



Solenoid vs. Strong Focusing



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If the solenoids field was perpendicular to the particle's motion,

its bending radius would be $\rho_z = \frac{p}{qB_z}$

$$\ddot{r} = -\left(\frac{qB_z}{2m\gamma}\right)^2 r = -\frac{qv_z}{m\gamma} B_z \frac{r}{4\rho_z}$$

Solenoid focusing is weak compared to the deflections created by a transverse magnetic field.

Transverse fields: $\vec{B} = B_x \vec{e}_x + B_y \vec{e}_y$

$$m\gamma \begin{pmatrix} \ddot{x} \\ \ddot{y} \\ \ddot{z} \end{pmatrix} = q \begin{pmatrix} \dot{x} \\ \dot{y} \\ \dot{z} \end{pmatrix} \times \begin{pmatrix} B_x \\ B_y \\ 0 \end{pmatrix} \Rightarrow \begin{pmatrix} \ddot{x} \\ \ddot{y} \end{pmatrix} = \frac{qv_z}{m\gamma} \begin{pmatrix} -B_y \\ B_x \end{pmatrix} \quad \text{Strong focusing}$$

Weak focusing < Strong focusing by about r/ρ



Solenoid Focusing



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Solenoid magnets are used in detectors for particle identification via $\rho = \frac{p}{qB}$

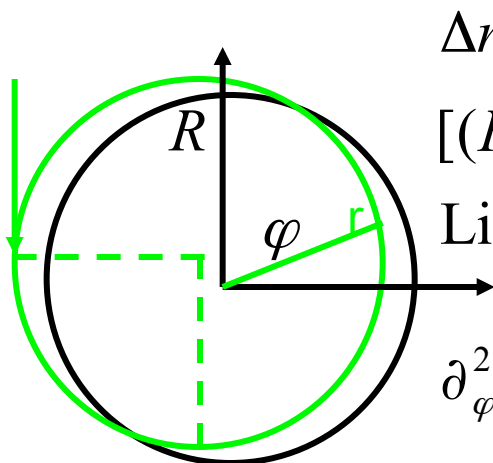
The solenoid's rotation $\dot{\varphi} = -\frac{qB_z}{2m\gamma}$ of the beam is often compensated by a reversed solenoid called compensator.

Solenoid or Weak Focusing:

Solenoids are also used to focus low γ beams:

$$\ddot{w} = -\left(\frac{qB_z}{2m\gamma}\right)^2 w$$

Weak focusing from natural ring focusing:



$$\Delta r = r - R$$

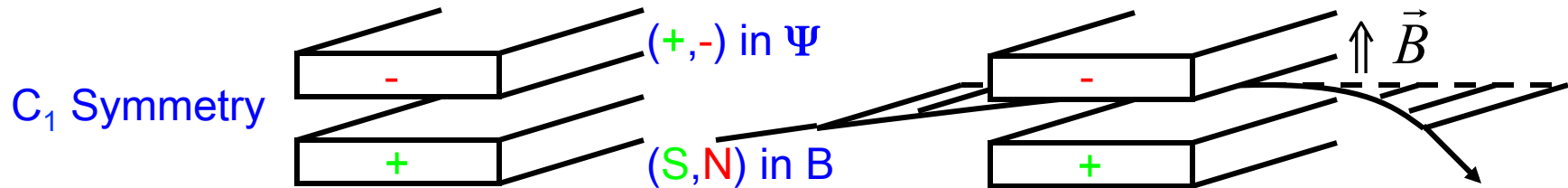
$$[(R + \Delta r) \cos \varphi - \Delta x_0]^2 + [(R + \Delta r) \sin \varphi - \Delta y_0]^2 = R^2$$

$$\text{Linearization in } \Delta: \Delta r = (\cos \varphi \Delta x_0 + \sin \varphi \Delta y_0)$$

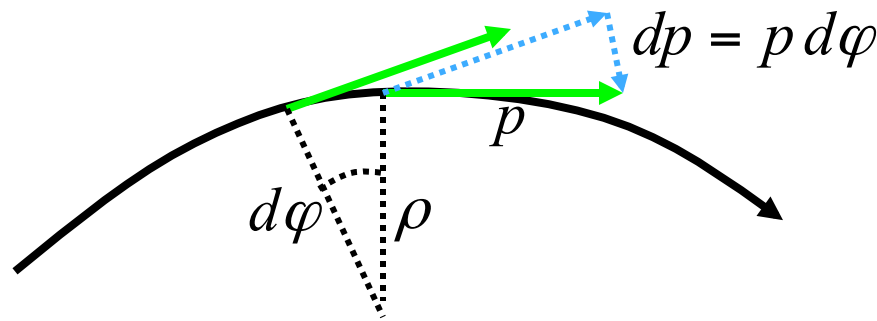
$$\partial_\varphi^2 \Delta r = -\Delta r \Rightarrow \Delta \ddot{r} = -\dot{\varphi}^2 \Delta r = -\left(\frac{v}{\rho}\right)^2 \Delta r = -\left(\frac{qB}{m\gamma}\right)^2 \Delta r$$



$$\psi = \Psi_1 \operatorname{Im}\{x - iy\} = -\Psi_1 \cdot y \Rightarrow \vec{B} = -\vec{\nabla} \psi = \Psi_1 \vec{e}_y \quad \text{Equipotential } y = \text{const.}$$



Dipole magnets are used for steering the beams direction



$$\frac{d\vec{p}}{dt} = q \vec{v} \times \vec{B} \Rightarrow \frac{dp}{dt} = qvB_{\perp} \Rightarrow \rho = \frac{dl}{d\varphi} = \frac{vdt}{dp/p} = \frac{p}{qB_{\perp}}$$

Bending radius: $\rho = \frac{p}{qB}$

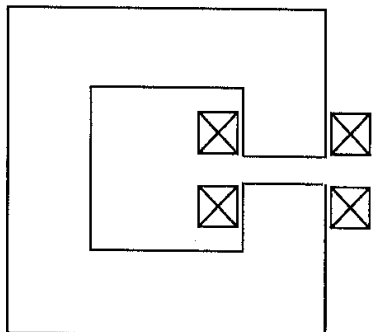


Different Dipoles

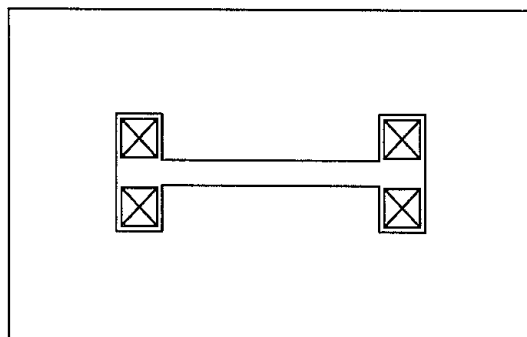


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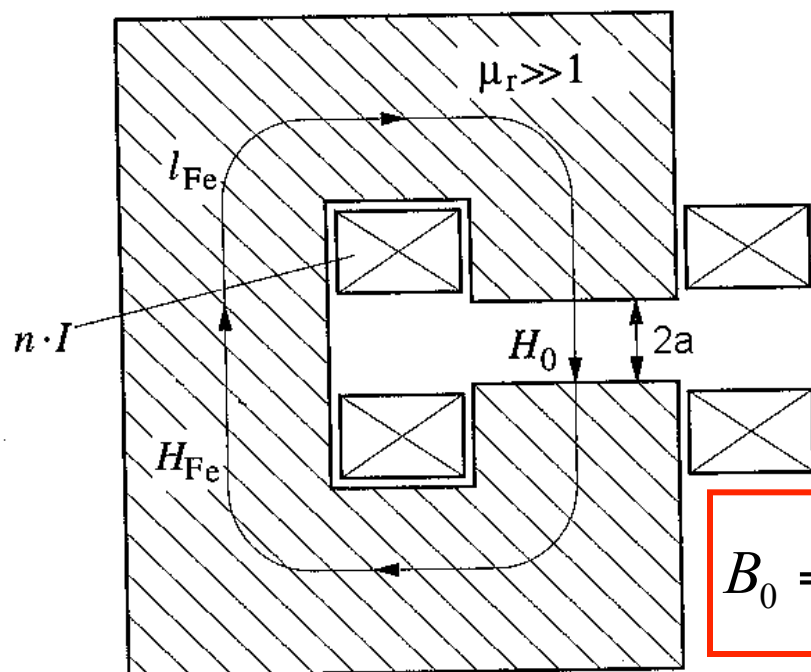
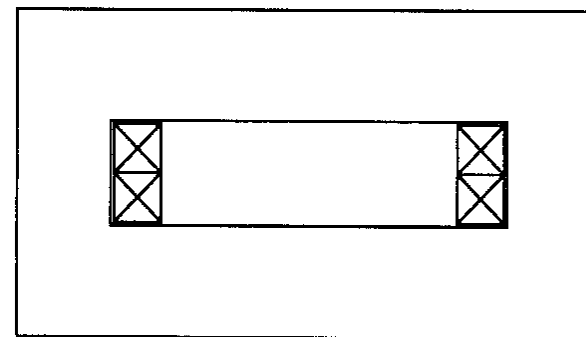
C-shape magnet:



H-shape magnet:



Window frame magnet:



$$\vec{B}_{\perp}(\text{out}) = \vec{B}_{\perp}(\text{in})$$

$$\vec{H}_{\perp}(\text{out}) = \mu_r \vec{H}_{\perp}(\text{in})$$

$$2nI = \oint \vec{H} \cdot d\vec{s} = H_{Fe} l_{Fe} + H_0 2a$$

$$= \frac{1}{\mu_r} H_0 l_{Fe} + H_0 2a \approx H_0 2a$$

$$B_0 = \mu_0 \frac{nI}{a}$$

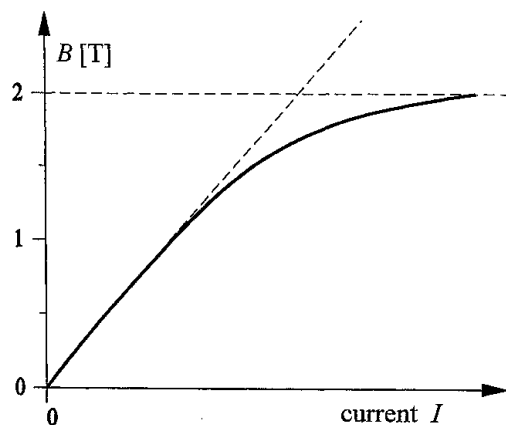
Dipole strength: $\frac{1}{\rho} = \frac{q\mu_0}{p} \frac{nI}{a}$



Dipole Fields



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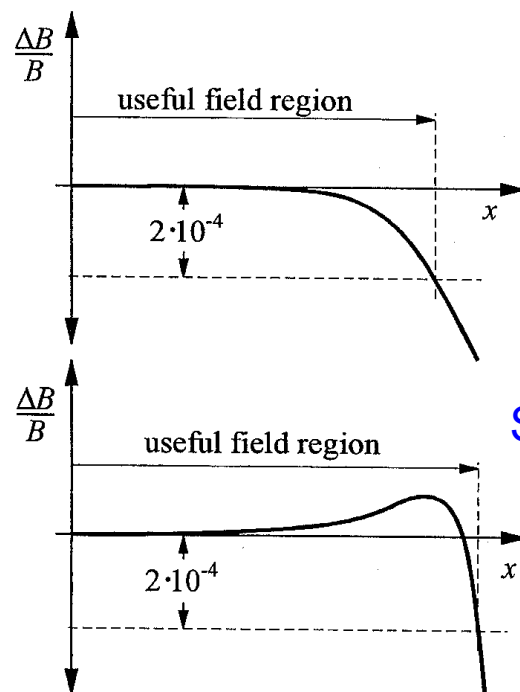
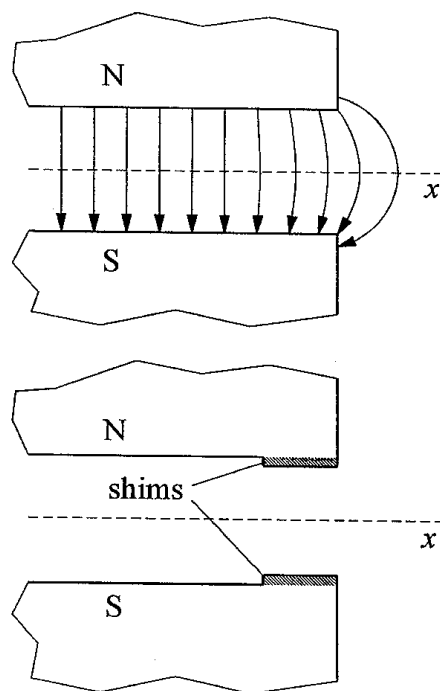


$B = 2 \text{ T}$: Typical limit, since the field becomes dominated by the coils, not the iron.

Limiting j for Cu is about 100 A/mm^2

$B < 1.5 \text{ T}$: Typically used region

$B < 1 \text{ T}$: Region in which $B_0 = \mu_0 \frac{nI}{a}$



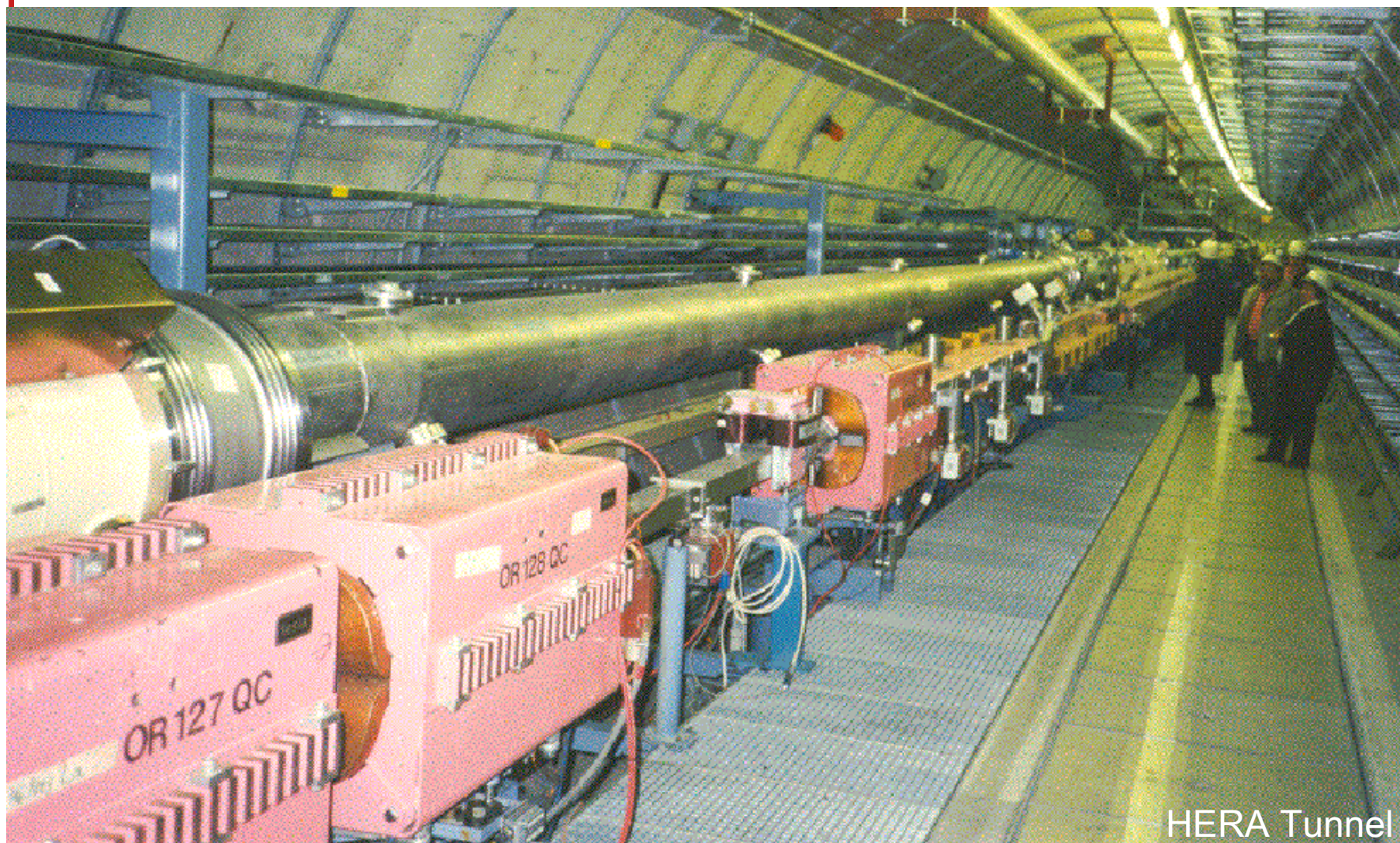
Shims reduce the space that is open to the beam, but they also reduce the fringe field region.



Where is the vertical Dipole?



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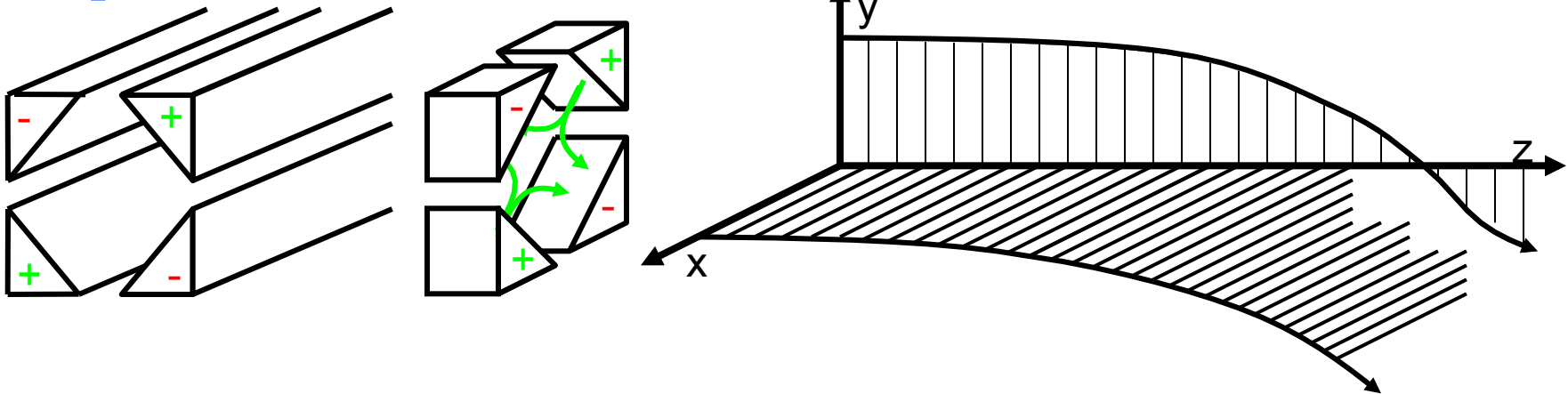


HERA Tunnel



$$\psi = \Psi_2 \operatorname{Im}\{(x - iy)^2\} = -\Psi_2 \cdot 2xy \Rightarrow \vec{B} = -\vec{\nabla}\psi = \Psi_2 2 \begin{pmatrix} y \\ x \end{pmatrix}$$

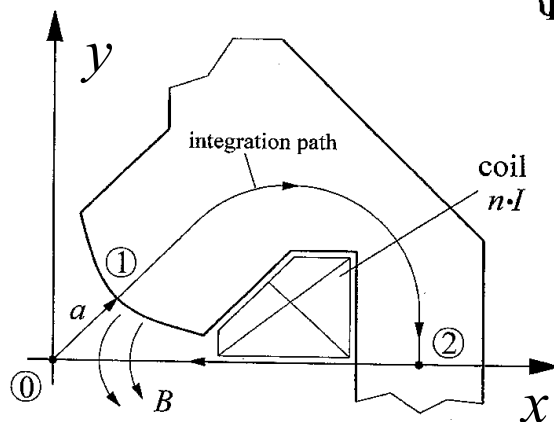
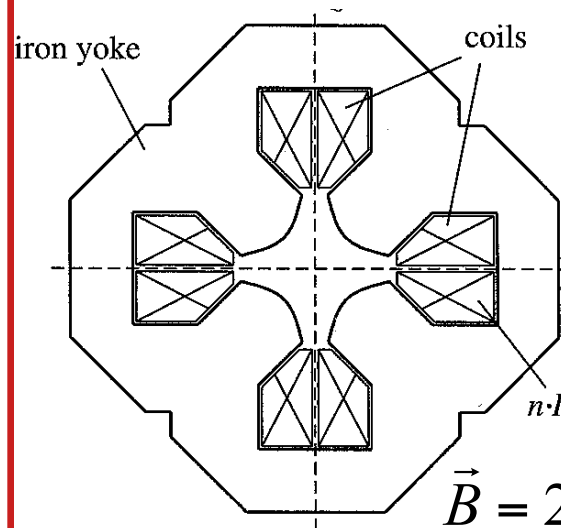
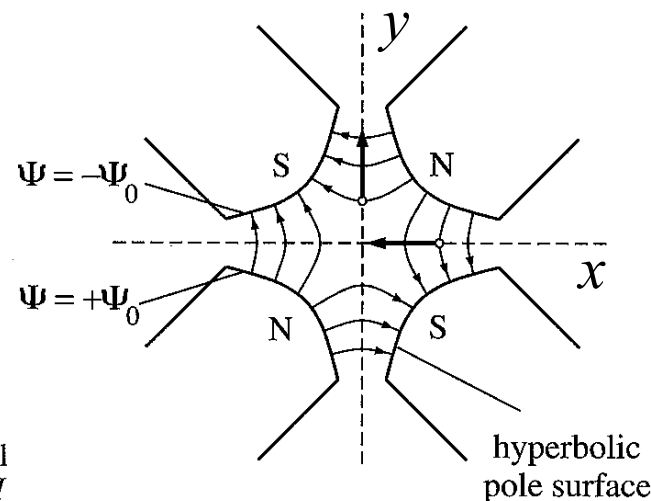
C_2 Symmetry



In a **quadrupole** particles are focused in one plane and defocused in the other plane. Other modes of **strong focusing** are not possible.

Quadrupole Fields

$$\psi = -\Psi_2 \cdot 2xy \Rightarrow \text{Equipotential: } x = \frac{\text{const.}}{y}$$



$$\vec{B} = 2\Psi_2 \begin{pmatrix} y \\ x \end{pmatrix} \Rightarrow \vec{B}(0 \mapsto 1) = 2\Psi_2 r \vec{e}_r$$

Quadrupole strength:

$$nI = \oint \vec{H} \cdot d\vec{s} \approx \int_0^a H_r dr = \Psi_2 \frac{a^2}{\mu_0}$$

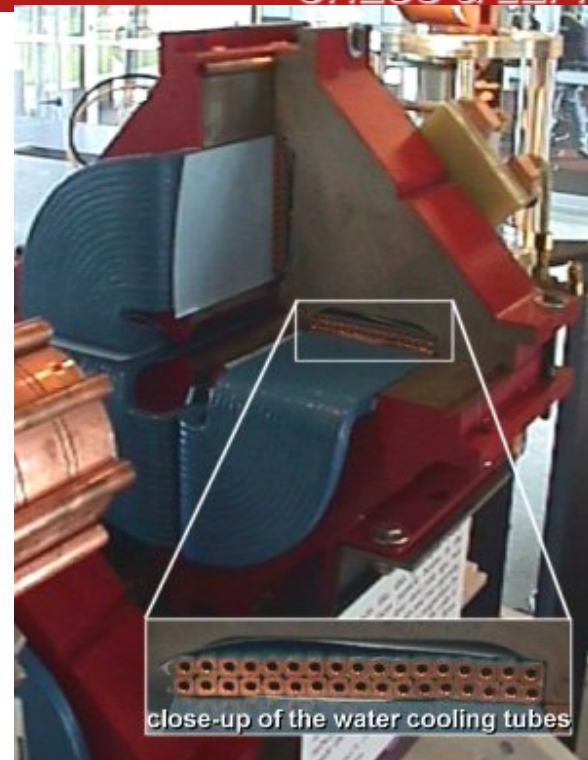
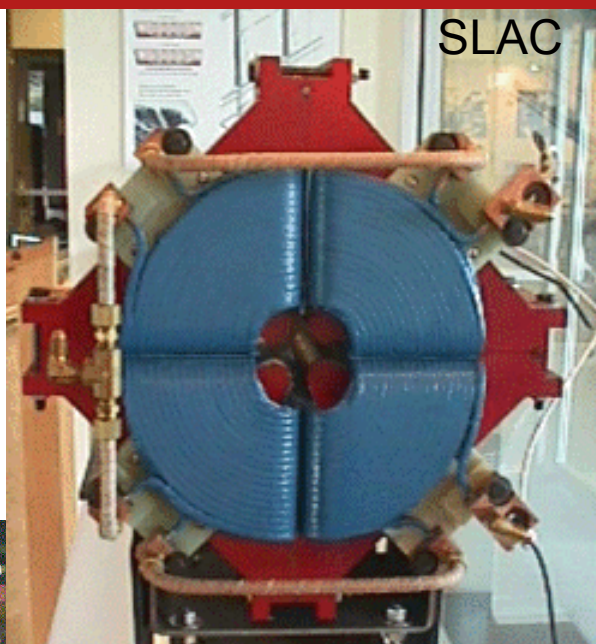
$$k_1 = \frac{q}{p} \partial_x B_y \Big|_0 = \frac{q\mu_0}{p} \frac{2nI}{a^2}$$



Real Quadrupoles



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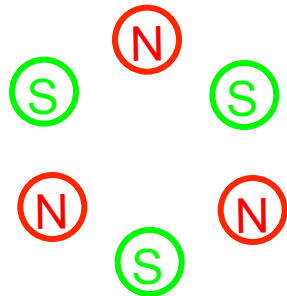


The coils show that this is an upright quadrupole not a rotated or skew quadrupole.



$$\psi = \Psi_3 \operatorname{Im}\{(x - iy)^3\} = \Psi_3 \cdot (y^3 - 3x^2y) \Rightarrow \vec{B} = -\vec{\nabla}\psi = \Psi_3 3 \begin{pmatrix} 2xy \\ x^2 - y^2 \end{pmatrix}$$

C_3 Symmetry



i) Sextupole fields hardly influence the particles close to the center, where one can linearize in x and y .

ii) In linear approximation a by Δx shifted sextupole has a quadrupole field.

$$\vec{B} = -\vec{\nabla}\psi = \Psi_3 3 \begin{pmatrix} 2xy \\ x^2 - y^2 \end{pmatrix}$$

iii) When Δx depends on the energy, one can build an **energy dependent quadrupole**.

$$x \mapsto \Delta x + x$$

$$\vec{B} \approx \Psi_3 3 \begin{pmatrix} 2xy \\ x^2 - y^2 \end{pmatrix} + 6\Psi_3 \Delta x \begin{pmatrix} y \\ x \end{pmatrix} + O(\Delta x^2)$$

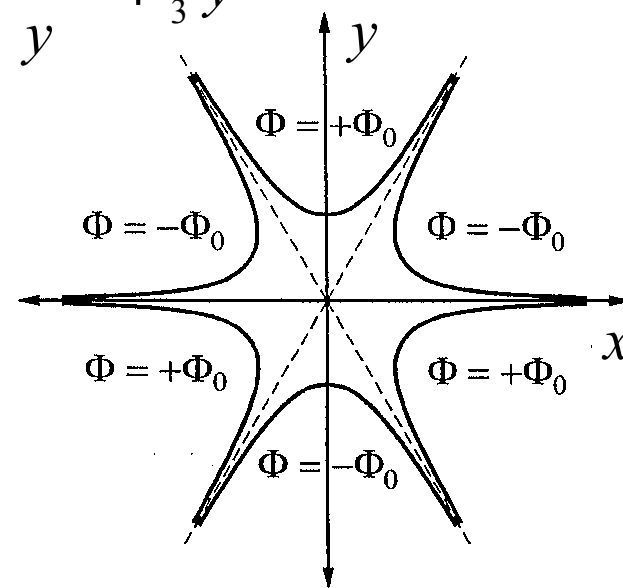
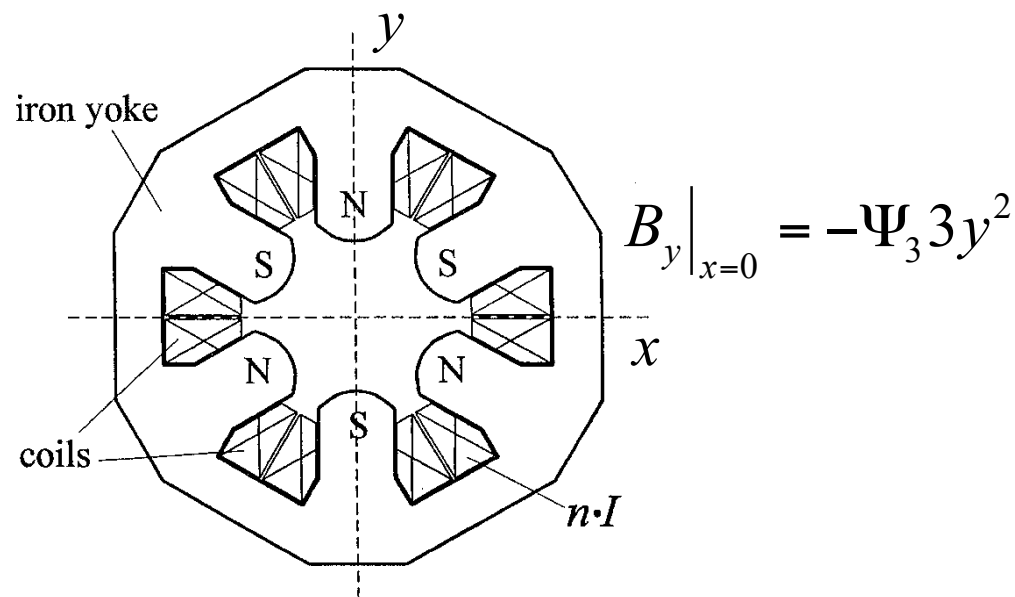


Sextupole Fields



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$$\psi = \Psi_2 \cdot (y^3 - 3x^2y) \Rightarrow \text{Equipotential: } x = \sqrt{\frac{\text{const.}}{y} + \frac{1}{3}y^2}$$



Quadrupole strength:

$$nI = \oint \vec{H} \cdot d\vec{s} \approx \int_0^a H_r dr = \Psi_3 \frac{a^3}{\mu_0}$$

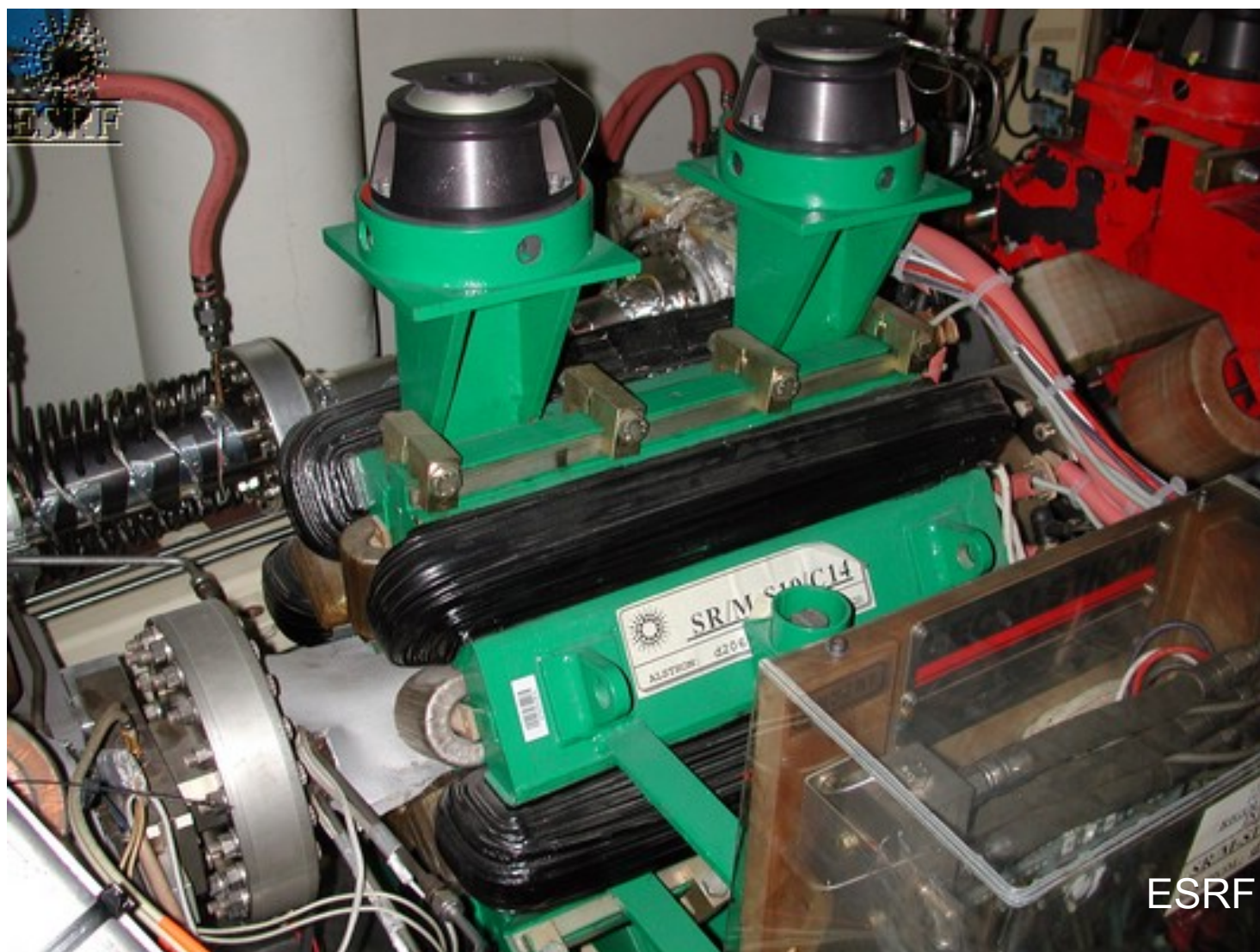
$$k_2 = \frac{q}{p} \partial_x^2 B_y \Big|_0 = \frac{q\mu_0}{p} \frac{6nI}{a^3}$$



Real Sextupoles



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ESRF



The CESR Tunnel



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Higher order Multipoles



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$$\psi = \Psi_n \operatorname{Im}\{(x - iy)^n\} = \Psi_n \cdot (\dots - i n x^{n-1} y) \Rightarrow \vec{B}(y=0) = \Psi_n n \begin{pmatrix} 0 \\ x^{n-1} \end{pmatrix}$$

Multipole strength: $k_n = \frac{q}{p} \partial_x^n B_y \Big|_{x,y=0} = \frac{q}{p} \Psi_{n+1} (n+1)! \text{ units: } \frac{1}{\text{m}^{n+1}}$

p/q is also called $B\rho$ and used to describe the energy of multiply charge ions

Names: dipole, quadrupole, sextupole, octupole, decapole, duodecapole, ...

Higher order multipoles come from

- Field errors in magnets
- Magnetized materials
- From multipole magnets that compensate such erroneous fields
- To compensate nonlinear effects of other magnets
- To stabilize the motion of many particle systems
- To stabilize the nonlinear motion of individual particles