

### 1. SHO propagator

We know that given the eigenfunctions and the eigenvalues we can construct the propagator:

$$K(x, t; x', t') = \sum_n \psi_n(x) \psi_n^*(x') e^{-iE_n(t-t')/\hbar} \quad (1)$$

Consider the reverse process (since the path integral approach gives  $K$  directly), for the case of the harmonic oscillator where

$$K(x, t; x', 0) = \sqrt{\frac{m\omega}{2\pi i\hbar \sin \omega t}} \exp \left\{ \frac{im\omega}{2\hbar \sin \omega t} [(x^2 + x'^2) \cos \omega t - 2xx'] \right\}$$

- (a) Set  $x = x' = t' = 0$ . By expanding both sides of Equation 1 you should find that  $E = \hbar\omega/2, 5\hbar\omega/2, 9\hbar\omega/2, \dots$ , etc. What happened to the levels in between?
- (b) Now consider the extraction of the eigenfunctions. Let  $x = x'$  and  $t' = 0$ . Find  $E_0, E_1, |\psi_0(x)|^2$ , and  $|\psi_1(x)|^2$  by expanding in powers of  $a = \exp(i\omega t)$ .

### 2. Sakurai, p. 149, problem 30

The propagator in position space is (2.5.26), namely,

$$K(\mathbf{x}'', t, \mathbf{x}', t_0) = \sum_{a'} \left\langle \mathbf{x}'' \mid \exp\left(\frac{-iHt}{\hbar}\right) \mid a' \right\rangle \left\langle a' \mid \exp\left(\frac{iHt_0}{\hbar}\right) \mid \mathbf{x}' \right\rangle = \langle \mathbf{x}'', t \mid \mathbf{x}', t_0 \rangle$$

The analogous propagator in momentum space is given by  $\langle \mathbf{p}'', t \mid \mathbf{p}', t_0 \rangle$ . Derive an explicit expression for  $\langle \mathbf{p}'', t \mid \mathbf{p}', t_0 \rangle$  for the free particle case.

### 3. Green's function

Derive the Green's function  $G(x, x', E)$  for a free particle in one dimension, where

$$\left( \frac{-\hbar^2}{2m} \frac{\partial^2}{\partial x^2} - E \right) G(x, x', E) = \delta(x - x')$$

### 4. Sakurai, p. 149, Problem 26

Consider a particle moving in one dimension under the influence of a potential  $V(x)$ . Suppose its wave function can be written as  $\exp[iS(x, t)/\hbar]$ . Prove that  $S(x, t)$  satisfies the classical Hamilton-Jacobi equation to the extent that  $\hbar$  can be regarded as small in some sense. Show how one may obtain the correct wave function for a plane wave by starting with the solution of the classical Hamilton-Jacobi equation with  $V(x)$  set equal to zero. Why do we get the exact wave function in this particular case?

### 5. Coherent state

A coherent state of a one-dimensional simple harmonic oscillator is defined to be an eigenstate of the (non-Hermitian) annihilation operator  $a$ :

$$a|\lambda\rangle = \lambda|\lambda\rangle,$$

where  $\lambda$  is, in general, a complex number.

- (a) Prove that

$$|\lambda\rangle = e^{-|\lambda|^2/2} e^{\lambda a^\dagger} |0\rangle$$

is a normalized coherent state.

- (b) Prove the minimum uncertainty relation for such a state.  
 (c) Write  $|\lambda\rangle$  as

$$|\lambda\rangle = \sum_{n=0}^{\infty} f(n) |n\rangle.$$

Show that the distribution of  $|f(n)|^2$  with respect to  $n$  is of the Poisson form. Find the most probable value of  $n$ , hence of  $E$ .

- (d) Show that a coherent state can also be obtained by applying the translation (finite-displacement) operator  $e^{-ipl/\hbar}$  (where  $p$  is the momentum operator, and  $l$  is the displacement distance) to the ground state.  
 (e) Show that the wave function of the coherent state is

$$\psi_\lambda(x) = \langle x | \lambda \rangle = \left(\frac{m\omega}{\pi\hbar}\right)^{1/4} e^{-\lambda^2/2} e^{-(m\omega/2\hbar)x^2} e^{\sqrt{(2m\omega/\hbar)}\lambda x}$$

Start by using  $a|\lambda\rangle = \lambda|\lambda\rangle$  in the coordinate representation. Fix the normalization by demanding that  $\langle\lambda'|\lambda\rangle = e^{\lambda'^*\lambda}$ . (Hint: The identity  $e^A e^B = e^B e^A e^{[A,B]}$  which is true if  $[A, B]$  commutes with  $A$  and  $B$  might be useful.) Show that  $\psi_\lambda(x, t)$  evolves with the time like classical coordinates given that  $|\lambda\rangle \rightarrow |\lambda e^{-i\omega t}\rangle$ .

### 6. Sakurai, p. 150, problem 35

Consider the Hamiltonian of a spinless particle of charge  $q$ . In the presence of a static magnetic field, the interaction terms can be generated by

$$\mathbf{P}_{\text{operator}} \rightarrow \mathbf{P}_{\text{operator}} - \frac{q\mathbf{A}}{c},$$

where  $\mathbf{A}$  is the appropriate vector potential. Suppose, for simplicity, that the magnetic field  $\mathbf{B}$  is uniform in the positive  $z$ -direction. Prove that the above prescription indeed leads to the correct expression for the interaction of the orbital magnetic moment  $(q/2mc)\mathbf{L}$  with the magnetic field  $\mathbf{B}$ . Show that there is also an extra term proportional to  $B^2(x^2 + y^2)$ , and comment briefly on its physical significance.

## 7. Lorentz force law

Show that if we modify the classical Lagrangian to include the interaction of a charged particle with a magnetic field so that

$$L \rightarrow L + \frac{q}{c} \mathbf{v} \cdot \mathbf{A}(x)$$

that Lagrange's equations yield the Lorentz force law

$$m \frac{d^2 \mathbf{x}}{dt^2} = \frac{q}{c} \mathbf{v} \times \mathbf{B}$$