the the 4-GeV positron ring. The calculation incorporates an analytical calculation of the synchrotron radiation rate from bending radii of the dipole magnets only and the distance of the wall from the beam in the horizontal plane containing the beam. The element type in which the photons hit the wall is color-coded as shown. The linear density per positron of incident photons is averaged over 1-m intervals. The rate is less than 2 photons/m/e+ in the arcs, with hot spots reaching 10 photons/m/e+ downstream of the damping wiggler sections. No photons are directly incident on the outer wall within 10 m of the interaction point.

PHOTON TRACKING CODE SYNRAD3D

Analysis of the the CESRTA measurements of electroncloud-induced coherent tune shifts motivated the development of the Monte Carlo-based photon scattering and tracking code Synrad3D [5], since the important contribution of electron cloud buildup in dipole magnets depends on a quantitative understanding of photoelectron production on the top and bottom of the beam pipe near the vertical plane containing the beam [6]. The code implements the tabulated information for X-ray specular and diffuse scattering from the LBNL database [7]. A lower limit on the absorbed photon energies of 4 eV was imposed for the simulations discussed here.

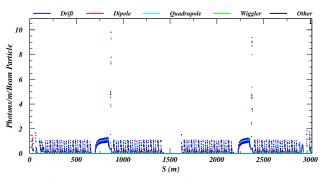


Figure 5: Analytic calculation of the synchrotron radiation pattern directly incident on the outside wall of the 4-GeVpositron ring. The linear density in photons/m/e+ is averaged over 1-m intervals and the lattice element type in which the photons strike the wall is coded as shown.

DISTRIBUTION OF ABSORBED PHOTONS

Figure 6 shows characteristics of photons absorbed around the entire ring, including the effects of photon scattering. Figure 6 a) shows the transverse absorption coordi-

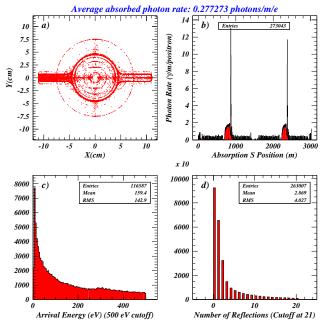


Figure 6: Distributions of absorbed photons around the ring including the effects of photon scattering. a) Scatter plot of vertical versus horizontal photon absorption coordinates integrated over the ring circumference, exhibiting the variety of transverse vacuum chamber profiles, including the antechambers. b) Distribution of absorbed photon locations in the arc coordinate S of the positron orbit, ranging up to 3016.31 m. c) Distribution of absorbed photon energies below 500 eV. d) Number of wall reflections undergone by the photons prior to absorption.

nates integrated over the ring circumference, exhibiting the variety of beam pipe cross sections in the vacuum system design, including the antechambers and the 1-cm-diameter pipe in the IR. Figure 6 b) shows the linear density per positron of absorbed photons. This density increases by about a factor of two to 2 photons/m/e+ in the wiggler regions when photon scattering is taken into account, and increases to 12-14 photons/m/e+ in the two hot spots downstream of the wiggler. The average energy of absorbed photons is about 160 eV, as shown in Figure 6 c). Figure 6 d) shows that about 90k of the 263k modeled photons are absorbed without previously having scattered. The average number of scatters around the ring prior to absorption is 6.1.

The rate of photons absorbed in the IR is of concern for electron cloud buildup in regions of high values for the beta functions [8]. The design of the final focus magnet system includes 8 superconducting magnets located with 3 m of the e+e- interaction point [9, 10]. Here we report on modeling results for the example of the magnet closest to the IP in the positron ring, the 33.4-cm-long QC1RP magnet, the center of which is located 935 mm from the IP. The field gradient in this quadrupole magnet is 68.7 T/m. The analytic calculation of the synchrotron radiation incident on the beam pipe wall and the photon scattering model with scattering turned off calculate a low rate of incident photons of about 0.04 photons/m/e+ originating in an anti-bend magnet 5 m upstream of QC1RP. With photon scattering enabled in the model, however, the absorbed photon density increases by a factor of nearly 30 to 1.1 photons/m/e+, comparable to the rates in the arcs. Figure 7 shows the distributions for photons absorbed in the walls. No photons are absorbed which

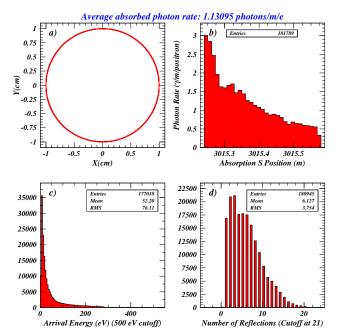


Figure 7: Distributions of absorbed photons in the finalfocus quadrupole magnet QC1RP including the effects of photon scattering.

have not previously scattered at least once, i.e. the directly incident photons strike the wall at such shallow incident angle and low energy that the absorption rate is negligible. A substantial fraction of the photons are absorbed in regions where the photoelectrons produced can be trapped in the quadrupole field. The average energy of these photons is about 50 eV. An initial study of the consequences of such a rate of absorbed photons for electron cloud buildup in the IR is presented in Ref. [11].

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