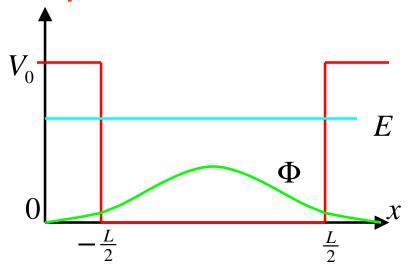
Stationary states in the square well

02/23/2005

$$\alpha = \sqrt{\frac{2m}{\hbar^2}(V_0 - E)}$$

$$k = \sqrt{\frac{2m}{\hbar^2}E}, \quad k_V = \sqrt{\frac{2m}{\hbar^2}V_0}$$



$$z = \frac{L}{2} : kAe^{-\alpha \frac{L}{2}} = kB \sin(k \frac{L}{2} + \varphi)$$

$$-\alpha Ae^{-\alpha \frac{L}{2}} = kB \cos(k \frac{L}{2} + \varphi)$$

$$\Rightarrow \begin{cases} A^{2}e^{-\alpha L}(k^{2} + \alpha^{2}) = k^{2}B^{2} \\ \frac{A^{2}}{B^{2}}e^{-\alpha L} = \frac{E}{V_{0}} = \frac{k^{2}}{k_{v}^{2}} \end{cases}$$

$$\Rightarrow \begin{cases} \frac{A^{2}}{B^{2}}e^{-\alpha L} = \frac{E}{V_{0}} = \frac{k^{2}}{k_{v}^{2}} \end{cases}$$

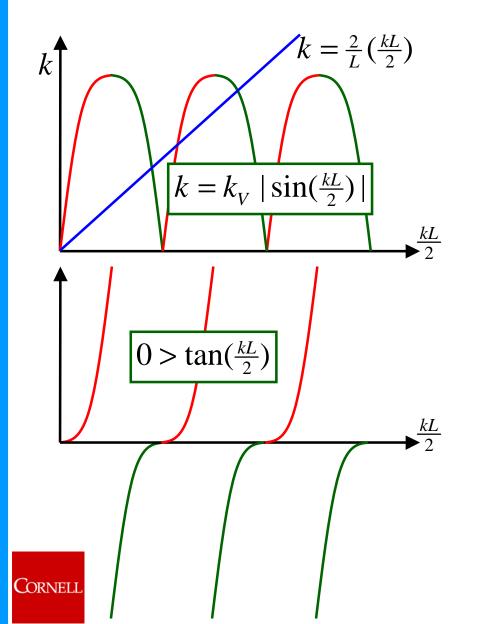
$$\Rightarrow \begin{cases} \frac{k}{k_{v}} = |\sin(\frac{kL}{2} + \varphi)| \\ \Rightarrow \tan(\frac{kL}{2} + \varphi) \end{cases}$$

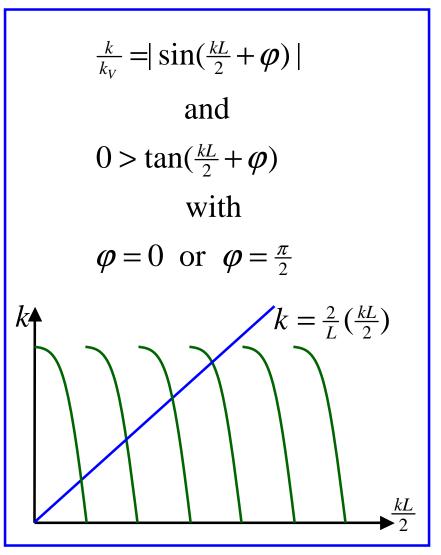
$$\Rightarrow \begin{cases} \frac{k}{k_{v}} = |\sin(\frac{kL}{2} + \varphi)| \\ \Rightarrow \tan(\frac{kL}{2} + \varphi) \end{cases}$$

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$$\left| \frac{k}{k_V} = -\operatorname{sign}[\cos(\frac{k}{k_V} \frac{k_V L}{2} + \boldsymbol{\varphi})] \sin(\frac{k}{k_V} \frac{k_V L}{2} + \boldsymbol{\varphi}) \right|$$

Graphical determination of possible wave numbers 02/23/2005



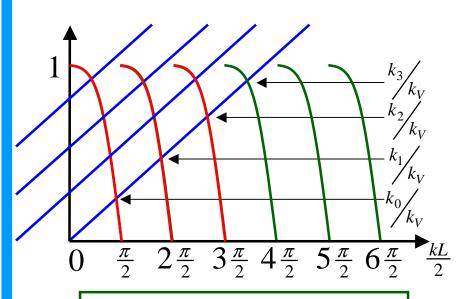


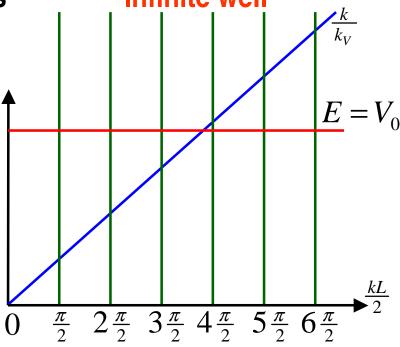
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Square well









Stationary states for the **finite** well:

$$\frac{k}{k_V} = \cos(\text{mod}[\frac{kL}{2}, \frac{\pi}{2}])$$

Stationary states for the **infinite** well:

$$k = n\frac{\pi}{L}, \mod[\frac{kL}{2}, \frac{\pi}{2}] = 0$$

The finite potential well with height V_0 has one stationary state more then the infinite potential well has states with energies below V_0 .

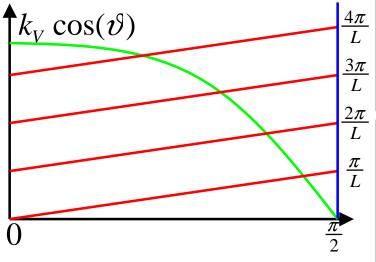


- Number of solutions = $\inf[\frac{k_V L}{\pi}] + 1 = \inf[\frac{L}{\pi} \sqrt{\frac{2m}{\hbar^2} V_0}] + 1$
- There is at least one bound state, and it is symmetric.

Graphic determination of states for the square well

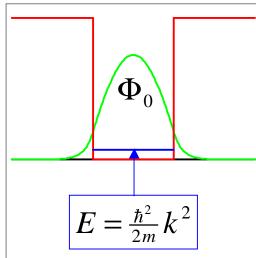
Finite well:

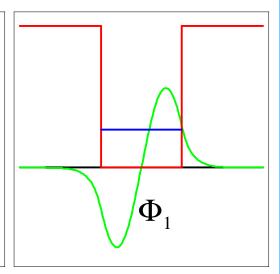
$$k = k_V \cos(\text{mod}[\frac{kL}{2}, \frac{\pi}{2}])$$

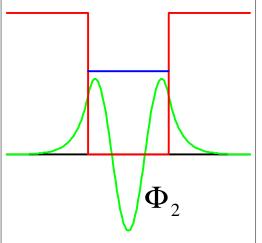


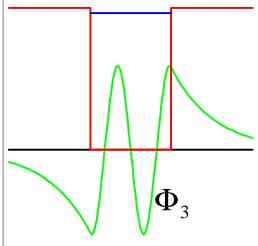
Infinite well:

$$k = n \frac{\pi}{L}$$









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