

Unexpected enhancements and reductions of rf spin resonance strengths

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We recently analyzed all available data on spin-flipping stored beams of polarized protons, electrons, and deuterons. Fitting the modified Froissart-Stora equation to the measured polarization data after crossing an rf-induced spin resonance, we found 10–20-fold deviations from the depolarizing resonance strength equations used for many years. The polarization was typically manipulated by linearly sweeping the frequency of an rf dipole or rf solenoid through an rf-induced spin resonance; spin-flip efficiencies of up to 99.9% were obtained. The Lorentz invariance of an rf dipole's transverse $\int Bdl$ and the weak energy dependence of its spin resonance strength \mathcal{E} together imply that even a small rf dipole should allow efficient spin flipping in 100 GeV or even TeV storage rings; thus, it is important to understand these large deviations. Therefore, we recently studied the resonance strength deviations experimentally by varying the size and vertical betatron tune of a 2.1 GeV/c polarized proton beam stored in COSY. We found no dependence of \mathcal{E} on beam size, but we did find almost 100-fold enhancements when the rf spin resonance was near an intrinsic spin resonance.

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I. INTRODUCTION

Polarized beam experiments have become an important part of the programs in storage rings such as the IUCF Cooler Ring [1], AmPS at NIKHEF [2], the MIT-Bates Storage Ring [3], COSY [4], LEP at CERN [5], RHIC at BNL [6] and HERA at DESY [7,8]. To reduce their systematic errors, polarized scattering experiments require frequent spin-direction reversals (spin flips) while the polarized beam is stored. Spin resonances [9–11] induced by either an rf solenoid or rf dipole can produce these spin flips in a well-controlled way [12–25].

In any flat storage ring or circular accelerator with no horizontal magnetic fields, each beam particle's spin precesses around the vertical fields of the ring's dipole magnets. The spin tune ν_s , which is the number of spin precessions during one turn around the ring, is proportional to the particle's energy

$$\nu_s = G\gamma, \quad (1)$$

where $G = (g - 2)/2$ is the particle's gyromagnetic anomaly ($G_p = 1.792847$, $G_e = 1.15967 \times 10^{-3}$, and $G_d = -0.142987$) and γ is its Lorentz energy factor. The vertical polarization can be perturbed by an rf magnet's horizontal rf magnetic field. This perturbation can induce an rf depolarizing resonance, which can flip the spin direction of stored polarized particles [12–25]; the reso-

nance's frequency is

$$f_r = f_c(k \pm \nu_s), \quad (2)$$

where f_c is the particle's circulation frequency and k is an integer.

Ramping an rf magnet's frequency through f_r can flip each particle's spin. The modified [15,19] Froissart-Stora equation [9] relates the beam's initial polarization P_i to its final polarization P_f after crossing the resonance,

$$P_f = P_i \left\{ (1 + \hat{\eta}) \exp \left[\frac{-(\pi \mathcal{E}_{FS} f_c)^2}{\Delta f / \Delta t} \right] - \hat{\eta} \right\}; \quad (3)$$

the parameter $\hat{\eta}$ is the limiting spin-flip efficiency and the ratio $\Delta f / \Delta t$ is the resonance crossing rate, where Δf is the ramp's frequency range during the ramp time Δt . Equation (3) is valid if Δf is larger than the spin resonance's width. We will later discuss how one obtains the spin resonance strength \mathcal{E}_{FS} from Eq. (3).

One can also obtain the resonance strength \mathcal{E}_{Bdl} due to an rf solenoid or rf dipole using

$$\text{solenoid: } \mathcal{E}_{Bdl} = \frac{1}{\pi\sqrt{2}} \frac{e(1+G)}{p} \int B_{\text{rms}} dl, \quad (4)$$

$$\text{dipole: } \mathcal{E}_{Bdl} = \frac{1}{\pi\sqrt{2}} \frac{e(1+G\gamma)}{p} \int B_{\text{rms}} dl, \quad (5)$$

where e is the particle's charge, p is its momentum, and $\int B_{\text{rms}} dl$ is the rf magnet's rms magnetic field integral in its rest frame [26–28]. These equations are for an ideal flat circular accelerator with a point rf magnet causing the only perturbation of the spin motion. Other authors [29,30] derived a similar equation that differed by a factor of 2 in the denominator

$$\text{dipole : } \mathcal{E}_{Bdl} = \frac{1}{\pi 2\sqrt{2}} \frac{e(1+G\gamma)}{p} \int B_{\text{rms}} dl. \quad (6)$$

II. ANALYSIS OF EXISTING DATA

To understand this factor of 2 difference, we made a compilation of all available experimental data [31] that allowed a simultaneous evaluation of the spin resonance strength \mathcal{E}_{Bdl} , obtained from Eq. (4) or (5), and the spin resonance strength \mathcal{E}_{FS} obtained from Eq. (3), the modified Froissart-Stora formula. This compilation indicated that for many experiments \mathcal{E}_{Bdl} and \mathcal{E}_{FS} disagree with both [26,27] and [29,30] by factors of 0.1, 10, or more.

The resonance strength \mathcal{E}_{FS} can be obtained by first measuring the final beam polarization P_f after ramping an rf magnet's frequency by a range Δf during a time Δt through a spin resonance; then P_f is plotted vs the different Δt values and the data are fit to Eq. (3) with \mathcal{E}_{FS} as a fit parameter. Examples of these curves are shown in Figs. 1 and 2 for protons [23] and deuterons [25], respectively. For protons, using Eq. (5) gave \mathcal{E}_{Bdl} of $(80 \pm 4) \times 10^{-6}$, while \mathcal{E}_{FS} was $(580 \pm 10) \times 10^{-6}$, which is about 7.3 times larger. For deuterons, \mathcal{E}_{Bdl} was $(16 \pm 1) \times 10^{-6}$, while

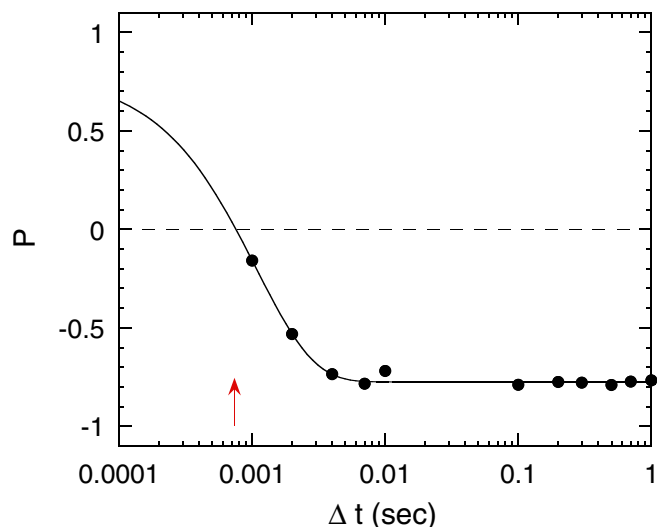


FIG. 1. (Color) Measured proton polarization at 2.1 GeV/c is plotted vs rf dipole ramp time Δt [23]. The rf dipole's frequency range Δf was 8 kHz; its $\int Bdl$ was 0.46 ± 0.03 T mm; thus, Eq. (5) gives $\mathcal{E}_{Bdl} = (80 \pm 4) \times 10^{-6}$. The fit to Eq. (3) gives $\mathcal{E}_{FS} = (580 \pm 10) \times 10^{-6}$; the arrow shows the Δt where $P = 0$.

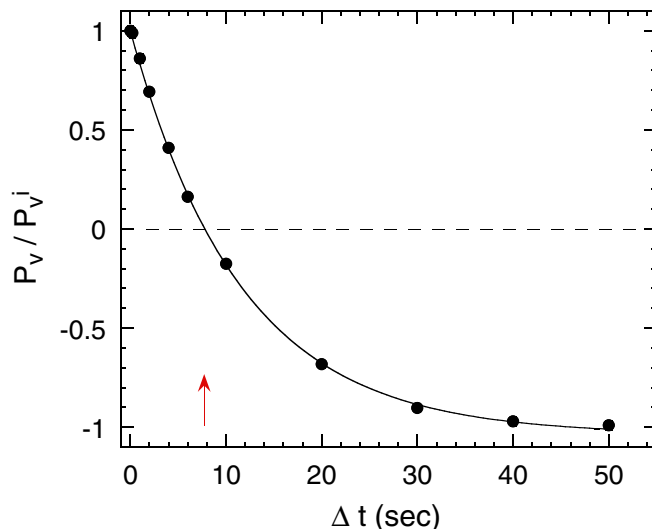


FIG. 2. (Color) Measured average vector deuteron polarization ratios at 1850 MeV/c are plotted vs rf dipole ramp time Δt [25]. The rf dipole's frequency range Δf was 200 Hz; its $\int Bdl$ was 0.54 ± 0.03 T mm; thus, Eq. (5) gives $\mathcal{E}_{Bdl} = (16 \pm 1) \times 10^{-6}$. The fit to Eq. (3) gives $\mathcal{E}_{FS} = (1.17 \pm 0.01) \times 10^{-6}$; the arrow shows the Δt where $P = 0$.

\mathcal{E}_{FS} was $(1.17 \pm 0.01) \times 10^{-6}$, which is about 14 times smaller. Note that these two experiments were both done with the same rf dipole at COSY; thus, these large strength deviations could not be due to incorrect calibrations of $\int Bdl$.

We examined all available publications and logbooks on spin-flipping experiments [31] that had both: information on the rf magnet's $\int Bdl$, and data for a polarization vs ramp time Δt curve. We also checked in each experiment for any other spin resonances near the rf-induced resonance. Fits to Eq. (3), the modified Froissart-Stora equation, gave the measured spin resonance strength for each experiment and its error. We also calculated the spin resonance strength from each rf magnet's $\int Bdl$ using Eq. (4) or (5). For the experiments at IUCF and MIT, there were no precise measurements of the rf magnet's $\int Bdl$; thus, we assumed a $\pm 10\%$ error. For the more recent experiments at COSY, the rf magnet's $\int Bdl$ was known with a $\pm 5\%$ uncertainty.

Then we took the ratio of the two spin resonance strengths, $\mathcal{E}_{FS}/\mathcal{E}_{Bdl}$, for each experiment and obtained its error using simple error propagation formulas. Figure 3 shows these ratios plotted vs the ramp time Δt at which P passes through 0 in each Δt curve. The dashed [26–28] and solid [29,30] lines show the predicted ratios.

III. COSY EXPERIMENT

To better understand these deviations, we recently studied experimentally the dependence of an rf resonance's strength on the beam size and on its distance from an

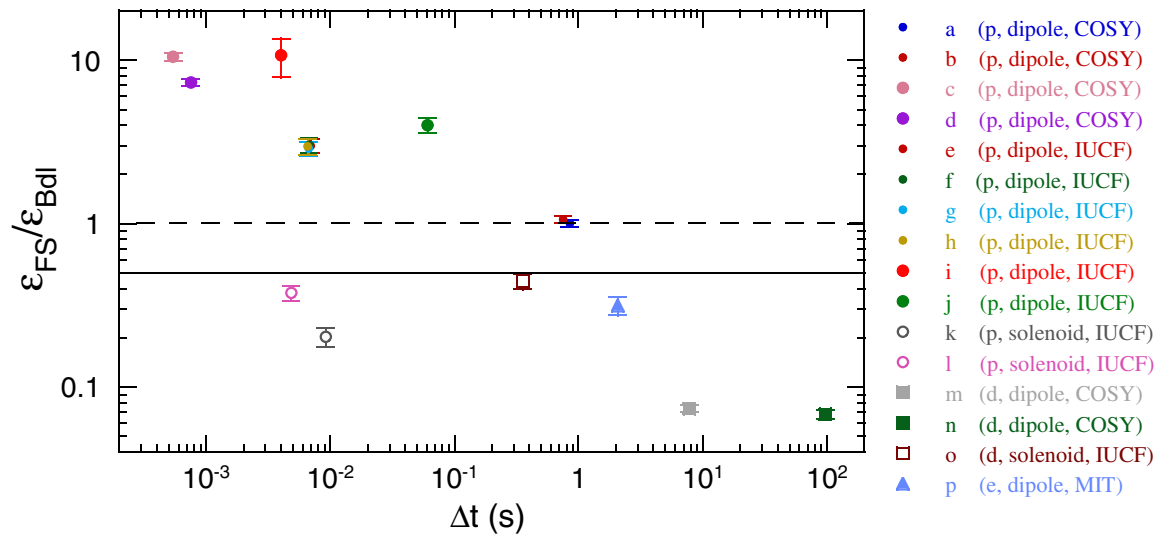


FIG. 3. (Color) Ratio of \mathcal{E}_{FS} to \mathcal{E}_{Bdl} is plotted vs the ramp time Δt giving $P = 0$ in each Δt curve [31]; for example, points (d) and (m) come from Figs. 1 and 2, respectively. \mathcal{E}_{FS} is the resonance strength obtained by fitting data in each Δt curve to Eq. (3), the modified Froissart-Stora equation; \mathcal{E}_{Bdl} is obtained using each data point's $\int Bdl$ in Eq. (4) or (5). The dashed and solid lines show the ratios predicted by [26–28] and [29,30], respectively. Note that points (a) and (b) overlap, as do points (e)–(h). The data for all points are listed in Table I.

intrinsic spin resonance. We did this by varying the size and the vertical betatron tune of a 2.1 GeV/c polarized proton beam stored in COSY.

The apparatus used for this experiment, including the COSY storage ring [32–35], the EDDA detector [36], the electron Cooler [37], the low energy polarimeter, the injector cyclotron, and the polarized ion source [38–40], are indicated in Fig. 4, along with the rf dipole. The beam emerging from the polarized H^- ion source was accelerated by the cyclotron to COSY's 45 MeV injection energy. Then the low energy polarimeter measured the beam's polarization before injection into COSY to monitor the stable operation and polarization of the ion source. The electron Cooler reduced the beam's size and momentum spread at injection energy. A 24.5 keV electron beam of about 170 mA, with a diameter of about 25 mm, cooled the protons in a 2-m-long cooling section; it took 10 s to cool the proton beam to its equilibrium emittances in both the longitudinal and transverse dimensions. This decreased its initial momentum spread by about a factor of 10 to $\Delta p/p = 9 \times 10^{-5}$ FWHM.

We manipulated the beam's polarization using a ferrite-core rf dipole, which contained a copper coil of 8 turns; their spacing was optimized to produce a uniform radial magnetic field. The rf dipole was part of an LC resonant circuit, which operated near $f_r = 902.6$ kHz, at an rf voltage of 1 kV rms giving an rf $\int B_{rms} dl$ of 0.19 ± 0.01 T mm.

We measured the polarization in COSY using the EDDA detector [4,36]; we reduced its systematic errors by cycling the polarized source between the up and down vertical

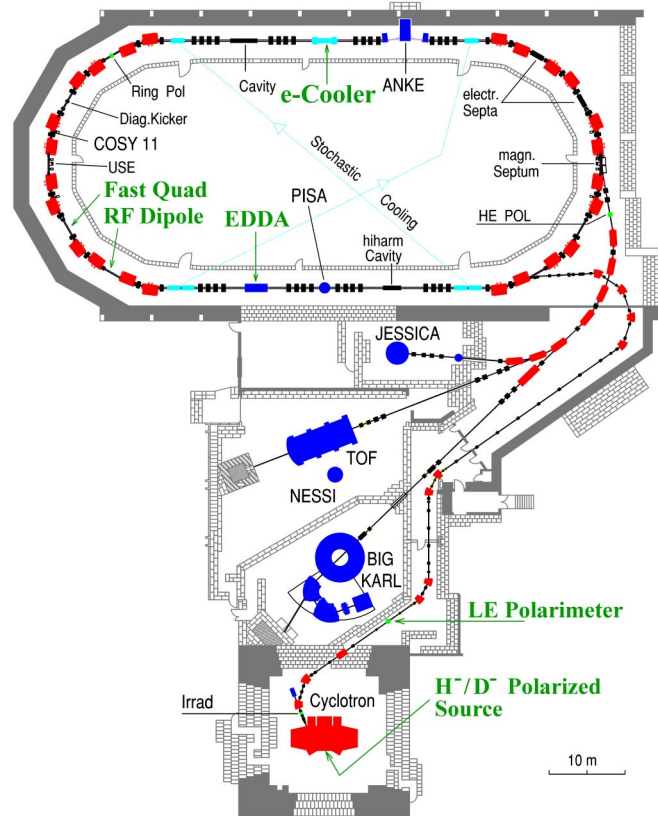


FIG. 4. (Color) Layout of the COSY storage ring, with its injector cyclotron and polarized ion source. Also note the rf dipole, the fast quadrupole, EDDA, the low energy (LE) polarimeter, and the electron Cooler.

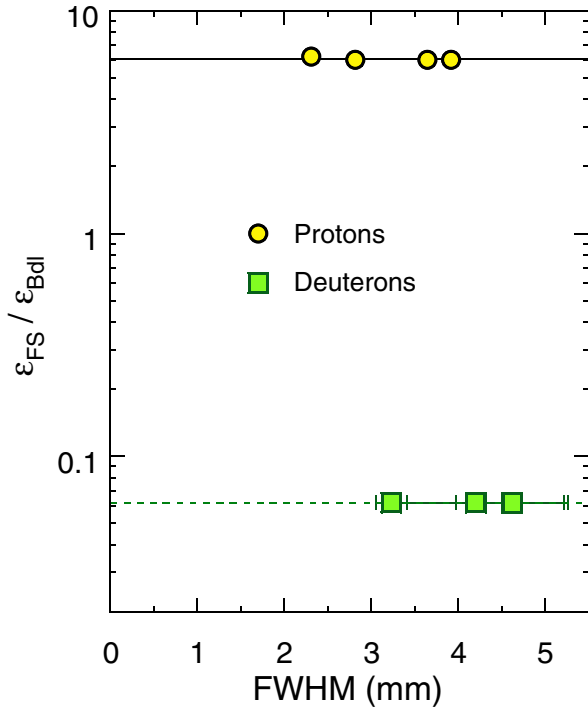


FIG. 5. (Color) Ratio of \mathcal{E}_{FS} to \mathcal{E}_{Bdl} is plotted vs the vertical beam size (FWHM). \mathcal{E}_{FS} is the resonance strength obtained by fitting the Δt curve for each data point to Eq. (3); \mathcal{E}_{Bdl} is obtained using each data point's $\int Bdl$ in Eq. (5). Fits to horizontal straight lines give resonance strength ratios of 6.07 ± 0.04 for protons and 0.0618 ± 0.0002 for deuterons. The ν_y values were 3.525 for protons and 3.60 for deuterons.

polarization states. The rf acceleration cavity was turned off and shorted during COSY's flattop; thus, there were no synchrotron sideband effects [14,41,42]. The measured flattop polarization, before spin manipulation, was typically 50% to 60%.

We determined the resonance's position and the upper limit on its width using the procedure described in [23]. Then, we obtained the resonance strength \mathcal{E}_{FS} , as in Figs. 1 and 2, for different fast quadrupole strengths. The fast quadrupole increased the beam's emittance at the start of COSY's flattop, before we manipulated the beam's polarization with the rf dipole. We measured beam's vertical profile for each fast quadrupole strength by moving a thin fiber target vertically through the beam, while measuring the current due to its scattering events as the fiber moved. The $\mathcal{E}_{FS}/\mathcal{E}_{Bdl}$ ratios are plotted against the beam's vertical size (FWHM) in Fig. 5. We also plot deuteron data obtained in a similar way during a December 2004 run [25]. The proton and deuteron data both show no dependence on the beam's vertical size. The fits to horizontal straight lines give resonance strength ratios of 6.07 ± 0.04 for protons and 0.0618 ± 0.0002 for deuterons.

We also obtained the resonance strength \mathcal{E}_{FS} , as in Figs. 1 and 2, for different values of the vertical betatron tune ν_y . We plotted the resulting $\mathcal{E}_{FS}/\mathcal{E}_{Bdl}$ ratios against ν_y in Fig. 6; we also plotted our 2004 deuteron data [25]. Notice the $\nu_s = 8 - \nu_y$ first-order intrinsic spin resonance for protons, which let us study the rf resonance strength in its proximity; there were also four third-order resonances.

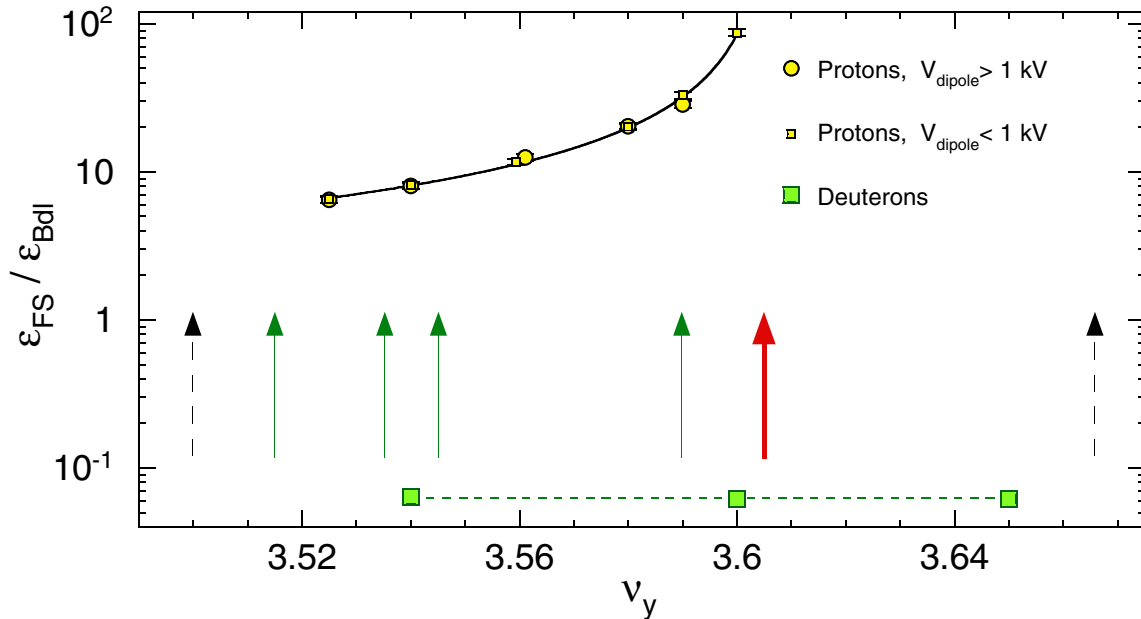


FIG. 6. (Color) Ratio of \mathcal{E}_{FS} to \mathcal{E}_{Bdl} is plotted vs the vertical betatron tune ν_y . \mathcal{E}_{FS} is the resonance strength obtained by fitting the data in the Δt curve for each ν_y to Eq. (3); \mathcal{E}_{Bdl} is obtained using each data point's $\int Bdl$ in Eq. (5). The proton data fit to Eq. (7), gives ν_r of 3.6060 ± 0.0005 , A of 0.44 ± 0.46 and B of 0.50 ± 0.03 . Fitting the deuteron data to a horizontal straight line gives a resonance strength ratio $\mathcal{E}_{FS}/\mathcal{E}_{Bdl} = 0.063 \pm 0.002$. The betatron beam resonances are shown by the black dashed arrows; the 1st and 3rd order proton spin resonances are shown by the red (bold) and green arrows, respectively.

Using references [43,44], one can obtain the dependence of $\mathcal{E}_{FS}/\mathcal{E}_{Bdl}$ on the distance between the vertical betatron tune ν_y and the rf spin resonance's tune $\nu_r \equiv k \pm f_r/f_c$, where k is an integer; this dependence is

$$\mathcal{E}_{FS}/\mathcal{E}_{Bdl} = A + \frac{B}{|\nu_y - \nu_r|}. \quad (7)$$

Fitting the proton data in Fig. 6 to Eq. (7) gave A of 0.44 ± 0.46 , B of 0.50 ± 0.03 , and ν_r of

3.6060 ± 0.0005 , which was very near the calculated value of 3.605 for the first-order intrinsic spin resonance. The parameter B depends on many details of the ring. The parameter A should give the predicted resonance strength ratio $\mathcal{E}_{FS}/\mathcal{E}_{Bdl}$ of either 1 [26–28] or 0.5 [29,30] when one is infinitely far from any intrinsic spin resonances. Clearly our data could not determine A with good precision.

The deuteron data in Fig. 6 show no dependence on ν_y ; its fit to a horizontal straight line gave a resonance strength ratio of 0.063 ± 0.002 , which agrees with the result from

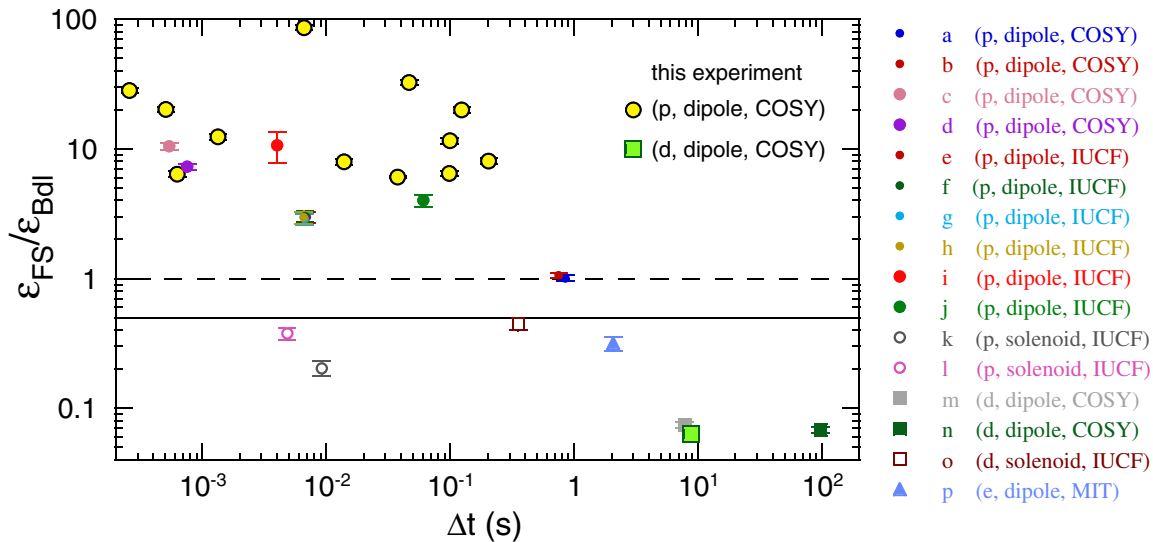


FIG. 7. (Color) Ratio of \mathcal{E}_{FS} to \mathcal{E}_{Bdl} is plotted vs the ramp time Δt giving $P = 0$ in each Δt curve [31]. \mathcal{E}_{FS} is the resonance strength obtained by fitting the data in each Δt curve to Eq. (3); \mathcal{E}_{Bdl} is obtained using each data point's $\int Bdl$ in Eq. (4) or (5). The dashed and solid lines show the ratios predicted by [26–28] and [29,30], respectively. Note that points (a) and (b) overlap, as do points (e)–(h). The data for all points are listed in Table I.

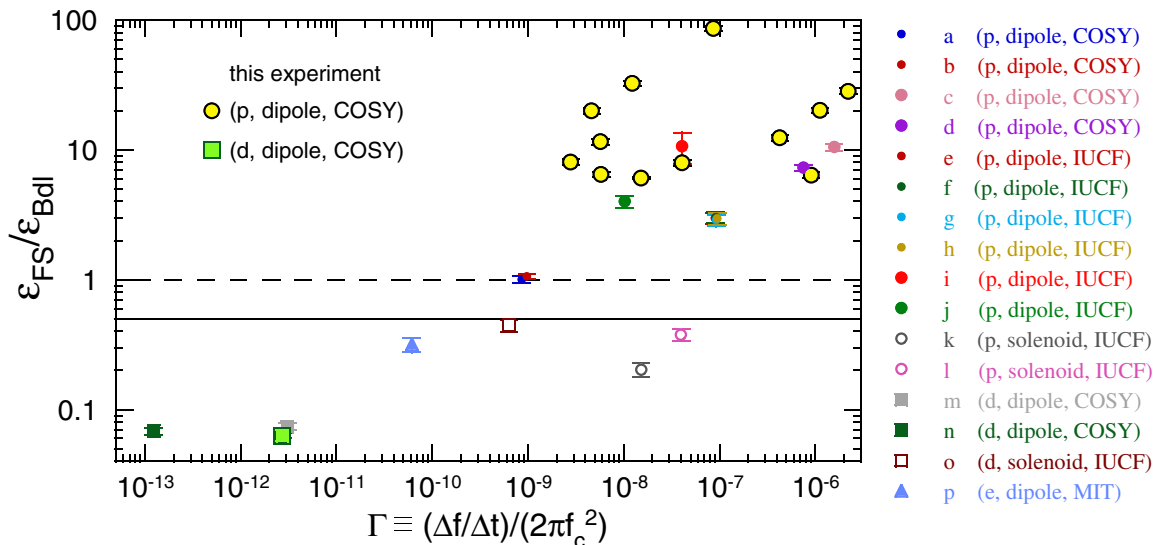


FIG. 8. (Color) Ratio of \mathcal{E}_{FS} to \mathcal{E}_{Bdl} is plotted vs the dimensionless crossing rate Γ giving $P = 0$ in each Δt curve [31]. \mathcal{E}_{FS} is the resonance strength obtained by fitting the data in each Δt curve to Eq. (3); \mathcal{E}_{Bdl} is obtained using each data point's $\int Bdl$ in Eq. (4) or (5). The dashed and solid lines show the ratios predicted by [26–28] and [29,30], respectively. Note that points (a) and (b) overlap, as do points (e)–(h). The data for all points are listed in Table I.

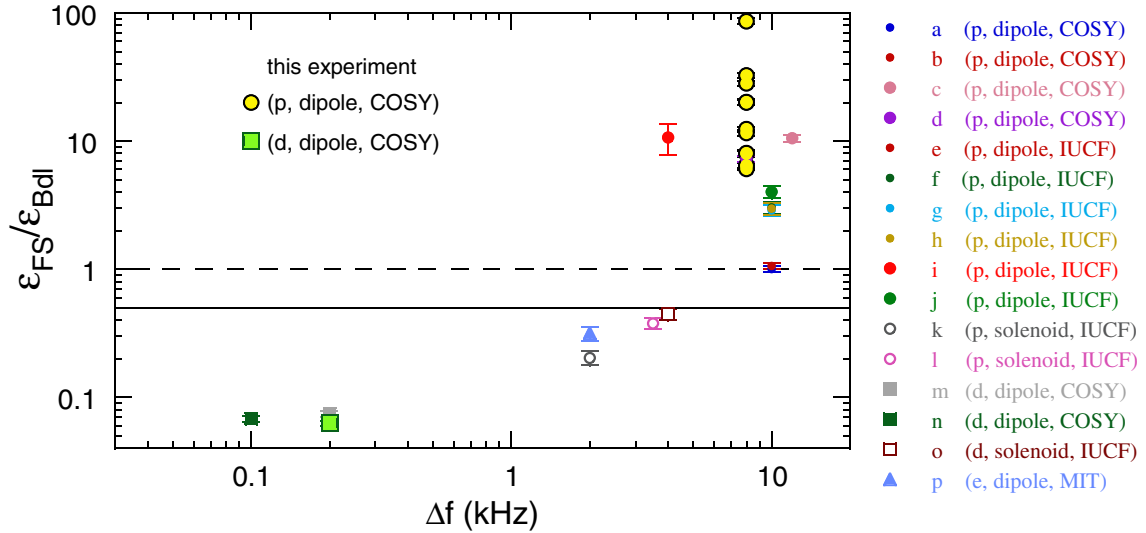


FIG. 9. (Color) Ratio of \mathcal{E}_{FS} to \mathcal{E}_{Bdl} is plotted vs the frequency ramp range Δf used for each Δt curve [31]. \mathcal{E}_{FS} is the resonance strength obtained by fitting the data in each Δt curve to Eq. (3); \mathcal{E}_{Bdl} is obtained using each data point's $\int Bdl$ in Eq. (4) or (5). The dashed and solid lines show the ratios predicted by [26–28] and [29,30], respectively. Note that points (a) and (b) overlap, as do points (e)–(h). The data for all points are listed in Table I.

TABLE I. Data for Figs. 3 and 7–9 along with relevant experimental parameters for this experiment and earlier experiments [31]. The letters (a)–(p) in column 1 stand for the references listed in [31], while the letters p , d , and e in column 3 denote protons, deuterons, and electrons, respectively. We assumed a $\pm 10\%$ error in the rf magnets' $\int Bdl$ for the experiments at IUCF and MIT, and a $\pm 5\%$ error in the rf dipole's $\int Bdl$ for our experiments at COSY. Note that, for the bottom 12 COSY proton points, ν_y was varied as explained in Figs. 5 and 6.

Ring	Magnet type	f_c (MHz)	p (GeV/c)	$\int B_{rms} dl$ (T mm)	\mathcal{E}_{Bdl} ($\times 10^{-6}$)	Flips	Δt ($P=0$) (ms)	Δf (kHz)	$\Gamma \equiv \frac{\Delta f / \Delta t}{2\pi f_c^2}$	\mathcal{E}_{FS} ($\times 10^{-6}$)	$\mathcal{E}_{FS}/\mathcal{E}_{Bdl}$ ($\times 10^{-6}$)
a	COSY p Dipole	1.471 17	1.941	0.11	19.59	11	849	10	8.66×10^{-10}	19.7 ± 0.3	1.00 ± 0.05
b	COSY p Dipole	1.491 89	2.100	0.46	79.80	11	0.541	12	1.59×10^{-6}	840 ± 30	10.5 ± 0.6
c	COSY p Dipole	1.491 89	2.100	0.46	79.80	11	0.541	12	1.59×10^{-6}	840 ± 30	10.5 ± 0.6
d	IUCF p Dipole	1.597 84	0.489	0.16	66.70	1	6.97	10	8.94×10^{-8}	199 ± 2	3.0 ± 0.3
e	IUCF p Dipole	1.597 84	0.489	0.16	66.70	1	6.97	10	8.94×10^{-8}	199 ± 2	3.0 ± 0.3
f	IUCF p Dipole	1.597 84	0.489	0.17	70.87	10	6.77	10	9.21×10^{-8}	202 ± 4	3.0 ± 0.3
g	IUCF p Dipole	1.597 84	0.489	0.17	70.87	10	6.72	10	9.28×10^{-8}	204 ± 3	2.9 ± 0.3
h	IUCF p Dipole	1.597 84	0.489	0.17	70.87	1	6.56	10	9.50×10^{-8}	210 ± 9	3.0 ± 0.3
i	IUCF p Dipole	1.965 8	0.649	0.042	14.03	1	4.03	4	4.09×10^{-8}	150 ± 40	11 ± 3
j	IUCF p Dipole	1.597 84	0.489	0.042	17.69	1	60.7	10	1.03×10^{-8}	71 ± 3	4.0 ± 0.4
k	IUCF p Solenoid	1.504 90	0.454	1.131	469.9	1	9.19	2	1.53×10^{-8}	95 ± 8	0.20 ± 0.03
l	IUCF p Solenoid	1.697 3	0.529	0.99	352.7	1	4.86	3.5	3.98×10^{-8}	133 ± 2	0.38 ± 0.04
m	COSY d Dipole	1.147 43	1.850	0.54	15.75	1	7.81×10^3	0.2	3.10×10^{-12}	1.166 ± 0.009	0.074 ± 0.004
n	COSY d Dipole	1.147 43	1.850	0.15	4.375	1	9.76×10^4	0.1	1.24×10^{-13}	0.298 ± 0.006	0.068 ± 0.004
o	IUCF d Solenoid	1.677 55	1.042	0.70	38.88	1	356	4	6.35×10^{-10}	17.3 ± 0.6	0.44 ± 0.05
p	MIT e Dipole	1.576 16	0.670	0.070	17.78	1	2.07×10^3	2	6.18×10^{-11}	5.6 ± 0.4	0.32 ± 0.04
This experiment											
COSY	d Dipole	1.147 43	1.850	0.60	17.50	1	8.81×10^3	0.2	2.74×10^{-12}	1.101 ± 0.006	0.063 ± 0.002
COSY	p Dipole	1.491 89	2.100	0.57	99.14	1	0.628	8	9.11×10^{-7}	634 ± 10	6.4 ± 0.3
				0.20	35.02	1	1.34	8	4.27×10^{-7}	434 ± 2	12.4 ± 0.6
				0.20	35.02	1	0.508	8	1.13×10^{-6}	705 ± 3	20 ± 1
				0.20	35.02	1	0.259	8	2.21×10^{-6}	988 ± 9	28.2 ± 1.4
				0.013	2.262	1	6.62	8	8.64×10^{-8}	195 ± 1	86 ± 4
				0.013	2.262	1	46.6	8	1.23×10^{-8}	73.6 ± 0.3	32.6 ± 1.6
				0.013	2.262	1	124	8	4.62×10^{-9}	45.2 ± 0.2	20 ± 1
				0.025	4.359	1	99.7	8	5.74×10^{-9}	50.3 ± 0.2	11.5 ± 0.6
				0.025	4.359	1	204	8	2.80×10^{-9}	35.2 ± 0.5	8.1 ± 0.4
				0.098	16.93	1	14.0	8	4.10×10^{-8}	135 ± 1	7.9 ± 0.4
				0.045	7.820	1	98.6	8	5.80×10^{-9}	50.6 ± 0.4	6.5 ± 0.3
				0.078	13.48	1	37.7	8	1.52×10^{-8}	81.9 ± 0.3	6.07 ± 0.04

Fig. 5. There were no intrinsic spin resonances for deuterons near the operating tunes; thus, this deviation can not be due to a nearby intrinsic spin resonance.

In an effort to experimentally determine what might be causing these huge deviations in $\mathcal{E}_{FS}/\mathcal{E}_{Bdl}$, we plotted the resonance strength ratios $\mathcal{E}_{FS}/\mathcal{E}_{Bdl}$, including the new experimental data, in Figs. 7–9 against three different parameters:

Fig. 7: ramp time Δt where $P = 0$ in each Δt curve,

Fig. 8: crossing rate Γ where $P = 0$ in each Δt curve,

Fig. 9: frequency ramp range Δf in each Δt curve.

The data seem scattered randomly in both Figs. 7 and 8. There is much less scatter in Fig. 9: all data with small Δf have $\mathcal{E}_{FS}/\mathcal{E}_{Bdl}$ values far below 1, while all data with large Δf have $\mathcal{E}_{FS}/\mathcal{E}_{Bdl}$ values near or above 1.

Clearly there still is no complete understanding of what is causing these very large resonance strength deviations. Figure 6 indicates that the large deviations above the predicted values of 1 [26–28] and 0.5 [29,30] are due to a nearby intrinsic spin resonance. Figure 9 suggests that the deviations below these lines could be due to the small Δf frequency ramp range. Figure 9 also shows that, for deuterons, all $\mathcal{E}_{FS}/\mathcal{E}_{Bdl}$ ratios are far below 1 for rf dipoles, but closer to 1 for rf solenoids. This may suggest anomalous behavior of spin-1 deuterons when spin-manipulated by dipoles.

Studying these deviations helped inspire a theoretical paper [45], which may be the first generalization of the Froissart-Stora equation [9] since its publication in 1960.

IV. SUMMARY AND CONCLUSION

In summary, by compiling existing data and fitting it to the Froissart-Stora equation, we first found 10–20-fold deviations from the spin resonance strength equation that has been used for many years. We recently studied these deviations experimentally at COSY. We found no dependence on the beam’s vertical size; however, we did find even larger deviations, of almost 100, when the rf spin resonance was near an intrinsic spin resonance. This proximity to an intrinsic resonance probably explains the 100-times-larger-than-predicted deviations. The 10–20-times-smaller-than-predicted deviations for deuterons could be due to the following: the small Δf sweep used to flip the deuteron spin; the resulting very slow crossing rate (Fig. 8); or some unexpected behavior of spin-1 deuterons in an rf dipole. We plan to test these and other possibilities at COSY, to learn if any are valid.

By adiabatically ramping the frequency of an rf dipole or rf solenoid through an rf-induced spin resonance, one can flip the polarization of a stored beam of protons, electrons or deuterons with an efficiency of up to 99.9%. An rf dipole only slightly stronger than our 0.5 Tmm dipole should allow efficient spin flipping of polarized protons in the 200 GeV RHIC and perhaps someday in the 7 TeV LHC.

Thus, these sometimes 100-fold deviations must be better understood both theoretically and experimentally.

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