



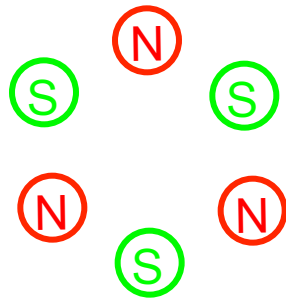
Sextupoles (revisited)



CHESS & LEPP

$$\psi = \Psi_3 \operatorname{Im}\{(x - iy)^3\} = \Psi_3 \cdot (y^3 - 3x^2y) \Rightarrow \vec{B} = -\vec{\nabla}\psi = \Psi_3 3 \begin{pmatrix} 2xy \\ x^2 - y^2 \end{pmatrix}$$

C_3 Symmetry



i) Sextupole fields hardly influence the particles close to the center, where one can linearize in x and y .

ii) In linear approximation a by Δx shifted sextupole has a quadrupole field.

$$\vec{B} = -\vec{\nabla}\psi = \Psi_3 3 \begin{pmatrix} 2xy \\ x^2 - y^2 \end{pmatrix}$$

iii) When Δx depends on the energy, one can build an **energy dependent quadrupole**.

$$x \mapsto \Delta x + x$$

$$\vec{B} \approx \Psi_3 3 \begin{pmatrix} 2xy \\ x^2 - y^2 \end{pmatrix} + 6\Psi_3 \Delta x \begin{pmatrix} y \\ x \end{pmatrix} + O(\Delta x^2)$$

$$k_2 = 3!\Psi_3 \Rightarrow k_1 = k_2 \Delta x$$



Chromaticity ξ = energy dependence of the tune

$$\nu(\delta) = \nu + \frac{\partial \nu}{\partial \delta} \delta + \dots$$

$$\xi = \frac{\partial \nu}{\partial \delta} \quad \text{with} \quad \nu = \frac{\mu}{2\pi}$$

Natural chromaticity ξ_0 = energy dependence of the tune due to quadrupoles only

$$\xi_{x0} = -\frac{1}{4\pi} \oint \beta_x(\hat{s}) k_1(\hat{s}) d\hat{s}$$

$$\xi_{y0} = \frac{1}{4\pi} \oint \beta_y(\hat{s}) k_1(\hat{s}) d\hat{s}$$

Particles with energy difference oscillate around the periodic dispersion leading to a quadrupole effect in sextupoles that also shifts the tune:

$$\xi_x = \frac{1}{4\pi} \oint \beta_x (-k_1 + \eta_x k_2) d\hat{s}$$

$$\xi_y = \frac{1}{4\pi} \oint \beta_y (k_1 - \eta_x k_2) d\hat{s}$$

Typically the the chromaticity ξ is chosen to be slightly positive, between 0 and 3.



$$\begin{pmatrix} x' \\ a' \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -K & 0 \end{pmatrix} \begin{pmatrix} x \\ a \end{pmatrix} + \begin{pmatrix} 0 \\ \Delta f \end{pmatrix}$$

$$\begin{pmatrix} x \\ a \end{pmatrix} = \sqrt{2J} \begin{pmatrix} \sqrt{\beta} & 0 \\ -\frac{\alpha}{\sqrt{\beta}} & \frac{1}{\sqrt{\beta}} \end{pmatrix} \begin{pmatrix} \sin(\psi + \phi_0) \\ \cos(\psi + \phi_0) \end{pmatrix} = \sqrt{2J} \underline{\beta} \vec{S}$$

This would be a solution with constant J and ϕ when $\Delta f=0$.

Variation of constants:

$$\frac{J'}{\sqrt{2J}} \underline{\beta} \vec{S} + \sqrt{2J} \phi_0' \begin{pmatrix} 0 & \sqrt{\beta} \\ -\frac{1}{\sqrt{\beta}} & -\frac{\alpha}{\sqrt{\beta}} \end{pmatrix} \vec{S} = \begin{pmatrix} 0 \\ \Delta f \end{pmatrix}$$

$$\frac{J'}{\sqrt{2J}} \vec{S} + \sqrt{2J} \phi_0' \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \vec{S} = \underline{\beta}^{-1} \begin{pmatrix} 0 \\ \Delta f \end{pmatrix} \quad \text{with} \quad \underline{\beta}^{-1} = \begin{pmatrix} \frac{1}{\sqrt{\beta}} & 0 \\ \frac{\alpha}{\sqrt{\beta}} & \sqrt{\beta} \end{pmatrix}$$

$$\frac{J'}{\sqrt{2J}} = \cos(\psi + \phi_0) \sqrt{\beta} \Delta f \quad , \quad \sqrt{2J} \phi_0' = -\sin(\psi + \phi_0) \sqrt{\beta} \Delta f$$

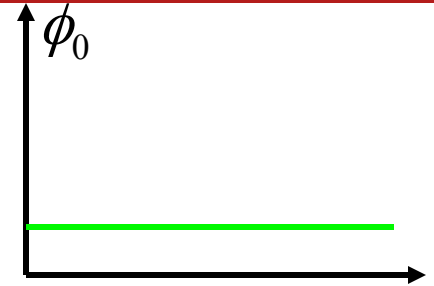


Simplification of linear motion

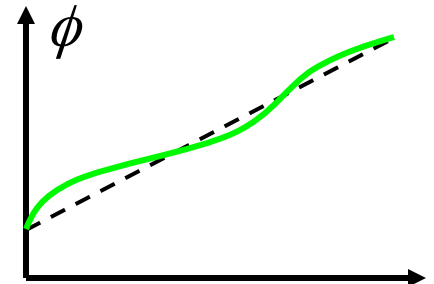


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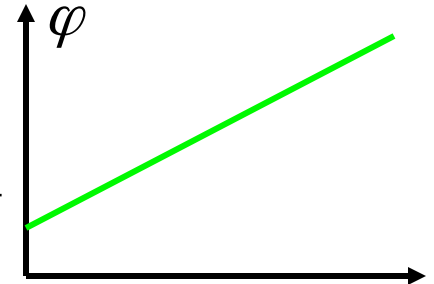
$$\begin{pmatrix} x \\ a \end{pmatrix} = \sqrt{2J} \begin{pmatrix} \sqrt{\beta} & 0 \\ -\frac{\alpha}{\sqrt{\beta}} & \frac{1}{\sqrt{\beta}} \end{pmatrix} \begin{pmatrix} \sin(\psi + \phi_0) \\ \cos(\psi + \phi_0) \end{pmatrix} \Rightarrow \begin{aligned} J' &= 0 \\ \phi_0' &= 0 \end{aligned}$$



$$\begin{pmatrix} x \\ a \end{pmatrix} = \sqrt{2J} \begin{pmatrix} \sqrt{\beta} & 0 \\ -\frac{\alpha}{\sqrt{\beta}} & \frac{1}{\sqrt{\beta}} \end{pmatrix} \begin{pmatrix} \sin \phi \\ \cos \phi \end{pmatrix} \Rightarrow \begin{aligned} J' &= 0 \\ \phi' &= \frac{1}{\beta} \end{aligned}$$



$$\begin{pmatrix} x \\ a \end{pmatrix} = \sqrt{2J} \begin{pmatrix} \sqrt{\beta} & 0 \\ -\frac{\alpha}{\sqrt{\beta}} & \frac{1}{\sqrt{\beta}} \end{pmatrix} \begin{pmatrix} \sin(\psi - \mu \frac{s}{L} + \varphi) \\ \cos(\psi - \mu \frac{s}{L} + \varphi) \end{pmatrix} \Rightarrow \begin{aligned} J' &= 0 \\ \varphi' &= \mu \frac{1}{L} \end{aligned}$$



$$\tilde{\psi} = \psi - \mu \frac{s}{L} \Rightarrow \tilde{\psi}(s + L) = \tilde{\psi}(s)$$

Corresponds to Floquet's Theorem