

PHYS 3323 - Formulae

$\mathbf{E}(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \int \frac{\rho(\mathbf{r}')}{\mathbf{r}^2} \hat{\mathbf{z}} d\tau'$	$\oint \mathbf{E} \cdot d\mathbf{a} = \frac{1}{\epsilon_0} Q_{enc}$ $\nabla \times \mathbf{E} = 0$	$\nabla \cdot \mathbf{E} = \frac{\rho}{\epsilon_0}$ $\mathbf{E} = -\nabla V$
$\oint \mathbf{E} \cdot d\mathbf{l} = 0$	$\nabla^2 V = -\frac{\rho}{\epsilon_0}$	$V(\mathbf{r}) = \int \frac{\rho(\mathbf{r}')}{\mathbf{r}^2} d\tau'$
$V(\mathbf{r}) = - \int_O^r \mathbf{E}(\mathbf{r}') \cdot d\mathbf{l}$		
$E_{above}^\perp - E_{below}^\perp = \frac{\sigma}{\epsilon_0}$		
At the surface of a conductor	$\mathbf{E} = \frac{\sigma}{\epsilon_0} \hat{\mathbf{n}}$	and $\sigma = -\epsilon_0 \frac{\partial V}{\partial n}$
$C = Q/V$	$W = \frac{1}{2} CV^2$	$\mathbf{F} = q\mathbf{E}$
$V_{dip} = \frac{1}{4\pi\epsilon_0} \frac{1}{r^2} \hat{\mathbf{r}} \cdot \int \mathbf{r}' \rho(\mathbf{r}') d\tau'$	$\mathbf{p} = \int \mathbf{r}' \rho(\mathbf{r}') d\tau'$	$V_{dip}(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \frac{\mathbf{p} \cdot \hat{\mathbf{r}}}{r^2}$
$\mathbf{E}_{dip} = \frac{p}{4\pi\epsilon_0 r^3} (2 \cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\theta})$		

Solution to Laplace equation in cartesian coordinates in two dimensions

$$V(x, y) = (Ae^{kx} + Be^{-kx})(C \sin ky + D \cos ky).$$

Solution to Laplace equation in spherical coordinates

$$V(r, \theta) = \sum_{l=0}^{\infty} \left(A_l r^l + \frac{B_l}{r^{l+1}} \right) P_l(\cos \theta).$$

Solution to Laplace equation in cylindrical coordinates (2 D)

$$V(\rho, \phi) = \sum_{m=1}^{\infty} (A_m \cos m\phi + B_m \sin m\phi)(C_m \rho^m + \frac{D_m}{\rho^m}) + A_0 + B_0 \ln(\rho)$$

Multipole expansion

$$V(\mathbf{r}) = \frac{1}{4\pi\epsilon_0} \sum_{n=0}^{\infty} \frac{1}{r^{n+1}} \int (r')^n P_n(\cos \alpha) \rho(\mathbf{r}') d\tau'$$

- Spherical:

- line element

$$d\mathbf{l} = dr \hat{\mathbf{r}} + rd\theta \hat{\theta} + r \sin \theta d\phi \hat{\phi}$$

- Gradient:

$$\nabla \mathbf{t} = \frac{\partial t}{\partial r} \hat{\mathbf{r}} + \frac{1}{r} \frac{\partial t}{\partial \theta} \hat{\theta} + \frac{1}{r \sin \theta} \frac{\partial t}{\partial \phi} \hat{\phi}$$

- Divergence:

$$\nabla \cdot \mathbf{v} = \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 v_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (\sin \theta v_\theta) + \frac{1}{r \sin \theta} \frac{\partial v_\phi}{\partial \phi}$$

- Cylindricall:

- Gradient:

$$\nabla \mathbf{t} = \frac{\partial t}{\partial r} \hat{\mathbf{r}} + \frac{1}{r} \frac{\partial t}{\partial \phi} \hat{\phi} + \frac{\partial t}{\partial z} \hat{\mathbf{z}}$$

- Divergence:

$$\nabla \cdot \mathbf{v} = \frac{1}{r} \frac{\partial (rv_r)}{\partial r} + \frac{1}{r} \frac{\partial v_\phi}{\partial \phi} + \frac{\partial v_z}{\partial z}$$

- Curl:

$$\nabla \times \mathbf{v} = \left(\frac{1}{r} \frac{\partial v_z}{\partial \phi} - \frac{\partial v_\phi}{\partial z} \right) \hat{\mathbf{r}} + \left(\frac{\partial v_r}{\partial z} - \frac{\partial v_z}{\partial r} \right) \hat{\phi} + \frac{1}{r} \left(\frac{\partial (rv_\phi)}{\partial r} - \frac{\partial v_r}{\partial \phi} \right) \hat{\mathbf{z}}$$

Electric Displacement

$$\begin{aligned}\rho &= \rho_b + \rho_f. \\ \epsilon_0 \nabla \cdot \mathbf{E} &= \rho = -\nabla \cdot \mathbf{P} + \rho_f. \\ \rho_f &= \nabla \cdot (\epsilon_0 \mathbf{E} + \mathbf{P}). \\ \mathbf{D} &\equiv \epsilon_0 \mathbf{E} + \mathbf{P}. \\ \nabla \cdot \mathbf{D} &= \rho_f. \\ \nabla \cdot \mathbf{P} &= -\rho_b \\ \mathbf{P} \cdot \hat{\mathbf{n}} &= \sigma_b\end{aligned}$$

Linear dielectric

$$\mathbf{P} = \epsilon_0 \chi \mathbf{E}$$

where χ is the electric susceptibility.

$$\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P} = \epsilon_0 (1 + \chi) \mathbf{P} = \epsilon \mathbf{E}.$$

Boundary conditions

$$\begin{aligned}D_{\text{above}}^\perp - D_{\text{below}}^\perp &= \sigma_f & \mathbf{D}_{\text{above}}^\parallel - \mathbf{D}_{\text{below}}^\parallel &= \mathbf{P}_{\text{above}}^\parallel - \mathbf{P}_{\text{below}}^\parallel \\ E_{\text{above}}^\perp - E_{\text{below}}^\perp &= \frac{1}{\epsilon_0} \sigma & \mathbf{E}_{\text{above}}^\parallel - \mathbf{E}_{\text{below}}^\parallel &= 0\end{aligned}$$

In linear dielectric

$$\rho_b = -\nabla \cdot \mathbf{P} = -\nabla \cdot \left(\epsilon_0 \frac{\chi}{\epsilon} \mathbf{D} \right) = -\left(\frac{\chi}{1 + \chi} \right) \rho_f.$$

the bound charge density is proportional to the free charge density.

$$\epsilon_{\text{above}} E_{\text{above}}^\perp - \epsilon_{\text{below}} E_{\text{below}}^\perp = \sigma_f.$$

Current

$$\begin{aligned}\mathbf{J} &= \sigma \mathbf{E} & V &= IR \\ \mathbf{F} &= \int (\mathbf{v} \times \mathbf{B}) dq = \int (\mathbf{v} \times \mathbf{B}) \lambda dl = \int (\mathbf{I} \times \mathbf{B}) dl. \\ \mathbf{F}_{\text{mag}} &= I \int (d\mathbf{l} \times \mathbf{B}). \\ F_{\text{mag}} &= \int (\mathbf{v} \times \mathbf{B}) \rho d\tau = \int (\mathbf{J} \times \mathbf{B}) d\tau. \\ \nabla \cdot \mathbf{J} &\equiv -\frac{\partial \rho}{\partial t}.\end{aligned}$$

Magnetic Field

Biot-Savart

$$\mathbf{B} = \frac{\mu_0}{4\pi} \int \frac{\mathbf{I} \times \hat{\mathbf{z}}}{r^2} dl' = \frac{\mu_0}{4\pi} I \int \frac{dl' \times \hat{\mathbf{z}}}{r^2}.$$

$$\mathbf{B}(\mathbf{r}) = \frac{\mu_0}{4\pi} \int \frac{\mathbf{J}(\mathbf{r}') \times \hat{\mathbf{z}}}{r^2} d\tau'.$$

Divergence

$$\nabla \cdot \mathbf{B} = \frac{\mu_0}{4\pi} \int \nabla \cdot \left(\frac{\mathbf{J}(\mathbf{r}') \times \hat{\mathbf{z}}}{r^2} \right) d\tau'.$$

$$\nabla \cdot \mathbf{B} = 0.$$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}$$

$$\oint \mathbf{B} \cdot d\mathbf{l} = \mu_0 I_{enc} + \mu_0 \epsilon_0 \frac{\partial \Phi_e}{\partial t}.$$

Vector potential

$$\mathbf{B} = \nabla \times \mathbf{A}.$$

$$\nabla \times \mathbf{B} = \nabla \times (\nabla \times \mathbf{A}) = \nabla(\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A} = \mu_0 \mathbf{J}$$

$$\mathbf{A}(\mathbf{r}) = \frac{\mu_0}{4\pi} \int \frac{\mathbf{J}(\mathbf{r}')}{r} d\tau'.$$

$$\mathbf{A} = \frac{\mu_0 I}{4\pi} \int \frac{1}{r} dl'.$$

Boundary Conditions

In linear media

$$\epsilon_1 E_1^\perp - \epsilon_2 E_2^\perp = \sigma_f, \quad \mathbf{E}_1^\parallel - \mathbf{E}_2^\parallel = 0$$

$$B_1^\perp - B_2^\perp = 0, \quad \frac{1}{\mu_1} \mathbf{B}_1^\parallel - \frac{1}{\mu_2} \mathbf{B}_2^\parallel = \mathbf{K}_f \times \hat{\mathbf{n}}$$

If there is no free charge or free current at the interface

$$\epsilon_1 E_1^\perp - \epsilon_2 E_2^\perp = 0, \quad \mathbf{E}_1^\parallel - \mathbf{E}_2^\parallel = 0$$

$$B_1^\perp - B_2^\perp = 0, \quad \frac{1}{\mu_1} \mathbf{B}_1^\parallel - \frac{1}{\mu_2} \mathbf{B}_2^\parallel = 0$$

The vector potential is continuous at a boundary.

Multipole Expansion

$$\begin{aligned}\mathbf{A}(\mathbf{r}) &= \frac{\mu_0 I}{4\pi} \int \frac{1}{r} d\mathbf{l}' = \frac{\mu_0 I}{4\pi} \sum_{n=0}^{\infty} \frac{1}{r^{n+1}} \oint (r')^n P_n(\cos \alpha) d\mathbf{l}'. \\ \mathbf{m} &= I \int d\mathbf{a}' \\ \mathbf{A}_{\text{dip}}(\mathbf{r}) &= \frac{\mu_0}{4\pi} \frac{\mathbf{m} \times \hat{\mathbf{r}}}{r^2}. \\ \mathbf{A}_{\text{dip}}(\mathbf{r}) &= \frac{\mu_0}{4\pi} \frac{m \sin \theta}{r^2} \hat{\phi}. \\ \mathbf{B} &= \nabla \times \mathbf{A} = \frac{\mu_0 m}{4\pi r^3} (2 \cos \theta \hat{\mathbf{r}} + \sin \theta \hat{\theta}).\end{aligned}$$

Torque and Force

$$\mathbf{N} = \mathbf{m} \times \mathbf{B}.$$

$$\mathbf{F} = \nabla(\mathbf{m} \cdot \mathbf{B})$$

Magnetization

The total current

$$\mathbf{J} = \mathbf{J}_b + \mathbf{J}_f$$

$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J} = \mu_0 (\nabla \times \mathbf{M} + \mathbf{J}_f)$$

$$\nabla \times \left(\frac{1}{\mu_0} \mathbf{B} - \mathbf{M} \right) = \mathbf{J}_f$$

$$\mathbf{H} = \frac{1}{\mu_0} \mathbf{B} - \mathbf{M}$$

$$\nabla \times \mathbf{H} = \mathbf{J}_f + \frac{\partial \mathbf{D}}{\partial t}$$

$$\oint \mathbf{H} \cdot d\mathbf{l} = I_{enc}^{free} + \frac{d}{dt} \int_S \mathbf{D} \cdot d\mathbf{a}$$

$$\mathbf{J}_b = \nabla \times \mathbf{M} \quad \mathbf{K}_b = \mathbf{M} \times \hat{\mathbf{n}}$$

Linear media:

$$\mathbf{M} = \chi_m \mathbf{H}$$

$$\chi_m = \frac{\mu}{\mu_0} - 1$$

$$\mathbf{B} = \mu \mathbf{H}$$

Boundary conditions:

$$\begin{aligned}B_{\text{above}}^\perp &= B_{\text{below}}^\perp & B_{\text{above}}^\parallel - B_{\text{below}}^\parallel &= \mu_0 K \\ \mathbf{A}_{\text{above}} &= \mathbf{A}_{\text{below}} & \mathbf{B}_{\text{above}} - \mathbf{B}_{\text{below}} &= \mu_0 (\mathbf{K} \times \hat{\mathbf{n}})\end{aligned}$$

Induced fields

$$\begin{aligned}\mathcal{E} &= \oint \mathbf{E} \cdot d\mathbf{l} = -\frac{d\Phi}{dt} \\ \oint \mathbf{E} \cdot d\mathbf{l} &= -\frac{d}{dt} \int \mathbf{B} \cdot d\mathbf{a} \\ \nabla \times \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t}. \\ \mathbf{E}(\mathbf{r}) &= \frac{1}{4\pi} \int \frac{-\frac{\partial \mathbf{B}}{\partial t} \times \hat{\mathbf{r}}}{r^2} d\tau' = -\frac{1}{4\pi} \frac{\partial}{\partial t} \int \frac{\mathbf{B} \times \hat{\mathbf{r}}}{r^2} d\tau \\ \int \mathbf{B} \cdot d\mathbf{l} &= \mu_0 \int \mathbf{J} \cdot d\mathbf{l} = \mu_0 I_{enc}\end{aligned}$$

Inductance

$$\begin{aligned}\Phi_2 &= M_{21}I_1 & M_{21} &= \frac{\mu_0}{4\pi} \oint \oint \frac{d\mathbf{l}_1 \cdot d\mathbf{l}_2}{r} \\ M_{21} &= M_{12}, & \mathcal{E}_2 &= -\frac{d\Phi_2}{dt} = -M \frac{dI_1}{dt} \\ \Phi &= LI, & \mathcal{E} &= -L \frac{dI}{dt}\end{aligned}$$

Maxwell Equations

$$\begin{aligned}\nabla \cdot \mathbf{E} &= \frac{1}{\epsilon_0} \rho & \nabla \cdot \mathbf{B} &= 0 \\ \nabla \times \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t} & \nabla \times \mathbf{B} &= \mu_0 \mathbf{J} + \mu_0 \epsilon_0 \frac{\partial \mathbf{E}}{\partial t}\end{aligned}$$

In matter

$$\begin{aligned}\nabla \cdot \mathbf{D} &= \rho_f & \nabla \cdot \mathbf{B} &= 0 \\ \nabla \times \mathbf{E} &= -\frac{\partial \mathbf{B}}{\partial t} & \nabla \times \mathbf{H} &= \mathbf{J}_f + \frac{\partial \mathbf{D}}{\partial t}\end{aligned}$$

Wave equation

$$\begin{aligned}\nabla^2 \mathbf{E} &= \mu_0 \epsilon_0 \frac{\partial^2 \mathbf{E}}{\partial t^2}, & \nabla^2 \mathbf{B} &= \mu_0 \epsilon_0 \frac{\partial^2 \mathbf{B}}{\partial t^2} \\ v &= \frac{1}{\sqrt{\epsilon \mu}}\end{aligned}$$

Waves and Waveguides

Phase velocity and group velocity (k is the propagation direction)

$$v_p = \frac{\omega}{k}, \quad v_g = \frac{d\omega}{dk}$$

Rectangular waveguide

$$k = \sqrt{(\omega/c)^2 - \pi^2[(m/a)^2 + (n/b)^2]}, \quad \omega_{mn} = c\pi\sqrt{(m/a)^2 + (n/b)^2}, \quad k = \frac{1}{c}\sqrt{\omega^2 - \omega_{mn}^2}$$

Energy and Momentum of EM fields

$$u = \frac{1}{2} \left(\epsilon_0 E^2 + \frac{1}{\mu_0} B^2 \right), \quad \mathbf{S} = \frac{1}{\mu_0} (\mathbf{E} \times \mathbf{B})$$

Momentum

$$\mathbf{g} = \epsilon_0 (\mathbf{E} \times \mathbf{B}), \quad \mathbf{p} = \int_V \mathbf{g} d\tau$$

Angular momentum

$$\mathbf{l} = \mathbf{r} \times \mathbf{g} = \epsilon_0 [\mathbf{r} \times (\mathbf{E} \times \mathbf{B})]$$

Legendre Polynomials

$$\begin{aligned} P_0(\cos \theta) &= 1 \\ P_1(\cos \theta) &= \cos \theta \\ P_2(\cos \theta) &= (3 \cos^2 \theta - 1)/2 \\ P_3(\cos \theta) &= (5 \cos^3 \theta - 3 \cos \theta)/2 \\ P_4(\cos \theta) &= (35 \cos^4 \theta - 30 \cos^2 \theta + 3)/8 \\ P_5(\cos \theta) &= (63 \cos^5 \theta - 70 \cos^3 \theta + 15 \cos \theta)/8 \end{aligned}$$

$$\begin{aligned} \int_{-1}^1 P_l(x) P_{l'}(x) dx &= \int_0^\pi P_l(\cos \theta) P_{l'}(\cos \theta) \sin \theta d\theta = \frac{2}{2l+1} \delta_{ll'} \\ \int_0^a \sin(n\pi y/a) \sin(n'\pi y/a) dy &= \frac{1}{2} \delta_{nn'} \end{aligned}$$