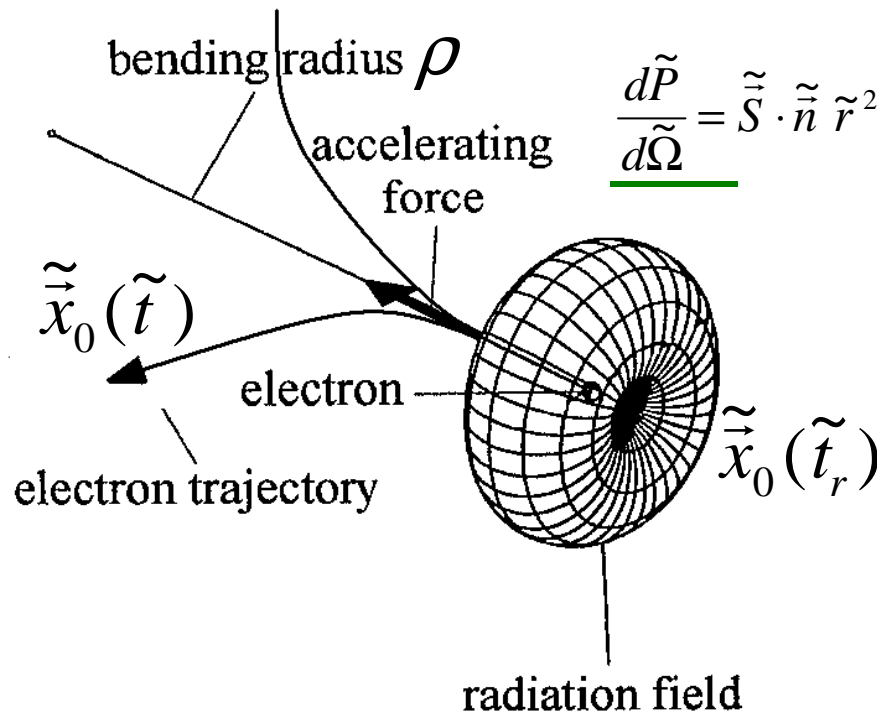




# Rest-Frame-Radiation Pattern



CHESS & LEPP



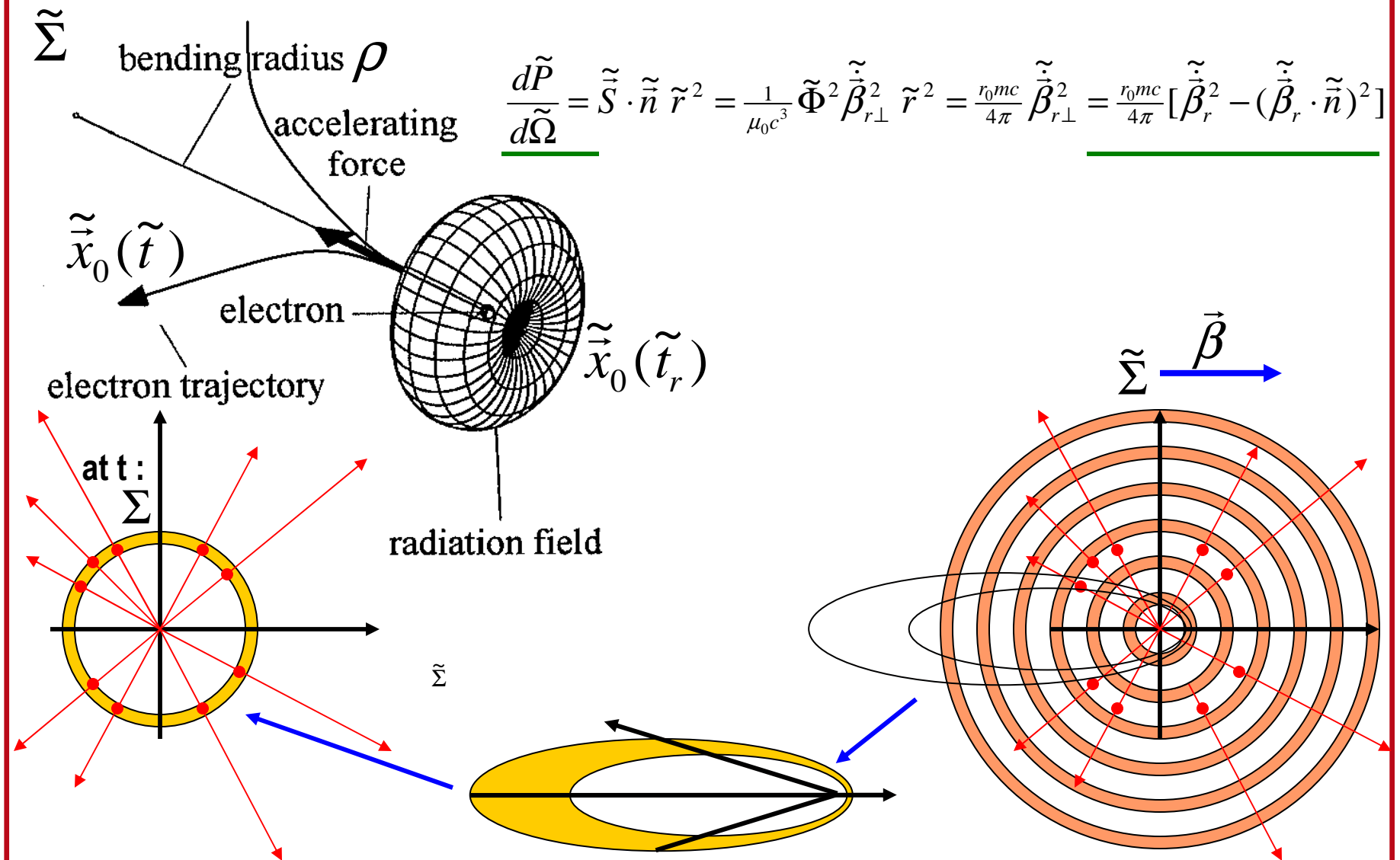
$$\frac{d\tilde{P}}{d\tilde{\Omega}} = \tilde{S} \cdot \tilde{n} \tilde{r}^2 = \frac{1}{\mu_0 c^3} \tilde{\Phi}^2 \tilde{\beta}_{r\perp}^2 \tilde{r}^2 = \frac{r_0 mc}{4\pi} \tilde{\beta}_{r\perp}^2 = \frac{r_0 mc}{4\pi} [\tilde{\beta}_r^2 - (\tilde{\beta} \cdot \tilde{n})^2]$$



# Rest-Frame-Radiation Pattern



CHESS & LEPP

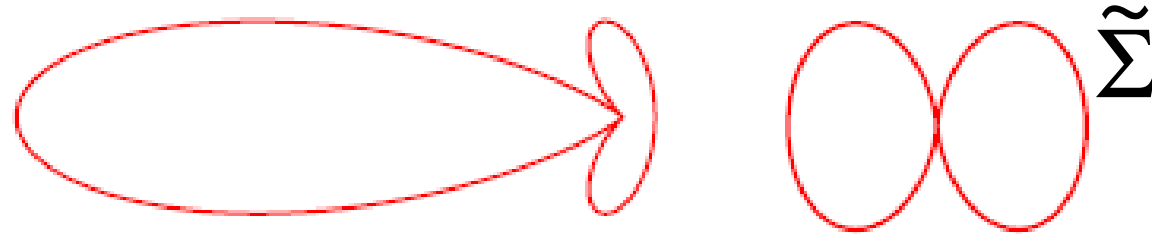




# Lab-Frame-Radiation Pattern



CHESS & LEPP

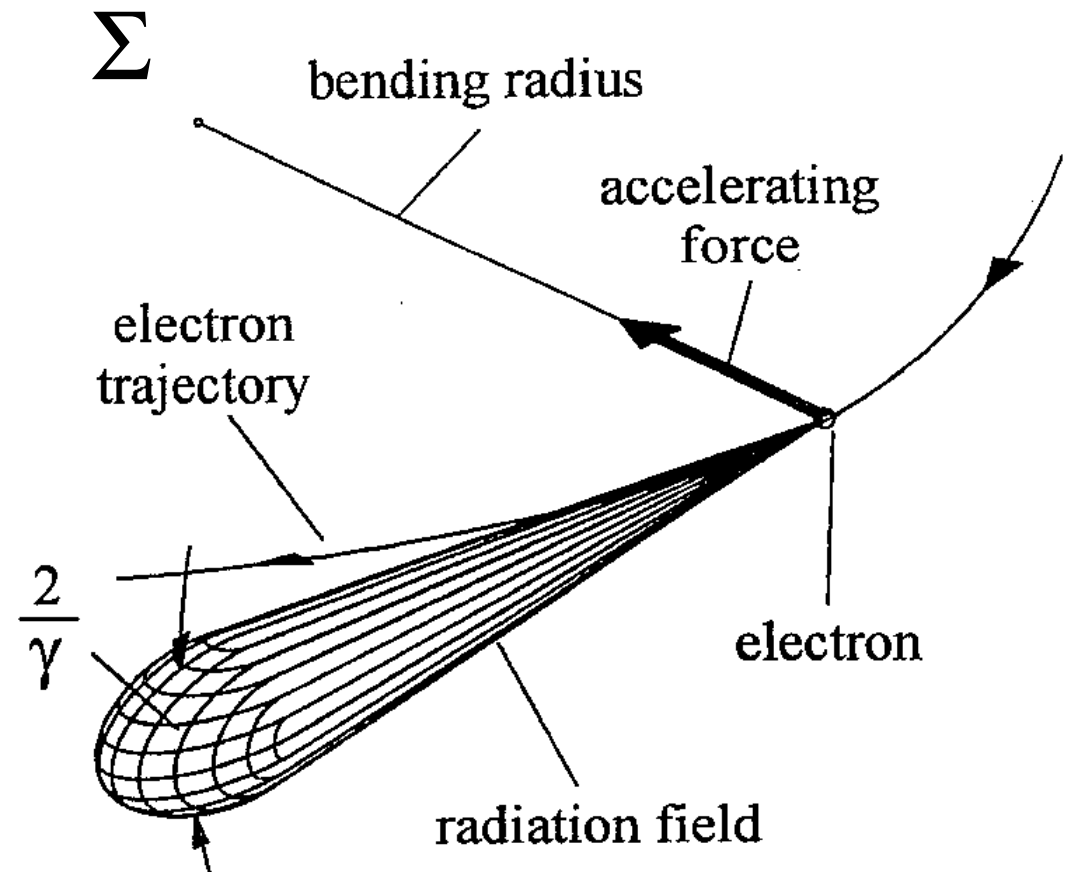


$$r^2 d\Omega = \tilde{r}^2 d\tilde{\Omega} \Rightarrow \frac{dP}{d\Omega} = \frac{d\tilde{P}}{d\tilde{\Omega}} \frac{r^2}{\tilde{r}^2}$$

$$\begin{pmatrix} ct \\ \vec{x} \end{pmatrix} = \begin{pmatrix} r \\ r\vec{n} \end{pmatrix} \Rightarrow \tilde{r} = \gamma(r - \vec{\beta} \cdot \vec{r})$$

$$\Rightarrow \frac{r^2}{\tilde{r}^2} = \frac{1}{\gamma^2(1 - \beta \cos \vartheta)^2}$$

$$\frac{dP}{d\Omega} = \frac{r_0 mc}{4\pi} \frac{\tilde{\beta}_r^2 - (\tilde{\beta}_r \cdot \tilde{\mathbf{n}})^2}{\gamma^2(1 - \beta \cos \vartheta)^2}$$



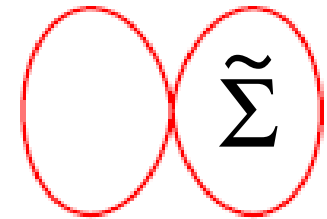
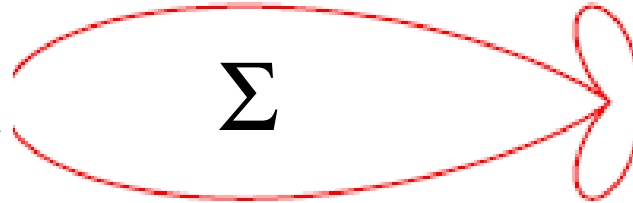


# Lab-Frame-Radiation Pattern



CHESS & LEPP

$$\begin{pmatrix} ct \\ \vec{x} \end{pmatrix} = \begin{pmatrix} r \\ r \vec{n} \end{pmatrix} \Rightarrow \begin{aligned} \tilde{r} &= \gamma(r - \vec{\beta} \cdot \vec{r}) \\ \tilde{r} \tilde{\vec{n}} &= \gamma(r \vec{n}_{\parallel} - \vec{\beta} r) + r \vec{n}_{\perp} \end{aligned}$$



$$\begin{pmatrix} \frac{1}{c} \dot{\vec{E}} \\ \dot{\vec{p}} \end{pmatrix} = \begin{pmatrix} 0 \\ mc \gamma^2 \dot{\vec{\beta}} \end{pmatrix} \Rightarrow \begin{aligned} 0 &= 0 \\ \dot{\vec{\beta}} &= \gamma^2 \dot{\vec{\beta}} \end{aligned}$$

$$\frac{dP}{d\Omega} = \frac{r_0 mc}{4\pi} \gamma^2 \frac{\dot{\vec{\beta}}_r^2 - (\dot{\vec{\beta}}_r \cdot \frac{r}{r} \vec{n}_{\perp})^2}{(1 - \beta \cos \vartheta)^2}$$

$$= \frac{r_0 mc}{4\pi} \gamma^2 \dot{\vec{\beta}}_r^2 \frac{(1 - \beta \cos \vartheta)^2 - \left(\frac{\sin \vartheta \cos \varphi}{\gamma}\right)^2}{(1 - \beta \cos \vartheta)^4} \approx \frac{r_0 mc}{4\pi} \gamma^2 \dot{\vec{\beta}}_r^2 \frac{\left(\frac{1}{2\gamma^2} + \frac{\vartheta^2}{2}\right)^2 - \frac{\vartheta^2}{\gamma^2} \cos^2 \varphi}{\left(\frac{1}{2\gamma^2} + \frac{\vartheta^2}{2}\right)^4}$$

$$= \begin{cases} \frac{3\pi}{2} \gamma^2 P \frac{(1 - \xi^2)^2}{(1 + \xi^2)^4} & \text{for horizontal} \\ \frac{3\pi}{2} \gamma^2 P \frac{(1 + \xi^2)^2}{(1 + \xi^2)^4} & \text{for vertical} \end{cases} \quad [\xi = \vartheta \gamma]$$

Conclusion: Angular power density falls off with  $\vartheta^{-4}$

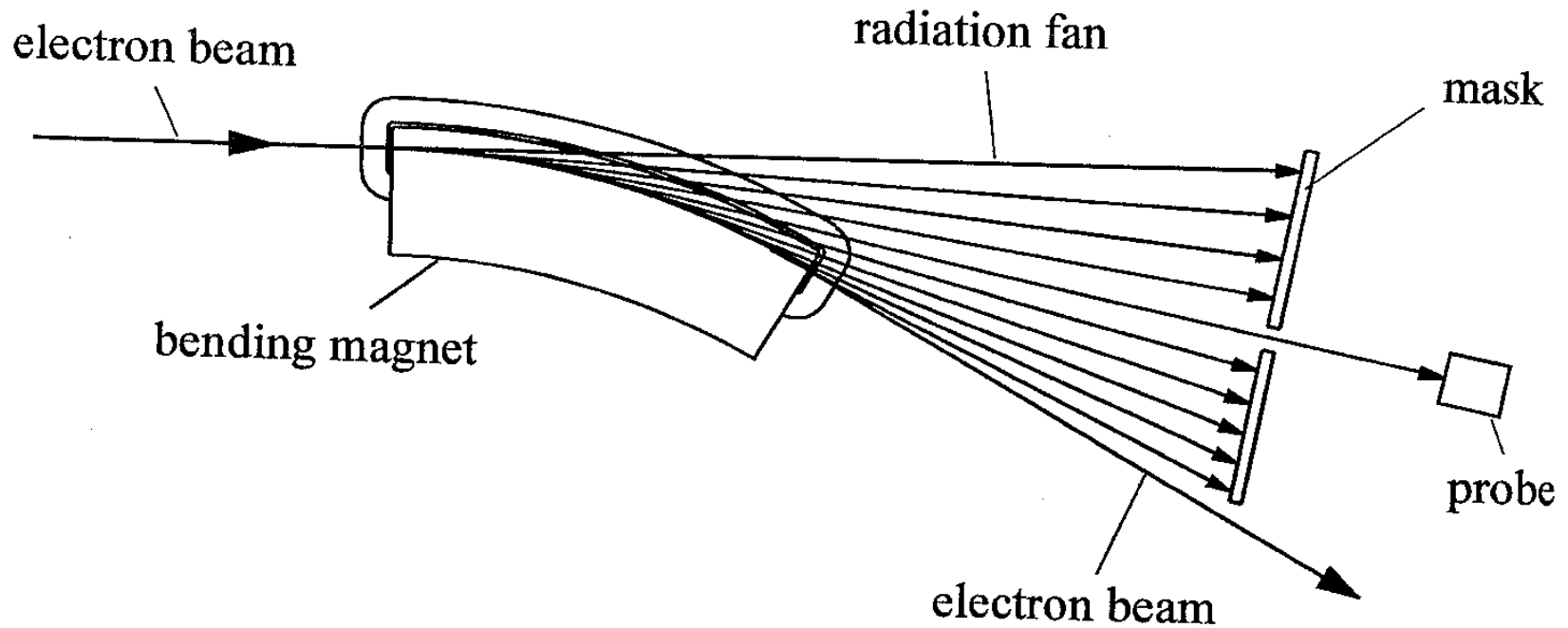
Peak angular power density growth with  $\gamma^6 / \rho^2$



## Long and Short Magnet Radiation



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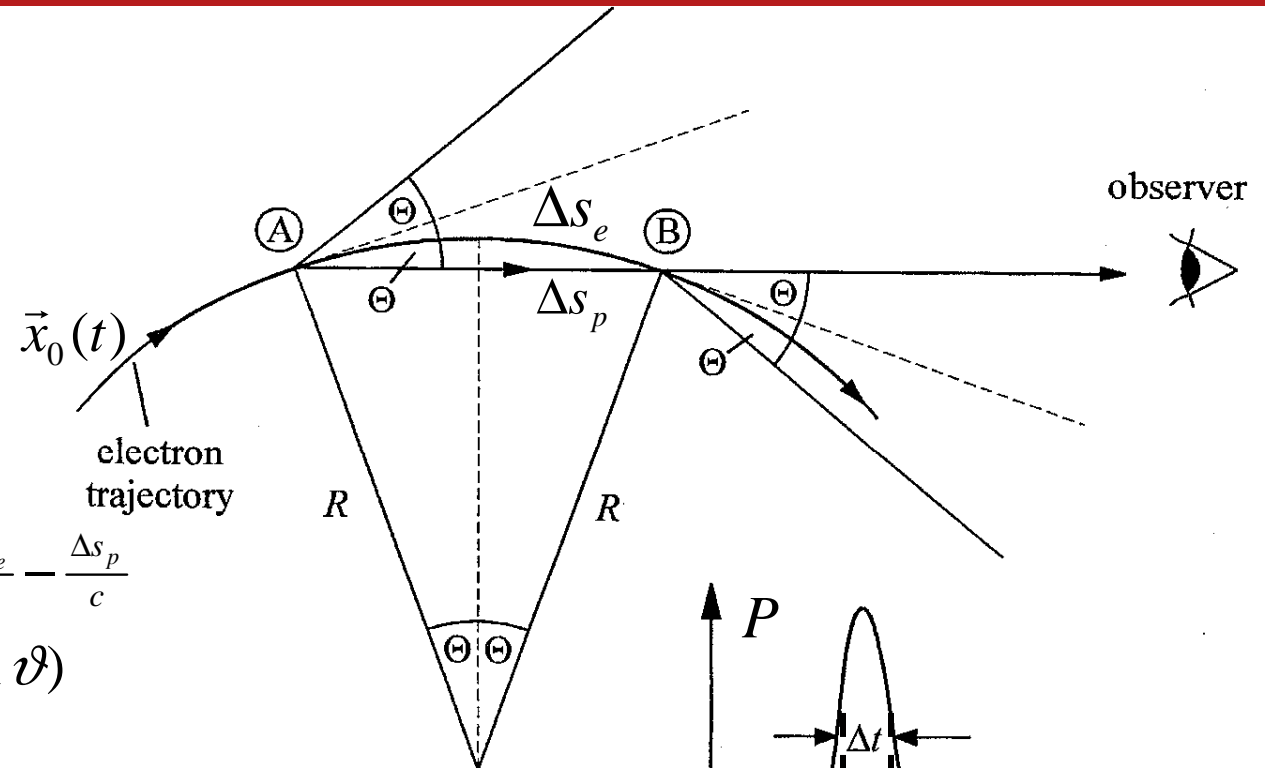
The exposure time at the probe is either given by the opening angle of radiation, or by the bend angle of the magnet, whichever is smaller.



# Spectrum of bend radiation



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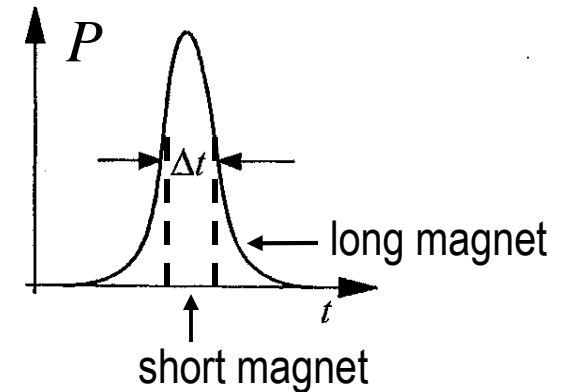


$$\Delta t = t_{electron} - t_{photon} = \frac{\Delta s_e}{v} - \frac{\Delta s_p}{c}$$

$$= \frac{1}{\beta c} (2\rho\vartheta - 2\beta\rho \sin \vartheta)$$

$$\approx \frac{2\rho}{c} \frac{1-\beta}{\beta} \vartheta \approx \frac{\rho\vartheta}{c\gamma^2}$$

$$\Delta\omega \approx \frac{1}{\Delta t} \approx \frac{c\gamma^2}{\rho\vartheta} \approx \begin{cases} \frac{c\gamma^3}{\rho} & \text{for a long magnet} \\ \frac{c\gamma^2}{L} & \text{for a short magnet} \end{cases}$$



Characteristic frequency:  $\omega_c \approx \frac{3}{2} \frac{c\gamma^3}{\rho}$



# Radiation Formulas to Remember



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$$\text{Characteristic frequency: } \omega_c \approx \frac{3}{2} \frac{c\gamma^3}{\rho}$$

Maximum spectral radiation power  $\frac{dP}{d\Omega}$  at  $\omega = \frac{1}{3} \omega_c$

Half radiation power above and half below  $\omega = \omega_c$

For a current  $I$  of electrons with:  $P = \frac{2r_e c}{3} mc^2 \frac{(\beta\gamma)^4}{\rho^2}$  in  $\Delta\psi = \gamma^{-1}$

Linear power density along a vacuum pipe (needs to be cooled):  $\frac{dP}{dl} = \frac{I}{q} \frac{2r_e}{3} mc^2 \frac{(\beta\gamma)^4}{\rho^2}$

Power density along a vacuum pipe of radius  $R$  (needs to be cooled):  $\frac{dP}{dl} = \frac{I}{q} \frac{2r_e}{3D} m\gamma c^2 \frac{(\beta\gamma)^4}{\rho^2}$

Linear photon flux density (causes gas desorption):  $\frac{d\dot{N}_p}{dl} \approx \frac{2r_e I}{3q} \frac{mc^2}{\hbar\omega_c} \frac{\gamma^4}{\rho^2} = \frac{4}{9} \frac{I}{q} \frac{r_e mc}{\hbar} \frac{\gamma}{\rho} = \frac{4}{9} \frac{I}{q} \frac{\alpha}{\rho \frac{1}{\gamma}}$

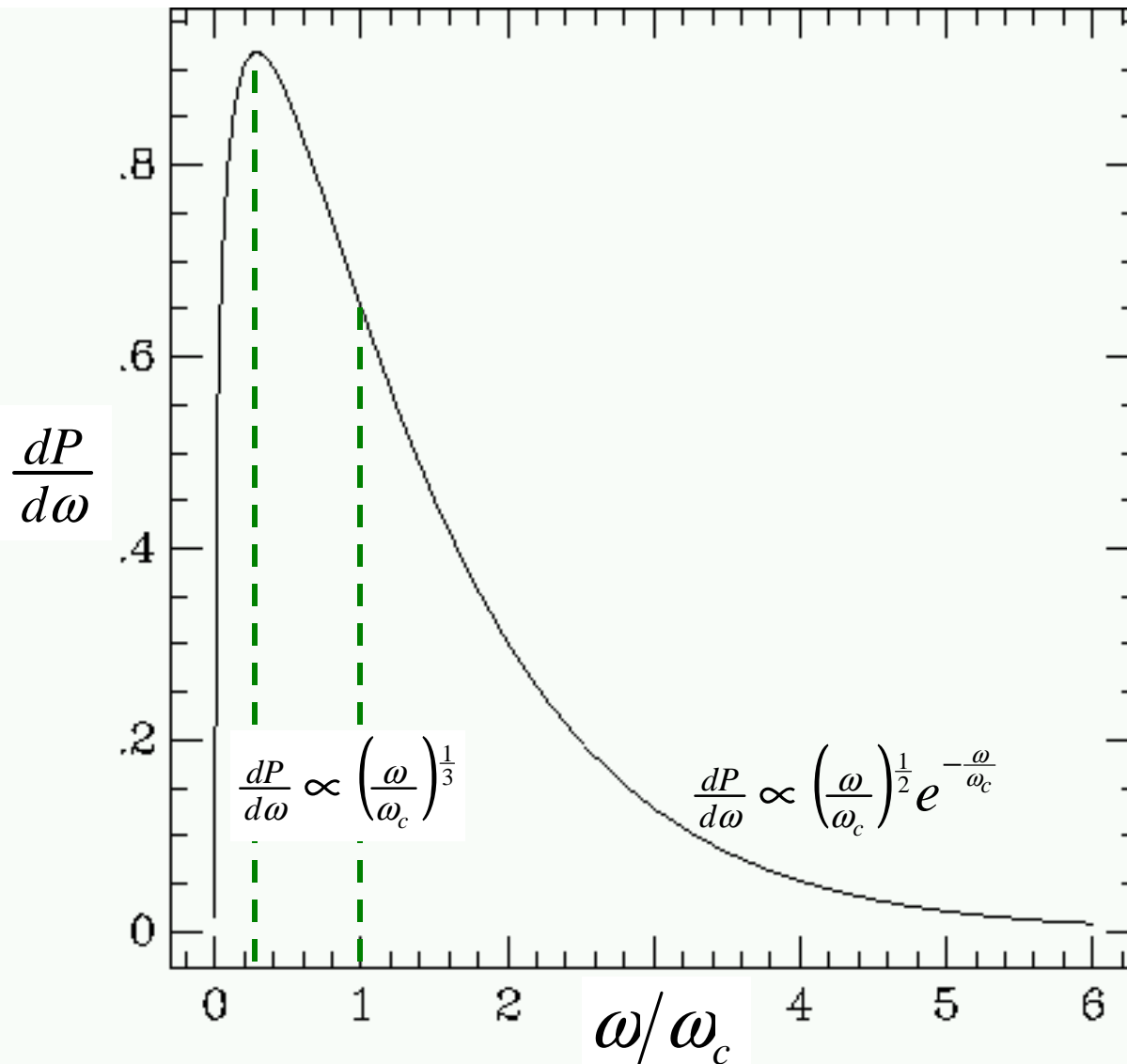
**Note:** Each electron radiates only  $\alpha=1/137$  photons per opening angle !



# Synchrotron-Radiation Spectrum



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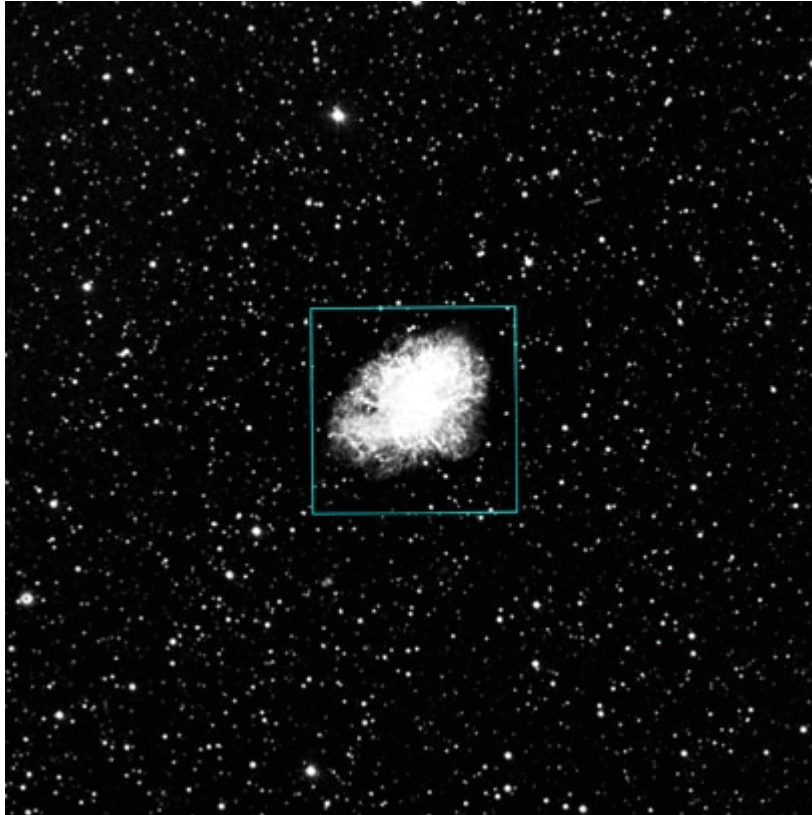


# The First Observed Synchrotron

with images from <http://www.astro.utu.fi/~cflynn/astroll/l4.html>



CHESS & LEPP



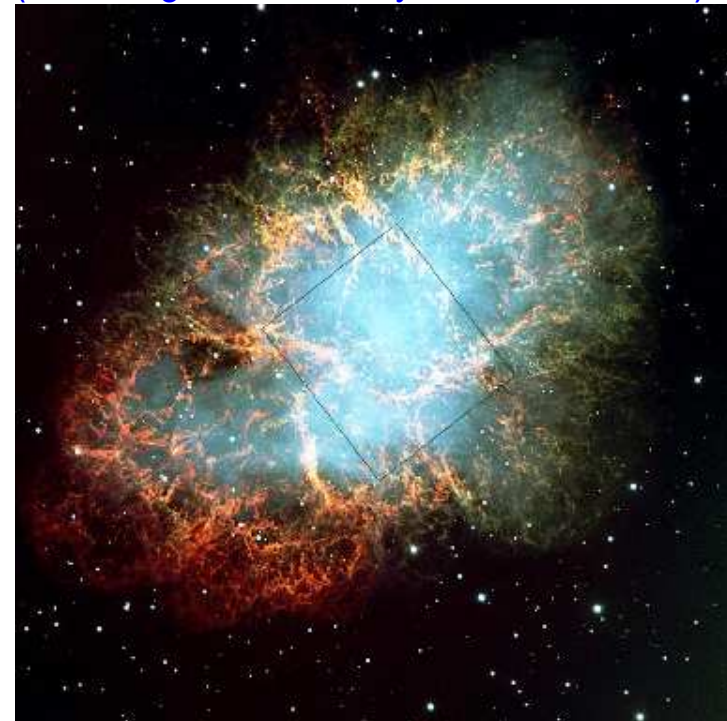
Emission from the filaments is mostly in Hydrogen and Oxygen lines, although the total optical emission of the remnant has a broad spectrum: Where is it from?

## The Crab nebula:

Remnant of the first observed supernova, 4<sup>th</sup> of July, China 1054

Visible by eye during daytime for a month, during night time for 2 years.

( The largest 4<sup>th</sup> of July fireworks, ever ! )

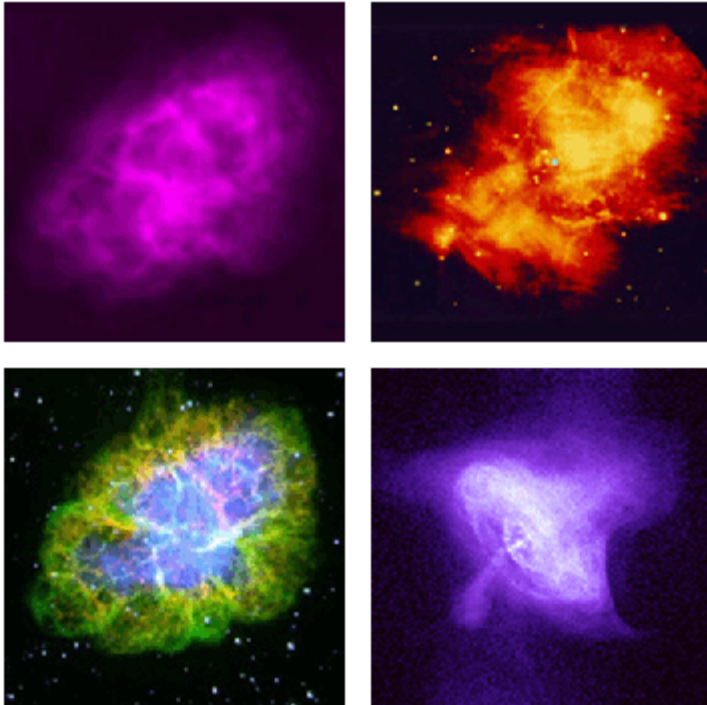




# Synchrotron Radiation in Space



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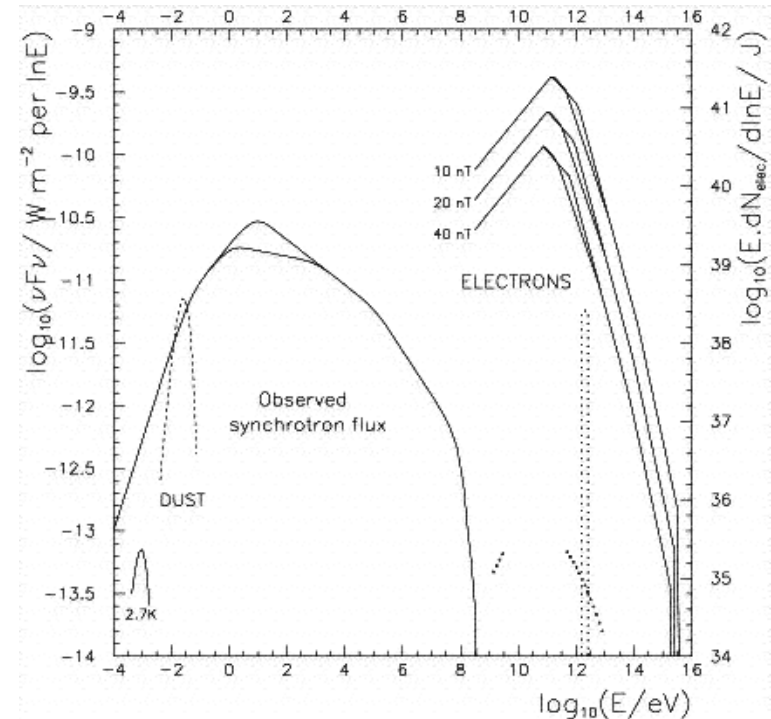
## The Crab nebula:

The spectrum has a  $\nu^{-2}$  law in the radio region.

The radiation is to about 7% polarized.

This was an early hint for synchrotron radiation.

The full spectrum clearly verified this conclusion.



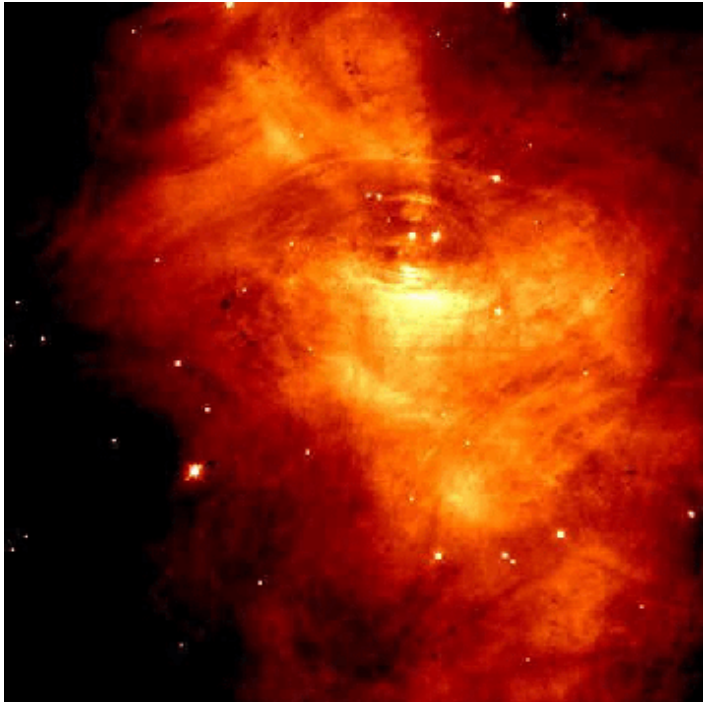
But where do all the electrons come from ?



# Synchrotron-Radiation Rings in Space



CHESS & LEPP



The source of the electrons is probably the central pulsar, or neutron star (seen here as one of the two bright stars near the center).

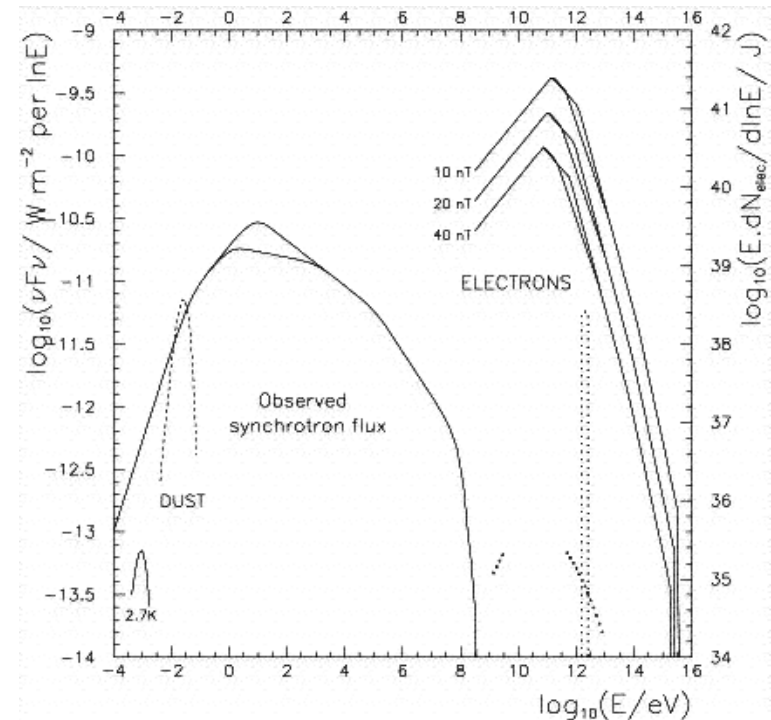
## The Crab nebula:

The spectrum has a power law in the radio region.

The radiation is to about 7% polarized.

This was an early hint for synchrotron radiation.

The full spectrum clearly verified this conclusion.



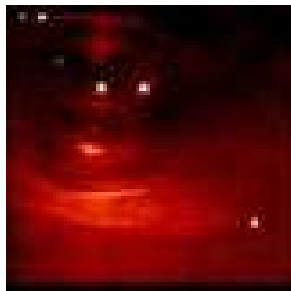
But where do all the electrons come from ?



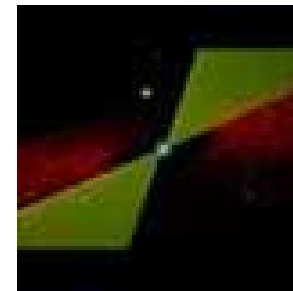
# Pulsars and Radiation Rings



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Moving radiation rings  
In the Crab nebula.



Animation explaining  
radiation rings.