

Cornell University



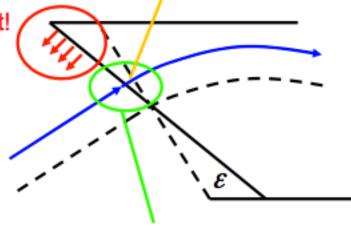
Top view:

x tan(ε)

Fringe field has a horizontal

field component!

Horizontal focusing with
$$\Delta x' = -x \frac{\tan(\varepsilon)}{\rho}$$



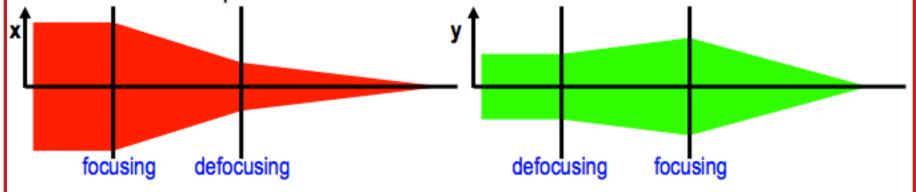
Extra bending focuses!

The longitudinal field above the enter plain $\Delta y' = y \frac{\tan(\varepsilon)}{\rho}$ defocuses, turns out to:

Quadrupole effect: focusing in x and defocusing in y or defocusing in x and focusing in y.



Transverse fields defocus in one plane if they focus in the other plane. But two successive elements, one focusing the other defocusing, can focus in both planes:





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CHESS & LEPP

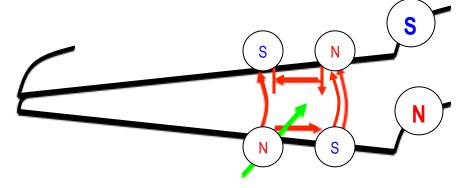
The isocyclotron with constant

$$\omega_z = \frac{q}{m_0 \gamma(E)} B_z(r(E))$$

Up to 600MeV but this vertically defocuses the beam.

Edge focusing is therefore used.







First Medical Applications

CHESS & LEPP

• 1939: Lawrence uses 60' cyclotron for 9MeV protons, 19MeV deuterons, and 35MeV 4He. First tests of tumor therapy with neutrons via d + t \mapsto n + α With 200-800keV d to get 10MeV neutrons.



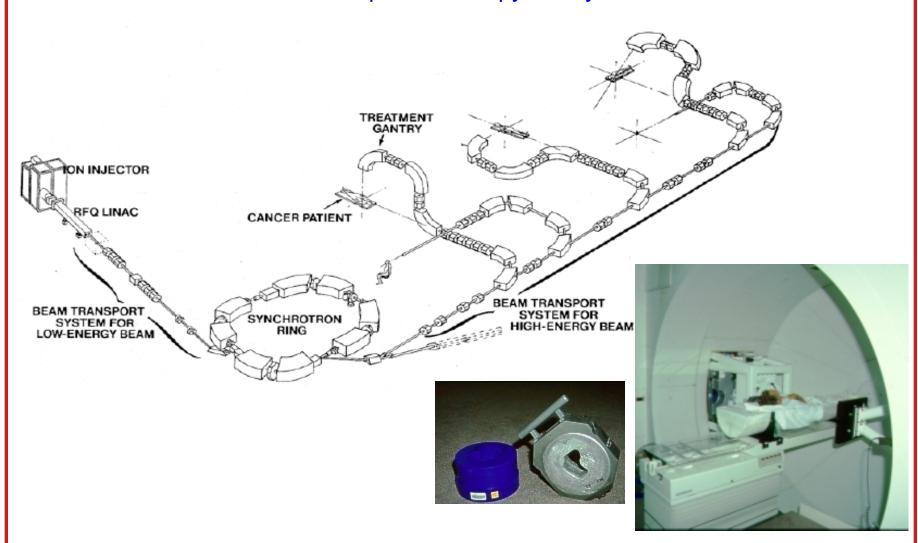




Modern Nuclear Therapy



The Loma Linda proton therapy facility

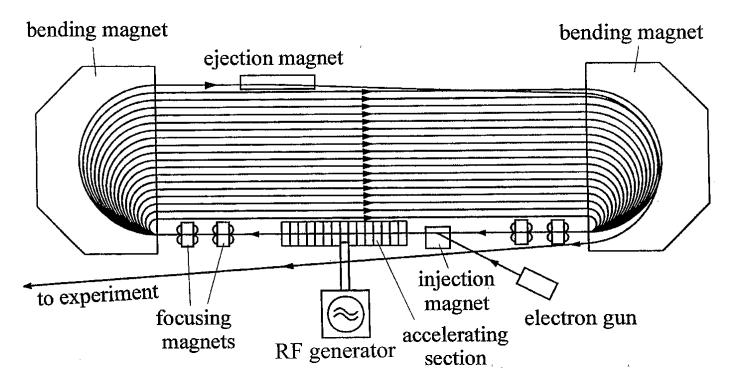




The microtron



- Electrons are quickly relativistic and cannot be accelerated in a cyclotron.
- •In a microtron the revolution frequency changes, but each electron misses an integer number of RF waves.



- Today: Used for medical applications with one magnet and 20MeV.
- •Nuclear physics: MAMI designed for 820MeV as race track microtron.



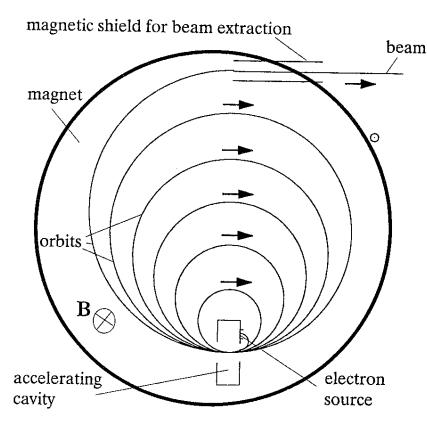
The microtron condition



CHESS & LEPP

 The extra time that each turn takes must be a multiple of the RF period.

$$\frac{dp}{dt} = qvB \Rightarrow \rho = \frac{dl}{d\varphi} = \frac{vdt}{dp/p} = \frac{p}{qB}$$



$$dp = pd\varphi$$

$$d\varphi = \frac{p}{qB}$$

$$\Delta t = 2\pi \left(\frac{\rho_{n+1}}{v_{n+1}} - \frac{\rho_n}{v_n}\right)$$

$$= \frac{2\pi}{qB} \left(m_0 \gamma_{n+1} - m_0 \gamma_n\right) = \frac{2\pi}{qBc^2} \Delta K$$

B=1T, n=1, and f_{RF} =3GHz leads to 4.78MeV This requires a small linear accelerator.

$$\Delta K = n \frac{qBc^2}{\omega_{RF}} \quad \text{for an integer n}$$

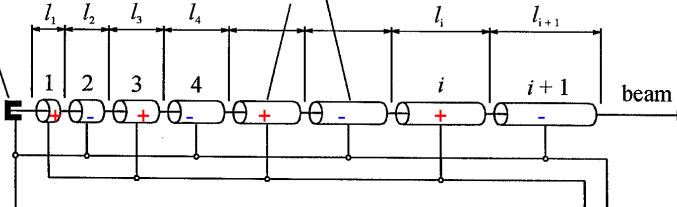


Wideroe linear accelerator



ion source

drift tubes (Faraday cage)





Wideroe

non-relativistic:

$$K_n = nqU_{\text{max}}\sin\psi_0 = \frac{1}{2}mv_n^2$$

$$l_n = \frac{1}{2} v_n T_{RF} = \frac{1}{2} \beta_n \lambda_{RF} \propto \sqrt{n}$$

Called the π or the $1/2\beta\lambda$ mode

RF generator

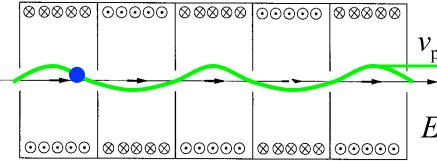


Accelerating cavities



1933: J.W. Beams uses resonant cavities for acceleration

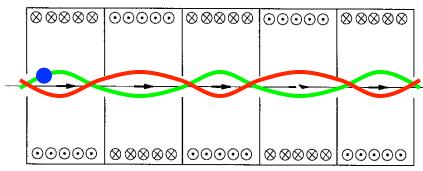
Traveling wave cavity:



$$v_{\text{phase}} = \frac{\omega}{k} = v_{\text{particle}}$$
 Here v=c for electrons

$$E(t,s) \approx E_{\text{max}} \sin(\omega t - ks)$$

Standing wave cavity:



$$\frac{\omega}{k} = v_{\text{particle}}$$

$$E(t,s) \approx E_{\text{max}} \sin(\omega t) \sin(ks)$$

$$E(\frac{s}{v_{\text{particle}}}, s) \approx E_{\text{max}} \sin^2(ks)$$

 π or the $1/2\beta\lambda$ mode

Transit factor (for this example):
$$\langle E \rangle = \frac{1}{\lambda_{RF}} \int_{0}^{\lambda_{RF}} E(\frac{s}{v_{\text{particle}}}, s) \, ds \approx \frac{1}{2} E_{\text{max}}$$

