

The effect of dipole and quadrupole HOMs on the performance of the Phase II ERL

Outline:

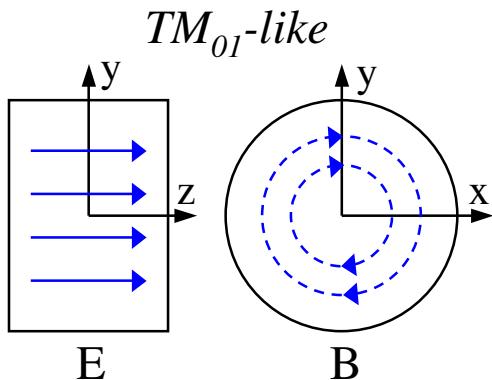
- Beam induced voltage: losses and kicks
- Beam breakup for quadrupole HOMs
- BBU statistics for “ERL in CESR tunnel”
- Discussion

Introduction

The “Big Four”

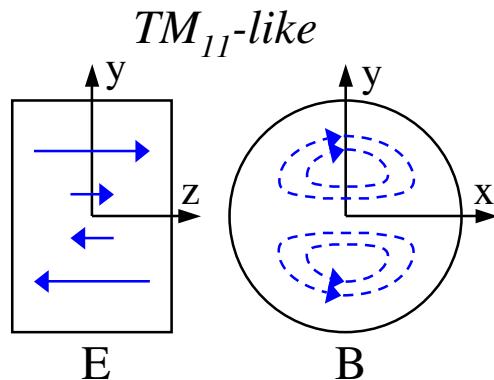
- Superposition
- Concept of normal modes
- Conservation of energy
- Causality

monopole ($m = 0$)



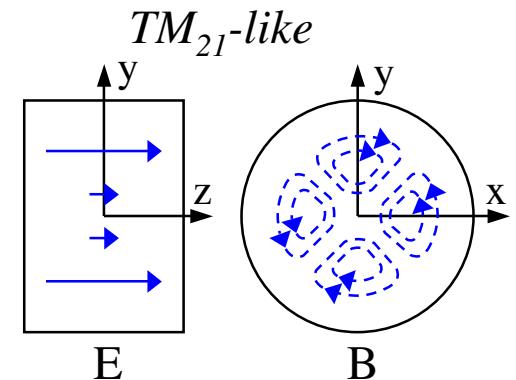
high losses, no kick

dipole ($m = 1$)



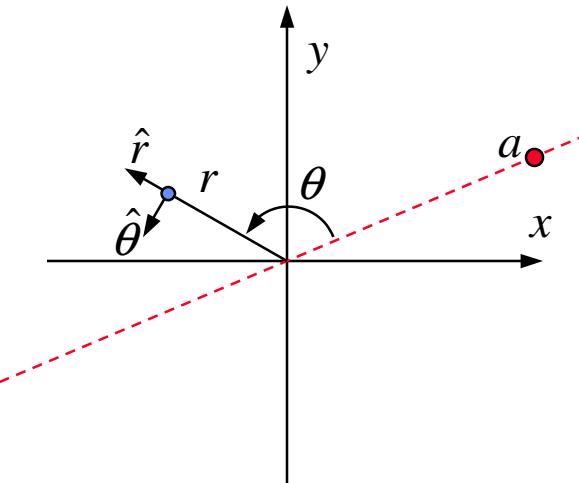
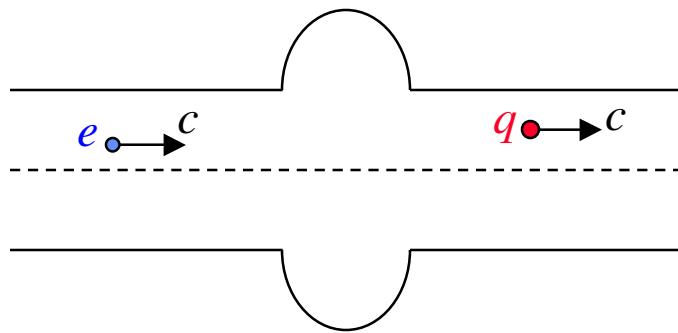
kick and losses when beam is not centered

quadrupole ($m = 2$)



kick, coupling and losses when beam is not centered

Wake formalism for point charge



$$\int_{-L/2}^{L/2} \vec{F}_\perp ds = -eqa^m W_m(z) mr^{m-1} (\hat{r} \cos m\theta - \hat{\theta} \sin m\theta)$$

$$\int_{-L/2}^{L/2} F_\parallel ds = -eqa^m W'_m(z) r^m \cos m\theta$$

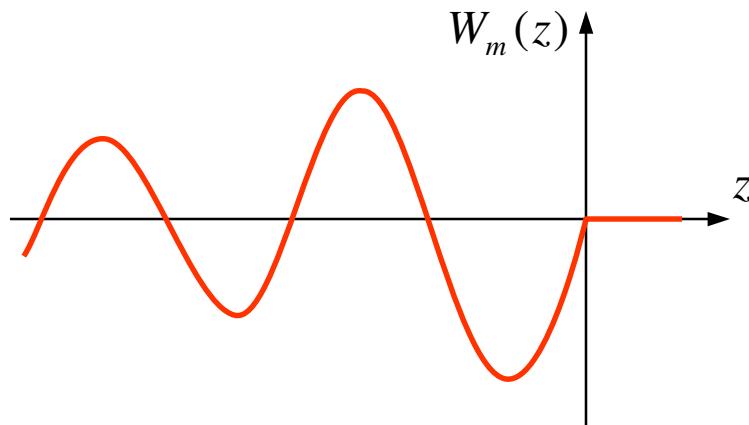
The above can be combined to say (Panofsky-Wenzel): $\nabla_\perp \int_{-L/2}^{L/2} F_\parallel ds = \frac{\partial}{\partial z} \int_{-L/2}^{L/2} \vec{F}_\perp ds$

Wake function for resonant mode

A mode characterized by $\left(\frac{R}{Q}\right)_m, Q, \omega$

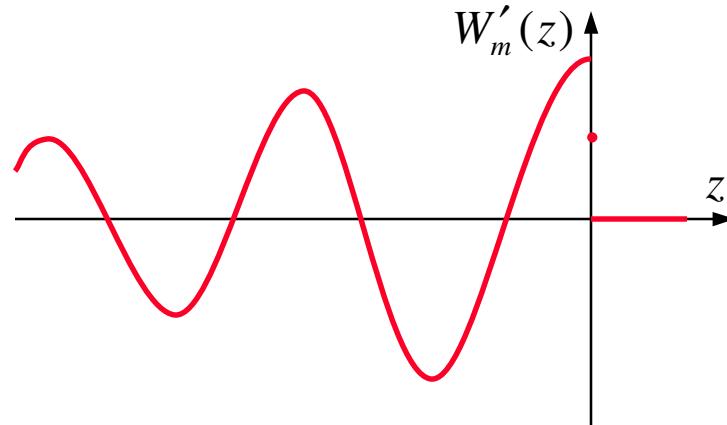
transverse

$$W_m = 2 \frac{ck_m}{\omega} e^{\frac{kz}{2Q}} \sin kz$$



longitudinal

$$W'_m \approx 2k_m e^{\frac{kz}{2Q}} \cos kz$$

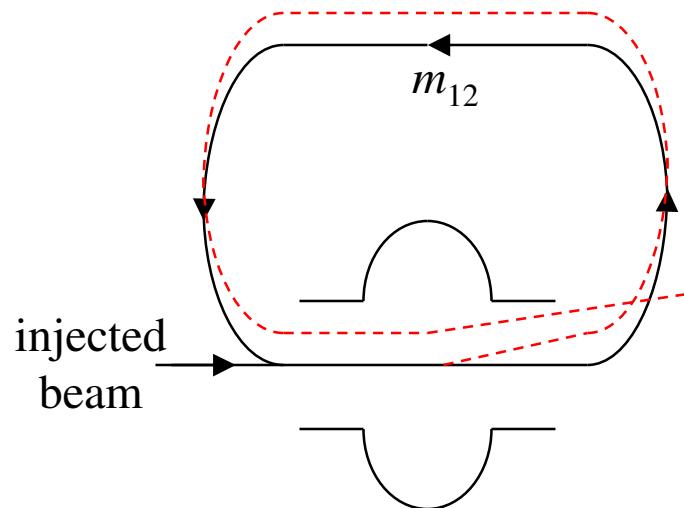


loss factor $k_m = \frac{\omega}{4} \left(\frac{R}{Q} \right)_m$

$\left(\frac{R}{Q} \right)_m$ in units of $\Omega \cdot m^{-2m}$

$k = \omega/c$ is mode's wave number

Beam breakup

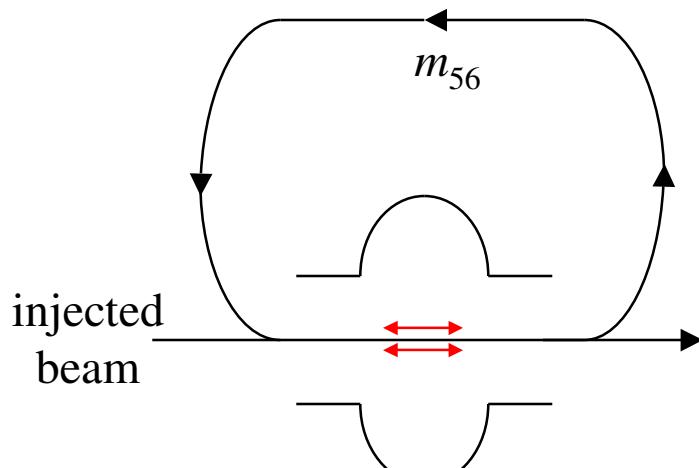


Transverse

2nd pass
deflected
beam

$$x^{(2)}(t) = \frac{e}{c} m_{12} V_x(t - t_r)$$

$$V_{1,x}(t) = - \int_{-\infty}^t W_1(t-t') [I^{(1)}(t') x^{(1)}(t') + I^{(2)}(t') x^{(2)}(t')] dt'$$



Longitudinal

2nd pass
modulated
beam

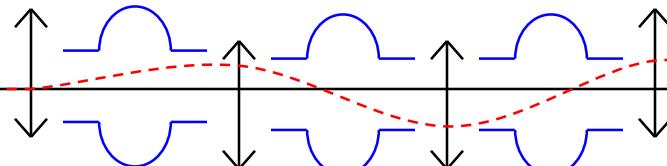
$$I_{th} = \frac{2Ec}{e(R/Q)_\lambda Q_\lambda \omega_\lambda} \frac{1}{R_{56}}$$

$$I_{th} = \frac{2Ec^2}{e(R/Q)_\lambda Q_\lambda \omega_\lambda} \frac{1}{R_{12}}$$

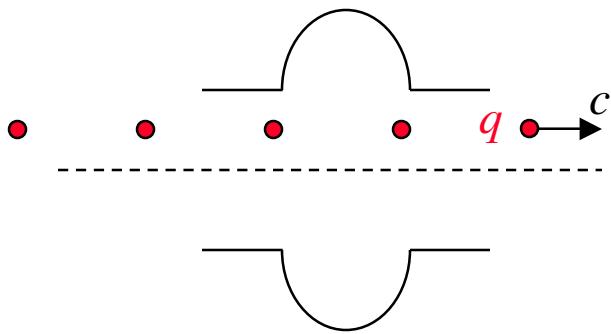
Transient and steady case parts of induced voltage

- Equation for transient part reproduces integral equation for BBU for aligned cavity
- Steady case part of beam induced voltage generates displaced orbit

$$-\frac{e}{c} m_{12} d \frac{\int_{-\infty}^t W_1(t-t') [I^{(2)}(t') + I^{(2)}(t'+t_r)] dt'}{1 + \frac{e}{c} m_{12} \int_{-\infty}^t W_1(t-t') I^{(2)}(t') dt'}$$

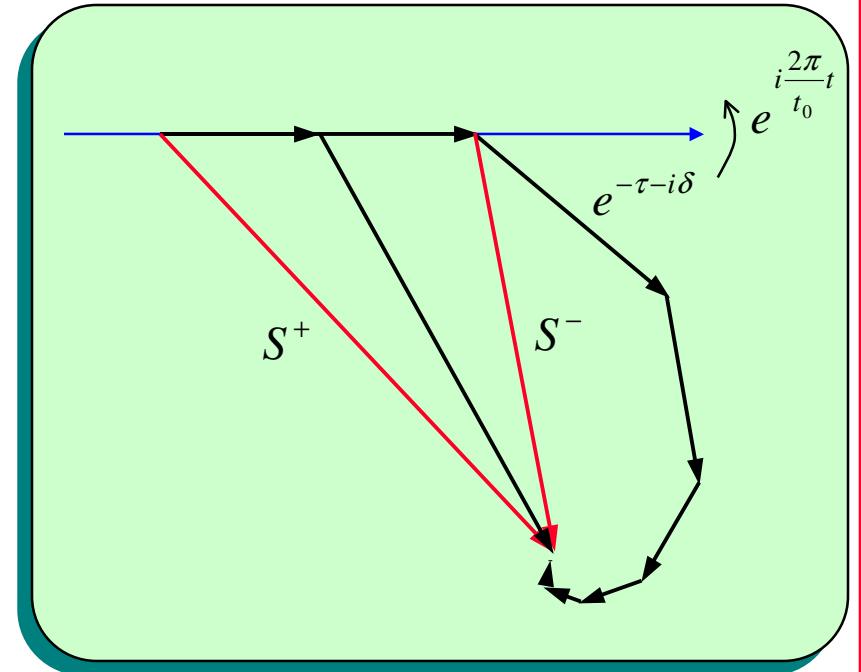


Beam induced voltage from infinite bunch train



$$V_{\parallel,m} = -qd^{2m} \frac{\omega}{2} \left(\frac{R}{Q} \right)_m \Re \left\{ \sum_{n=0}^{\infty} e^{-n\frac{\omega t_0}{2Q}} e^{-in\omega t_0} - \frac{1}{2} \right\}$$

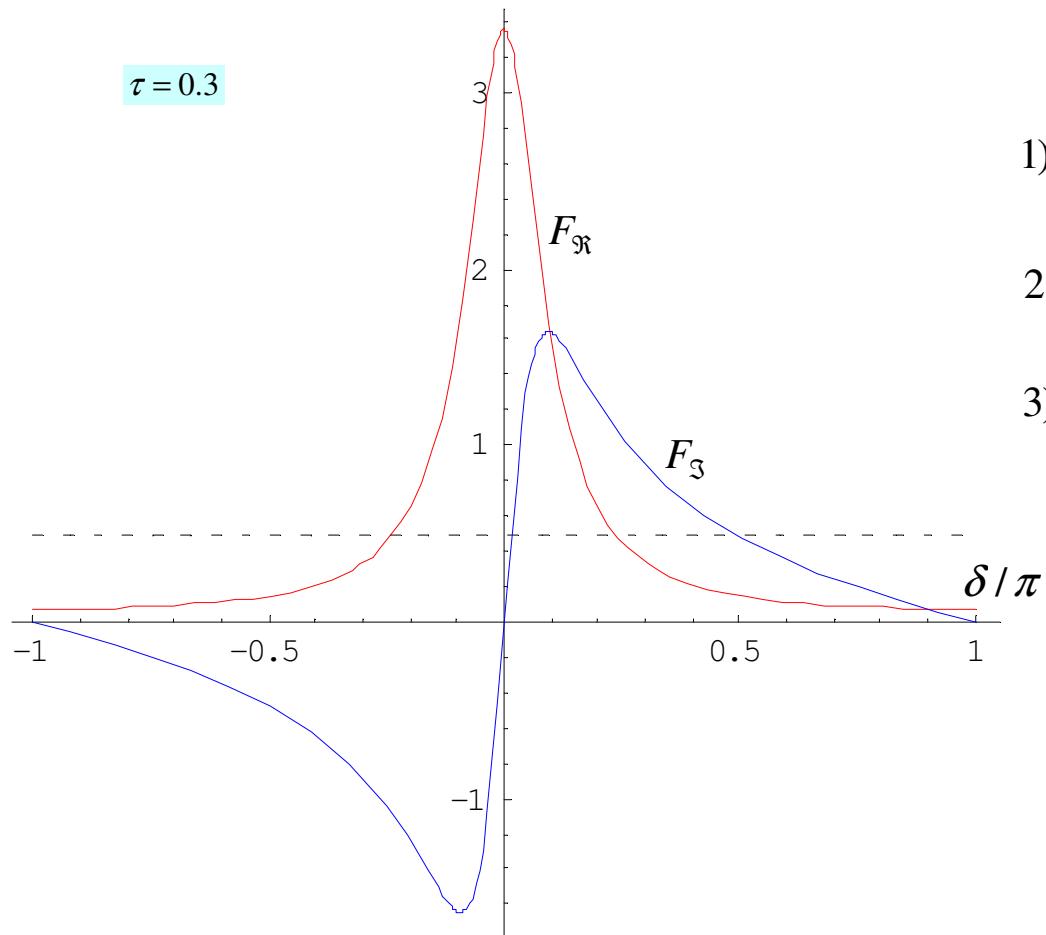
$$V_{\perp,m} = -\frac{m}{2} qd^{2m-1} c \left(\frac{R}{Q} \right)_m \Im \left\{ \sum_{n=0}^{\infty} e^{-n\frac{\omega t_0}{2Q}} e^{-in\omega t_0} - \frac{1}{2} \right\}$$



$$\sum_{n=0}^{\infty} e^{-n\frac{\omega t_0}{2Q}} e^{-in\omega t_0} - \frac{1}{2} = S^+(\tau, \delta) - \frac{1}{2} = F_{\Re}(\tau, \delta) - iF_{\Im}(\tau, \delta)$$

$$\tau = \frac{\omega t_0}{2Q}, \quad \delta = \omega t_0$$

Resonance function



Real part

$$1) \quad \frac{1}{2\pi} \int_{-\pi}^{\pi} F_{\Re}(\tau, \delta) d\delta = \frac{1}{2}$$

$$2) \quad \max F_{\Re}(\tau, \delta) = F_{\Re}(\tau, 0) = \frac{1}{\tau}$$

$$3) \quad F_{\Re} = \frac{1}{2} \text{ when } \delta = \pm\sqrt{2\tau} = \pm\sqrt{\frac{\omega t_0}{Q}}$$

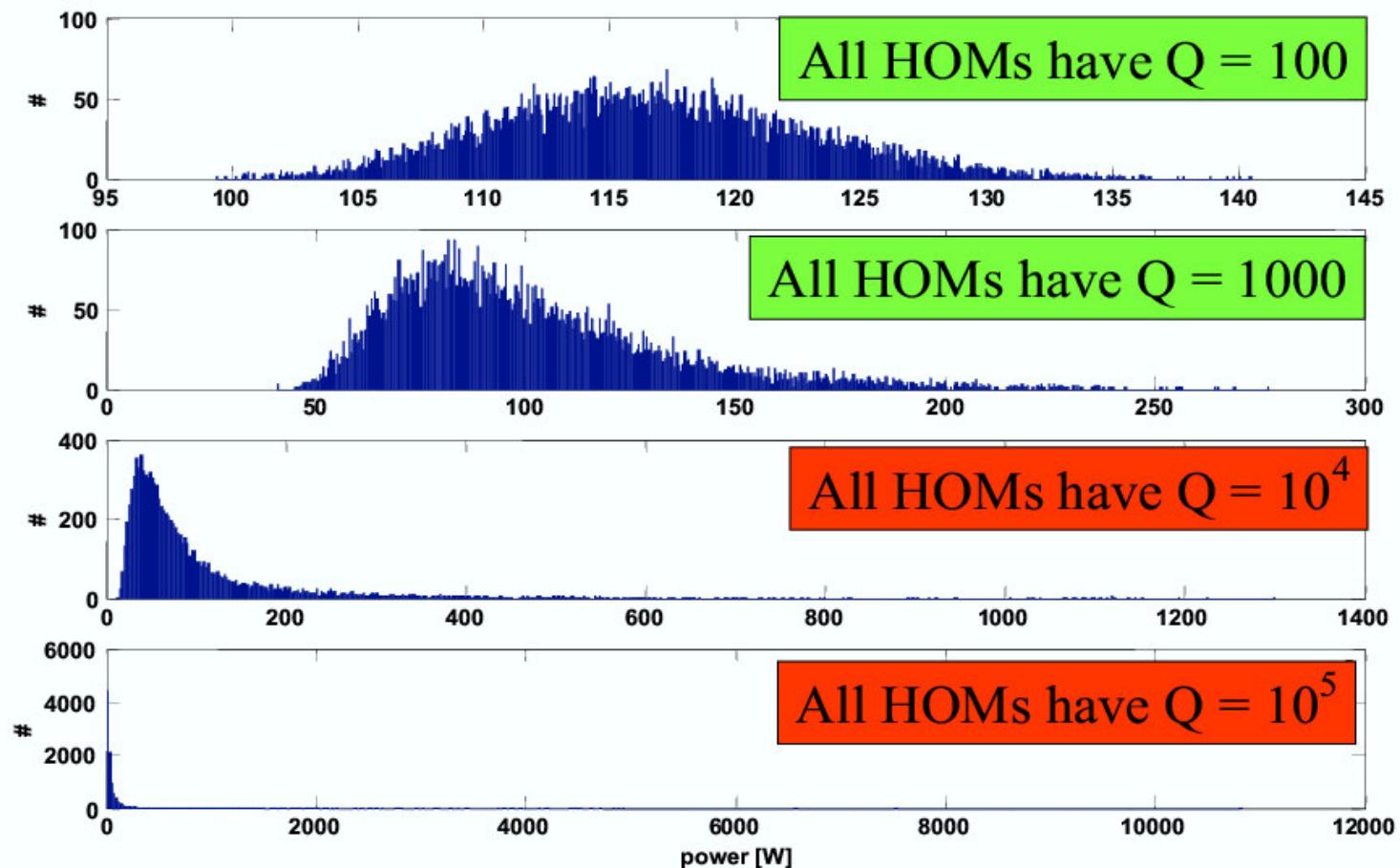
Imaginary part

$$1) \quad \frac{1}{2\pi} \int_{-\pi}^{\pi} F_{\Im}(\tau, \delta) d\delta = 0$$

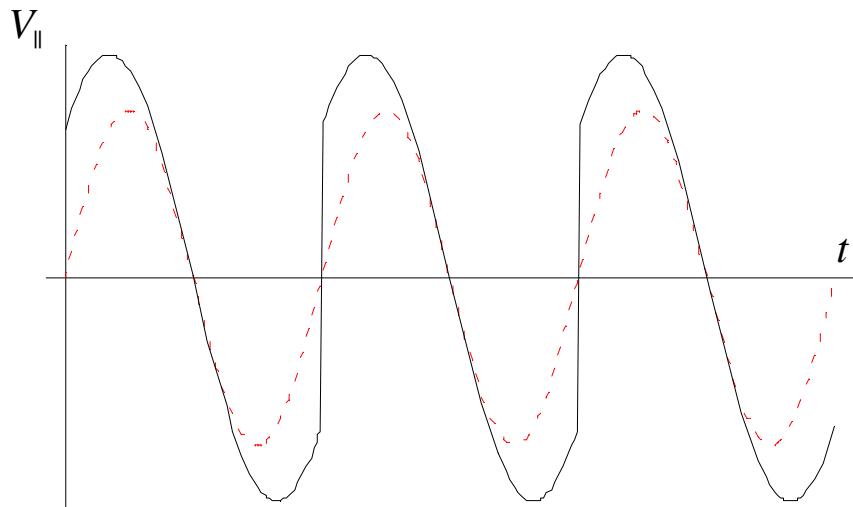
$$2) \quad \max |F_{\Im}(\tau, \delta)| \approx |F_{\Im}(\tau, \pm \frac{\tau}{\sqrt{2}})| \approx \frac{1}{2\tau}$$

$$3) \quad |F_{\Im}| \geq \frac{1}{2} \text{ when } -\frac{\pi}{2} \leq \delta \leq \frac{\pi}{2}$$

*A Simple Model: 1000 Monopoles with random f's
Total HOM Monopole Power for random Sets of Frequencies*



Beam induced voltages

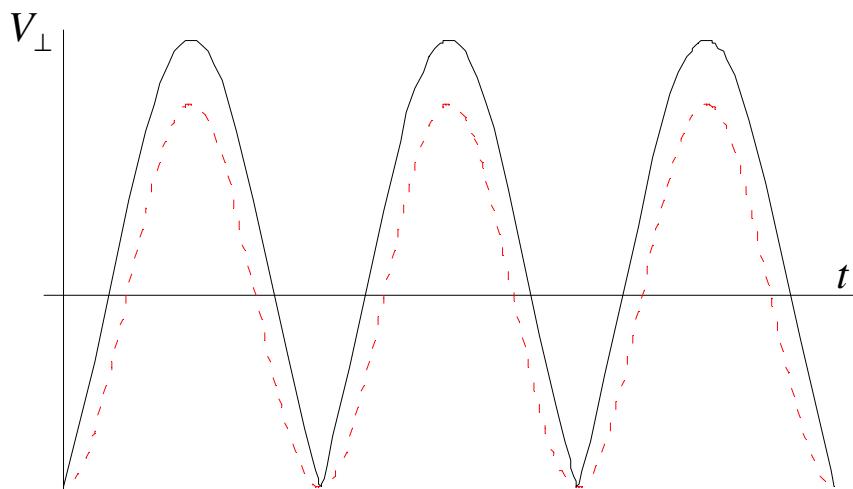


Typical

$$\langle V_{\parallel,m} \rangle = -k_m q d^{2m} = -q d^{2m} \frac{\omega}{4} \left(\frac{R}{Q} \right)_m$$

At resonance

$$|V_{\parallel,m}| = \left(\frac{R}{Q} \right)_m Q d^{2m} I$$



Typical

$$\langle V_{\perp,m} \rangle = 0$$

At resonance

$$|V_{\perp,m}| = \frac{m}{2} \left(\frac{R}{Q} \right)_m Q k^{-1} d^{2m-1} I$$

Dipole HOMs from TESLA TDR for 100 mA ERL

$R/Q, \frac{\Omega}{\text{cm}^2}$	Q
11.21	50000
8.69	70000
15	50000
15.51	20000
6.54	50000
1.72	100000
1.75	75000
0.76	70000
1.05	100000
2.16	20000
0.5	100000
0.46	50000
0.39	25000
0.77	10000

$\max V_{\perp}, \frac{\text{V}}{\text{mm}}$
1.03E+04
1.12E+04
1.38E+04
5.69E+03
6.00E+03
3.16E+03
2.41E+03
9.76E+02
1.93E+03
7.93E+02
9.18E+02
4.22E+02
1.79E+02
1.41E+02

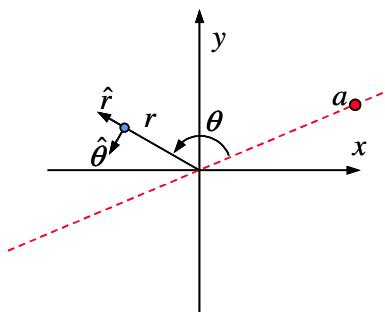
$\max V_{\parallel}, \frac{\text{V}}{\text{mm}^2}$	$\max P, \frac{\text{W}}{\text{mm}^2}$
1121	224.2
1216.6	243.32
1500	300
620.4	124.08
654	130.8
344	68.8
262.5	52.5
106.4	21.28
210	42
86.4	17.28
100	20
46	9.2
19.5	3.9
15.4	3.08

Transverse multipass beam breakup

dipole mode

$$\Delta V_{1,x} = -qW_1(z)\langle x_a \rangle$$

$$\Delta V_{1,y} = -qW_1(z)\langle y_a \rangle$$



quadrupole mode

focusing

$$\Delta \frac{\partial}{\partial x} V_{2,x} = -\Delta \frac{\partial}{\partial y} V_{2,y} = -2qW_2(z) [\langle x_a \rangle^2 - \langle y_a \rangle^2 + \langle x_a^2 \rangle - \langle y_a^2 \rangle]$$

coupling

$$\Delta \frac{\partial}{\partial y} V_{2,x} = \Delta \frac{\partial}{\partial x} V_{2,y} = -2qW_2(z) 2\langle x_a \rangle \langle y_a \rangle$$

↓

$$I_1 \approx \frac{2\omega}{em_{12}(R/Q)_1 Q}$$

$$m_{12} = \sqrt{\frac{\beta_{x,n}^{(1)}}{p^{(1)}} \frac{\beta_{x,n}^{(2)}}{p^{(2)}}} \sin \Delta \psi_x$$

Quadrupole HOM induced voltage for 100 mA ERL

focusing effect:

$$\max \frac{\partial}{\partial x} V_{2,x} = \left(\frac{R}{Q} \right)_m Q k^{-1} I d^2$$

$$\max \frac{\partial}{\partial x} V_{2,x} \sim 10^4 \frac{\text{V}}{\text{m}} \cdot d^2 (\text{mm}) \quad \text{focus.length } \frac{p}{e \frac{\partial}{\partial x} V_x} \geq 10^3 \text{ m /}d^2 (\text{mm})$$

coupling effect:

$$\Delta \epsilon_{n,c} = \frac{e}{mc^2} \sigma_x \sigma_y \frac{\partial}{\partial y} V_{2,x} \leq 1\% \cdot d^2 (\text{mm})$$

$$\epsilon_n^2 = \epsilon_{n,0}^2 + \Delta \epsilon_{n,c}^2$$

Beam breakup due to quadrupole HOMs

a) aligned cavity $\sigma_{x,y} \gg d$

$$I_{2a)} \approx \frac{2\omega}{e(R/Q)_2 Q} \frac{p^{(1)} p^{(2)}}{2m_e c [\epsilon_{x,n} \beta_{x,n}^{(1)} \beta_{x,n}^{(2)} \sin 2\Delta\psi_x + \epsilon_{y,n} \beta_{y,n}^{(1)} \beta_{y,n}^{(2)} \sin 2\Delta\psi_y]}$$

$$\frac{I_{2a)}}{I_1} \sim \frac{b^2}{4\sigma_{x,y}^2}$$

a) misaligned cavity $d \gg \sigma_{x,y}$

$$I_{2b)} \approx \frac{2\omega}{e(R/Q)_2 Q m_{12}} \frac{1}{4d^2}$$

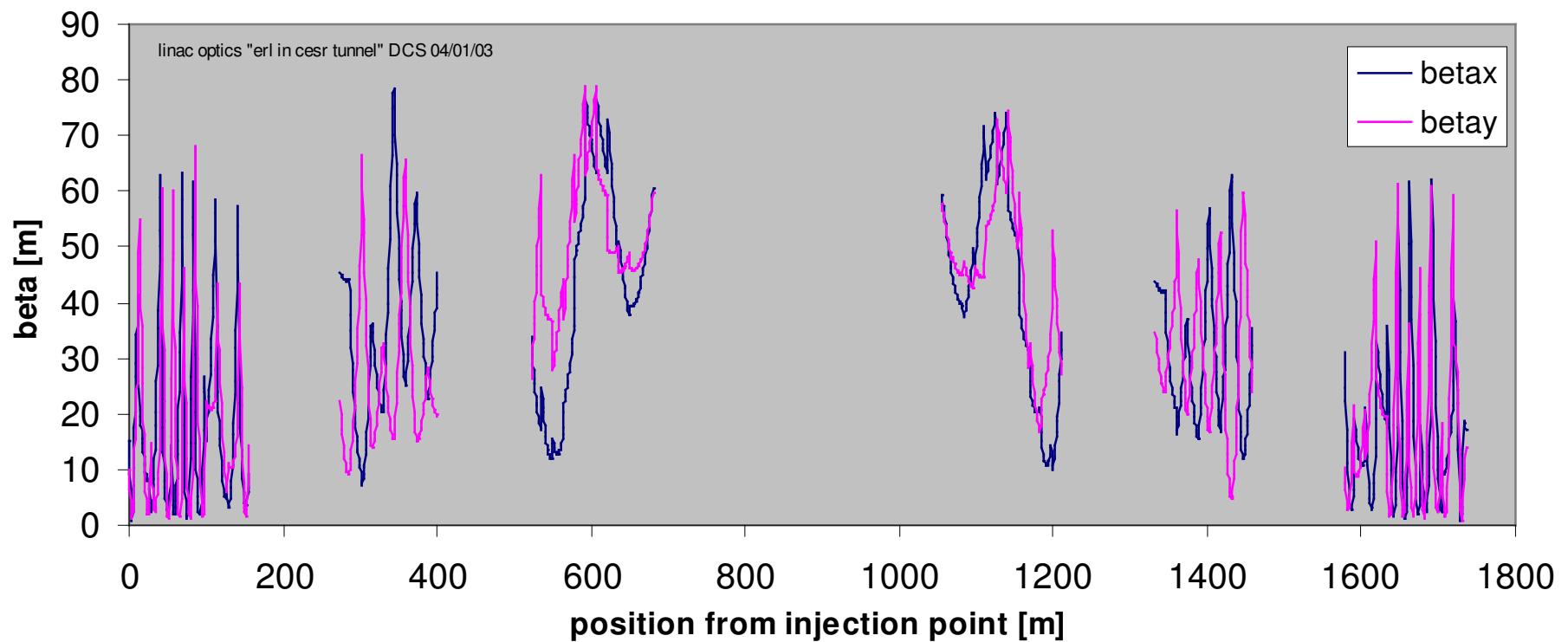
$$\frac{I_{2b)}}{I_1} \sim \frac{b^2}{4d^2}$$

bi - numeric tool for transverse BBU

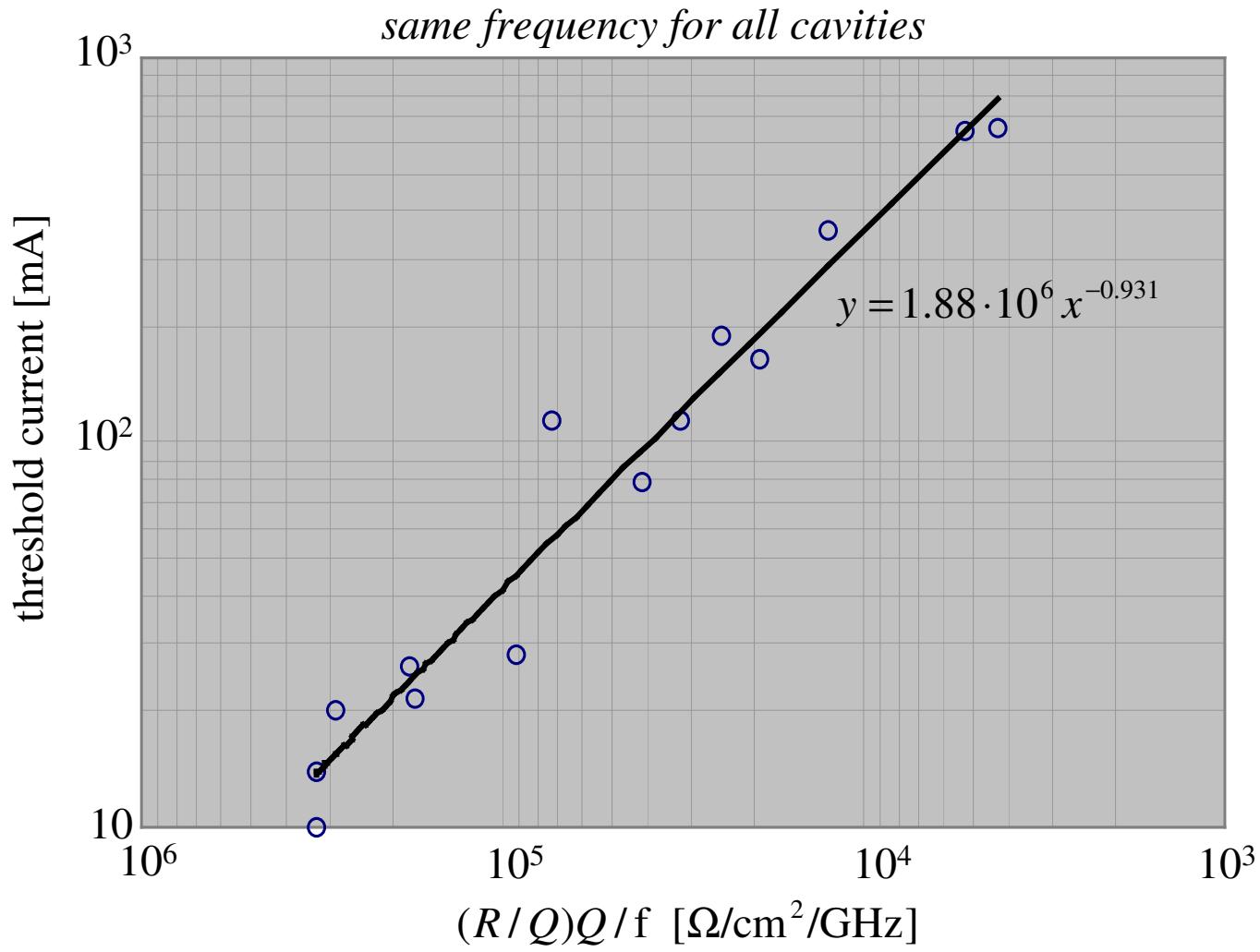
Features:

- written in C++
- cleaner algorithm than TDBBU
- faster than TDBBU (a single 5 GeV ERL run takes less than a minute; execution time is estimated to be 7-9 times faster when no coupling is present; with coupling it is estimated to be at least 4 times faster)
- allows any ERL topology
- easier to use

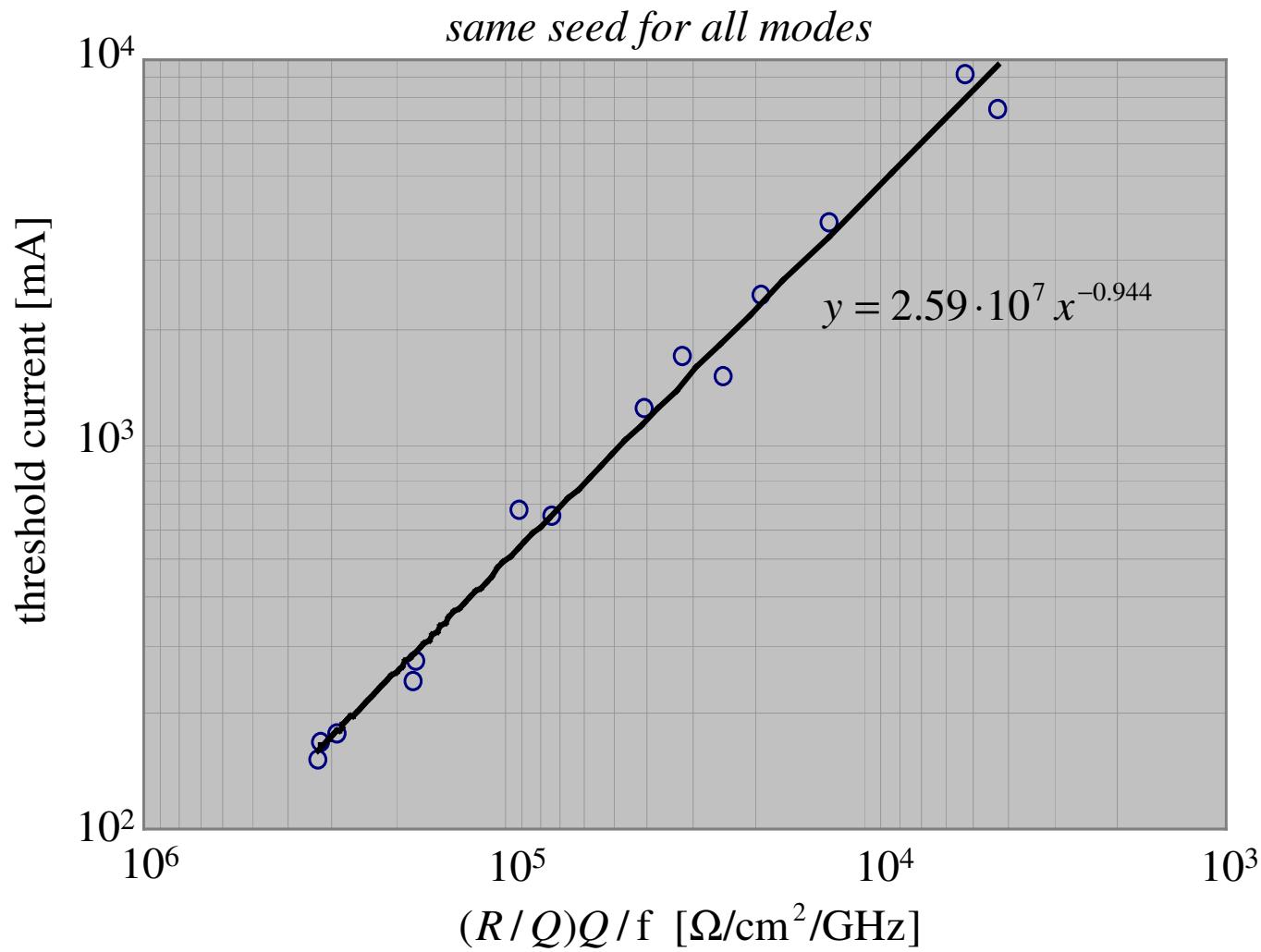
Linac optics (DCS, 04/01/03)



Worst 14 dipole modes from TESLA TDR

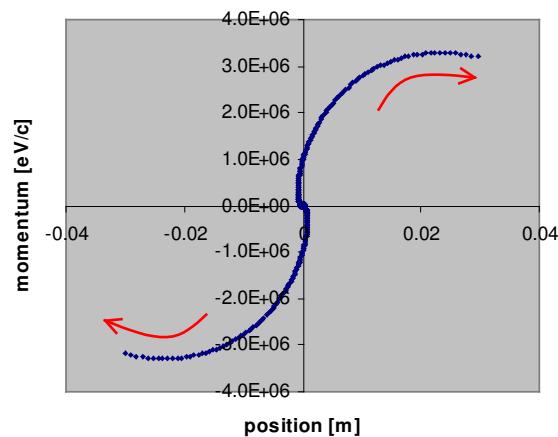


Uniform frequency spread of 10 MHz

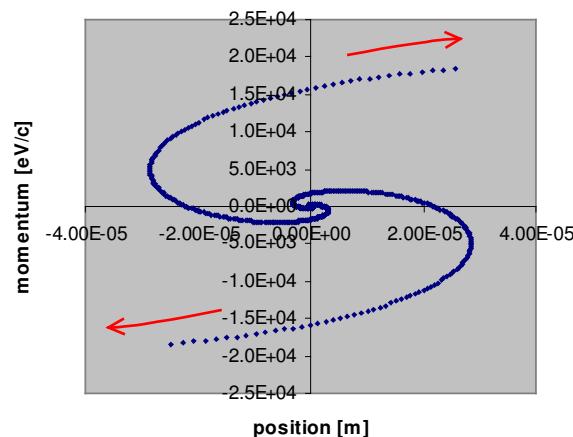


Fixed current (30 mA)

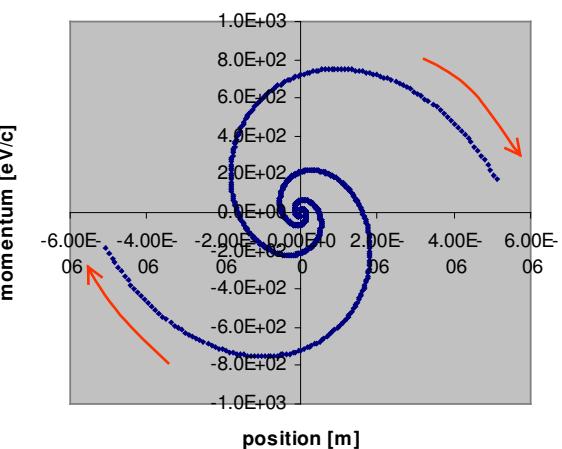
rms = 0 Hz



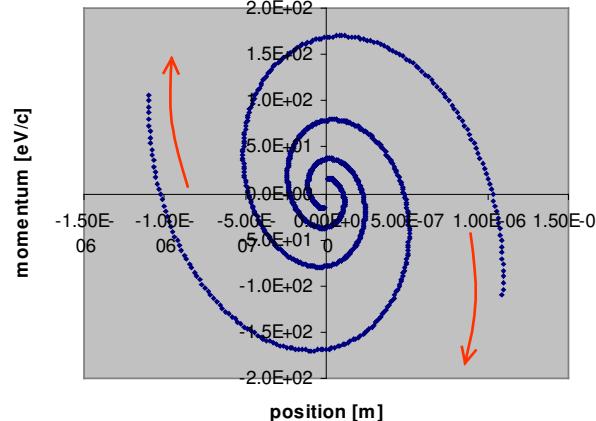
rms = 33 kHz



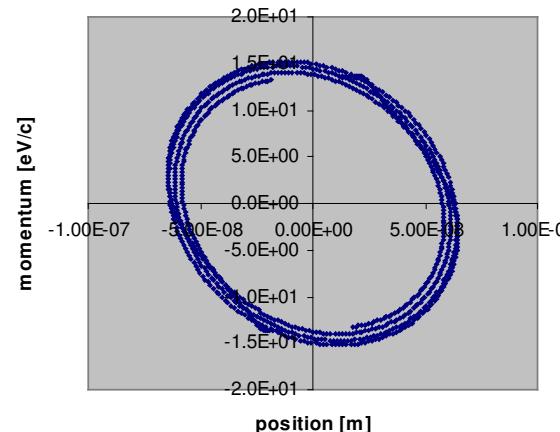
rms = 42 kHz



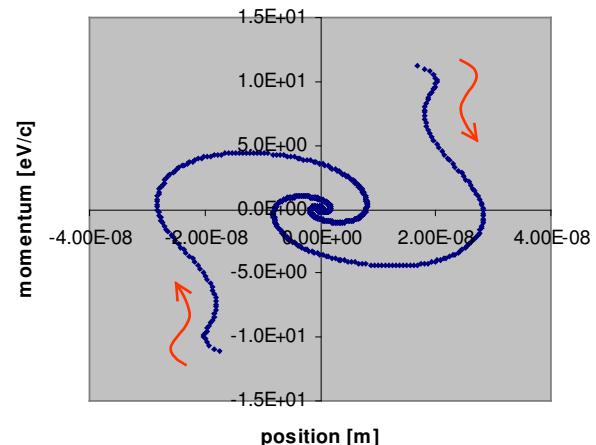
rms = 46 kHz



rms = 53 kHz



rms = 67 kHz



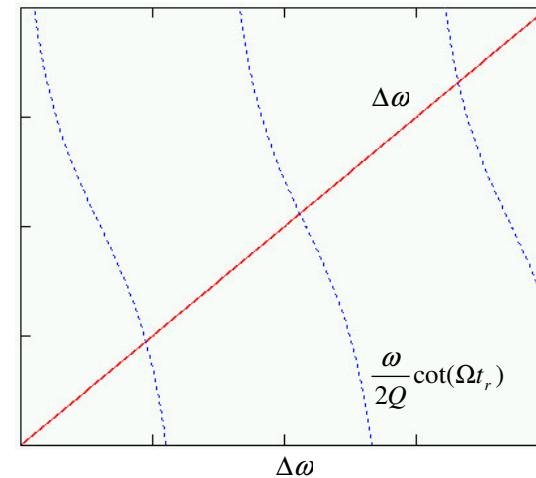
Stabilizing effect of HOM frequency spread

- Effectively decreases Q of the mode: $Q \sim \omega/\Delta\omega_{1/2}$
- Limited in its effect by $\sim \omega/Q$ of the HOM

Linearized solution for simplest case

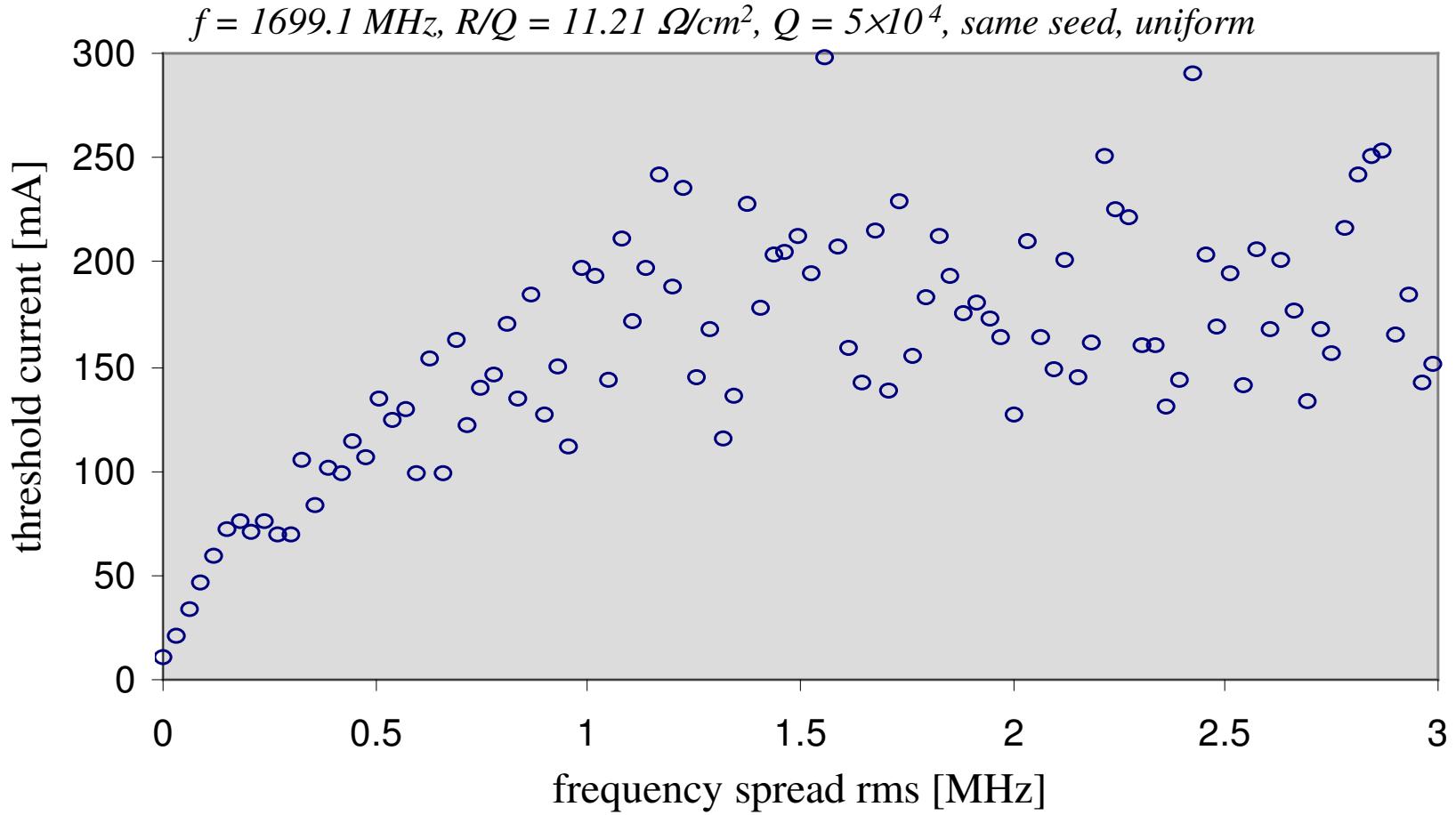
$$I_1 \approx \frac{-2c}{e(R/Q)Qkm_{12} \sin(\Omega t_r)},$$

$$\Delta\omega = \frac{\omega}{2Q} \cot(\Omega t_r)$$

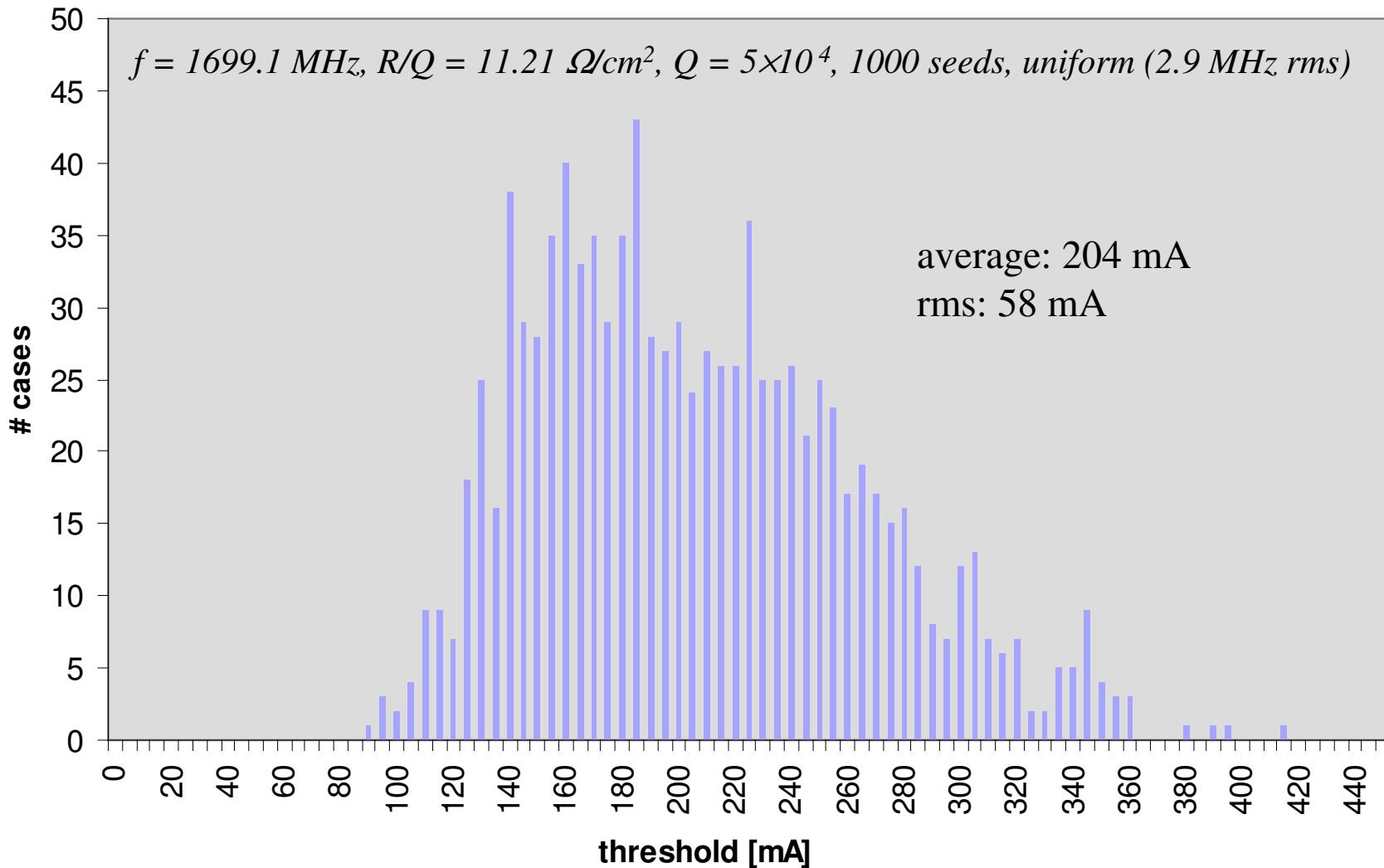


Periodic in $\Delta\omega = \Omega - \omega$ with $\omega_r = 2\pi/t_r$

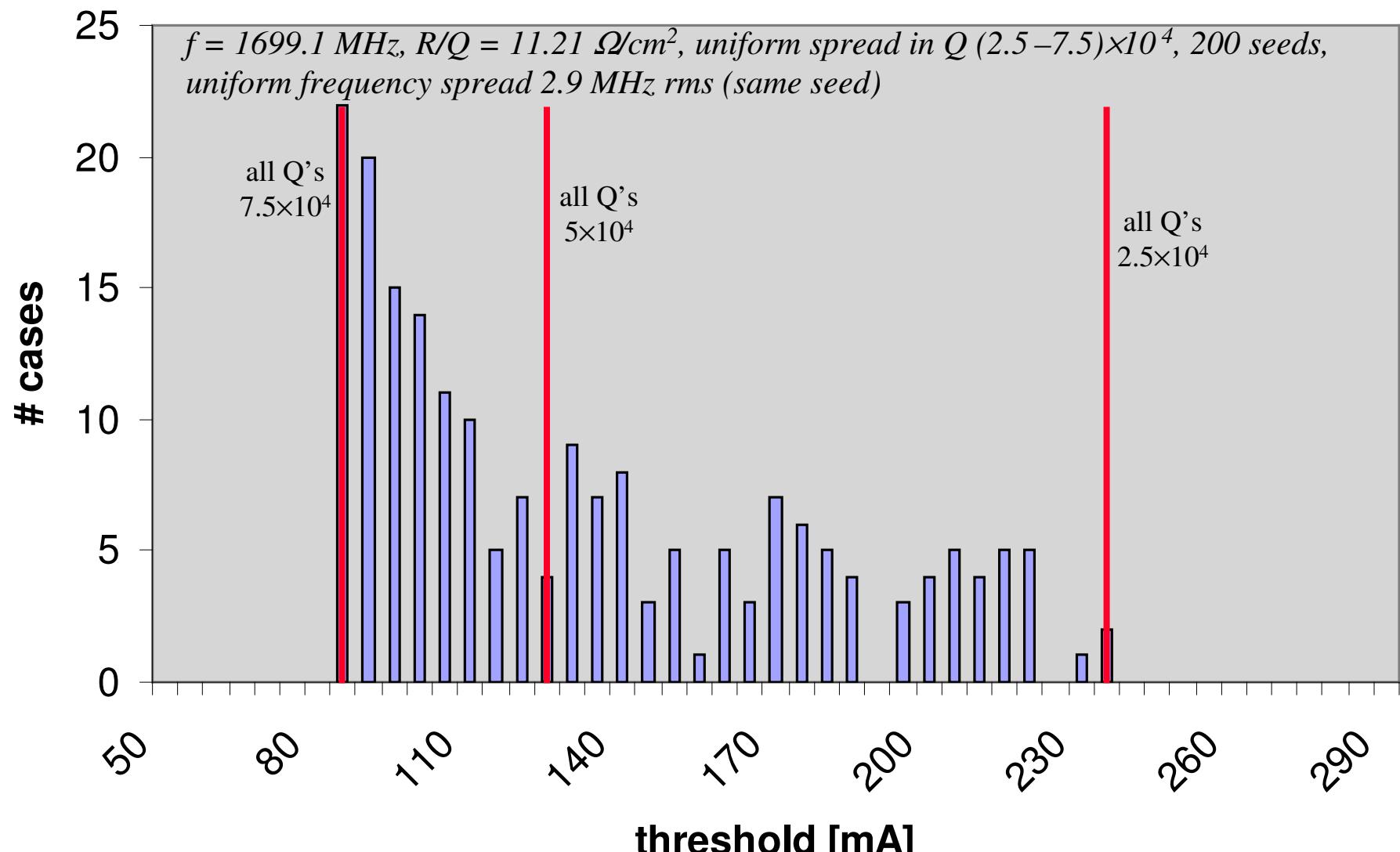
Threshold vs. frequency spread



Threshold for different seeds (10 MHz uniform)

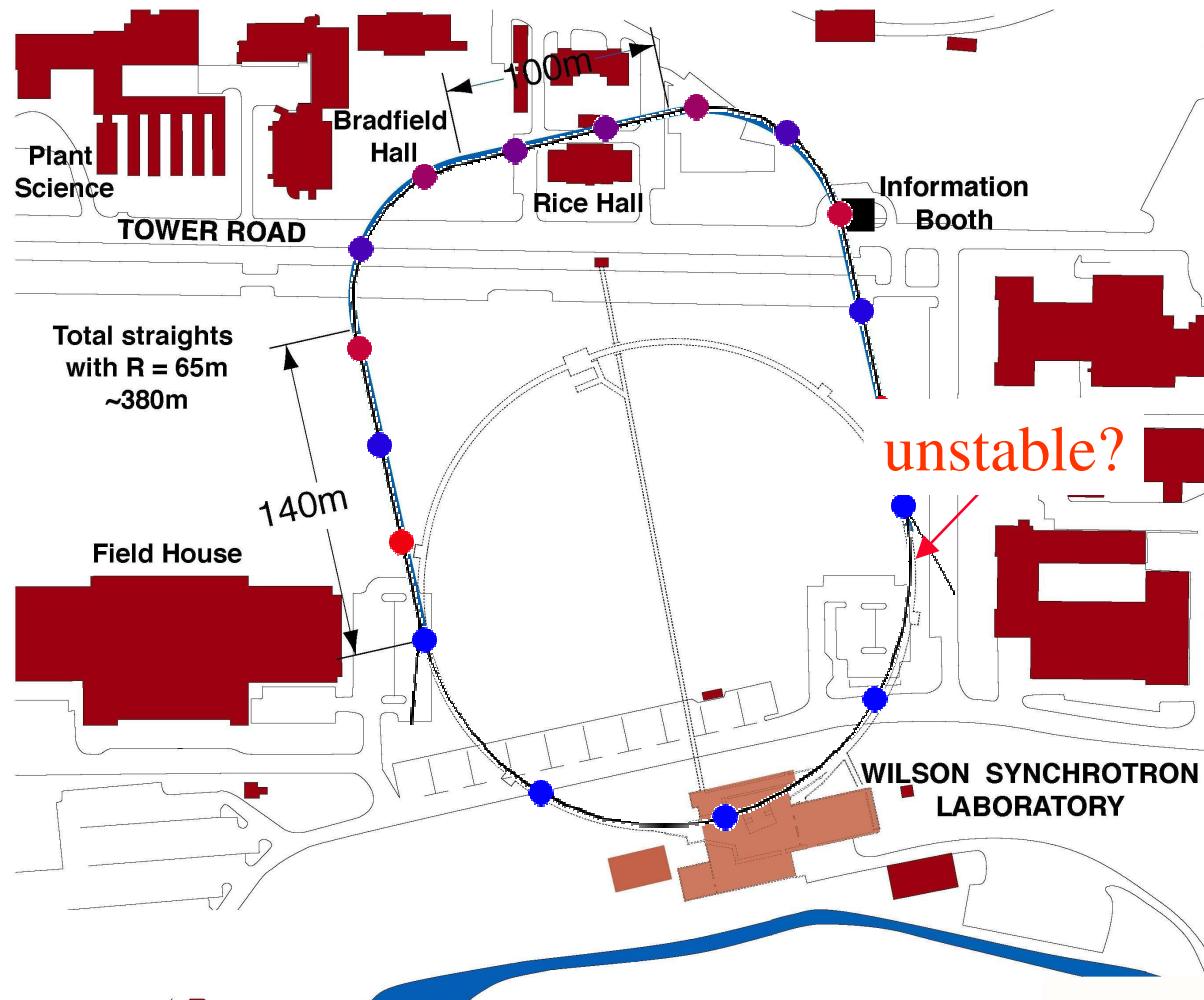


Spread in Q's



Direction we are heading to?

- Refrigerator
16.4 MW
- RF power
1.1 MW
- Dumped power
1 MW
- 7 undulators in
the current design



Two-recirculation option

- Need half the linac
 - ✓ linac cost
 - ✓ refrigeration
- BBU is currently limited to ~ 20 mA
- Current in linac is now 400 mA (difficult even for storage rings), e.g. beam induced effects go up by a factor of 4

BBU in two-pass ERL: gruesome look [from ERL memo 02-4]

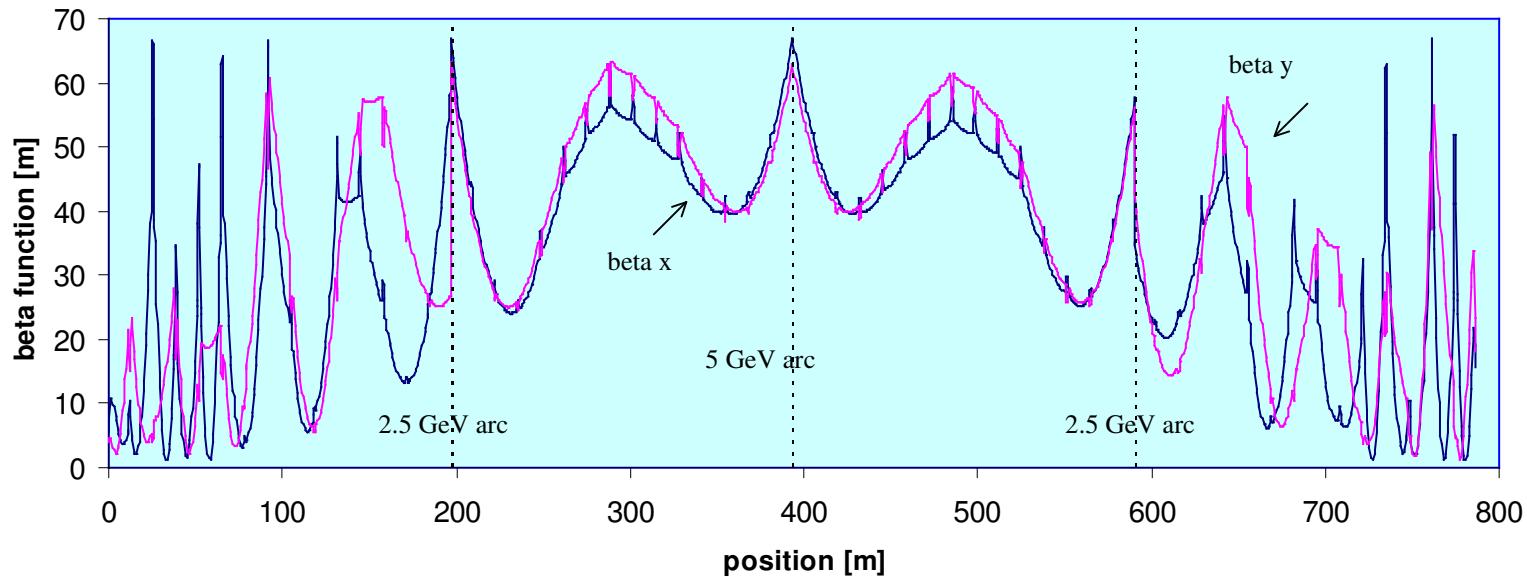


Table 1.
Results of TDBBU runs for 1-pass and 2-pass 5 GeV ERL.
HOM table (TESLA TDR 03/2001)

f (MHz)	R/Q (Ohm)	Q	(R/Q)*Q	1-pass 5 GeV ERL	2-pass 5 GeV ERL	Improved by
				BBU (mA)	BBU (mA)	a factor of
1699	88.40	5.00E+04	4.42E+06	160	20	8.00E+02
1873	56.39	7.00E+04	3.95E+06	190	25	1.30E+03
2575	51.50	5.00E+04	2.57E+06	115	15	9.00E+02
1725	118.64	2.00E+04	2.37E+06	135	15	5.00E+02
1864	42.84	5.00E+04	2.14E+06	> 200	40	2.00E+03
1880	11.08	1.00E+05	1.11E+06	> 200	90	8.00E+04
...	> 200	> 100	1.3

* BBU th >=100 mA

Trapped in the ~ 20 MW ERL?

Suggestion: drop the average current (~ 30 mA) and do two-recirculations

Pros: real savings (both construction & operation)
higher injection energy (~ 20 MeV)
lower space charge (improved brilliance)
more room for undulators

Cons: lower flux per meter of insertion device

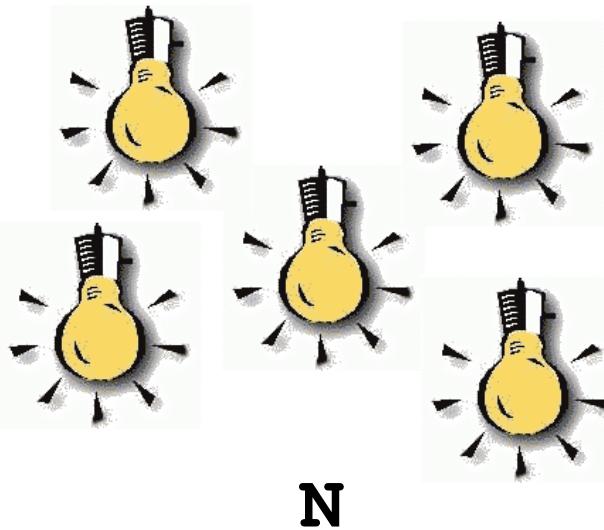
Maxim time

Never say “never”

Flux: how many bulbs does it takes?



N \times FLUX



BRILLIANCE

How much flux do people really care for?

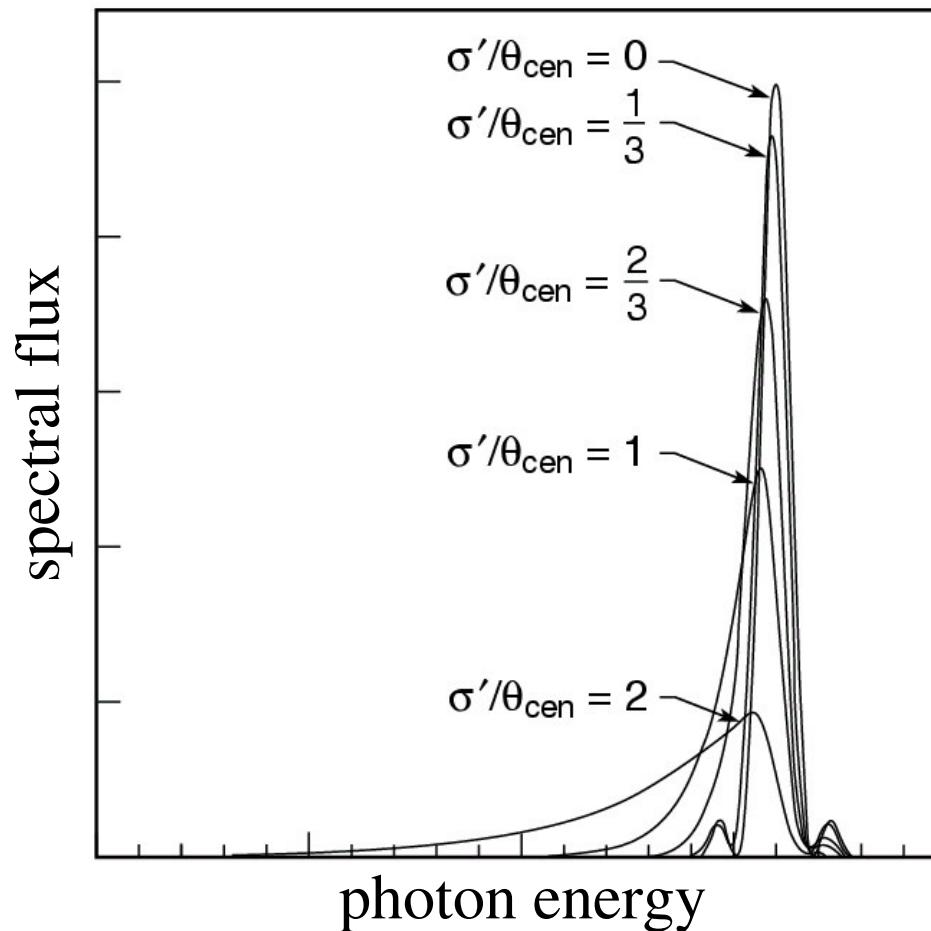
- At 100 mA one has $\sim \underline{1 \text{ kW}}$ of X-ray flux per meter of undulator
- At 30 mA, one would still have $\sim \underline{300 \text{ W}}$ per meter

$$\mathbf{FLUX} \times (\Delta\omega/\omega) \times (\varepsilon_{\text{ph}}/\varepsilon_{\perp})^2$$

spatial filtering with pin-hole for undulator does both

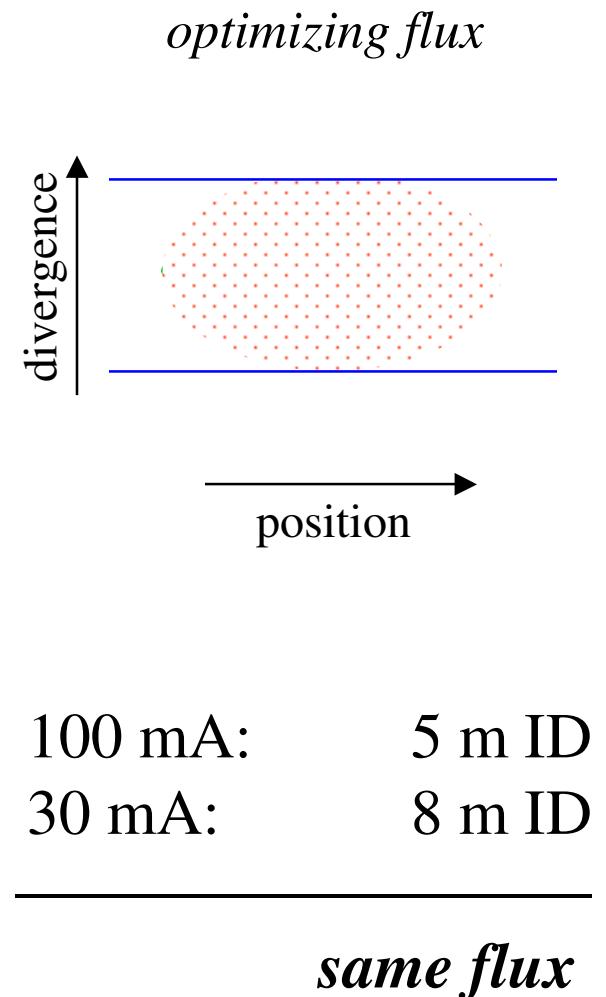
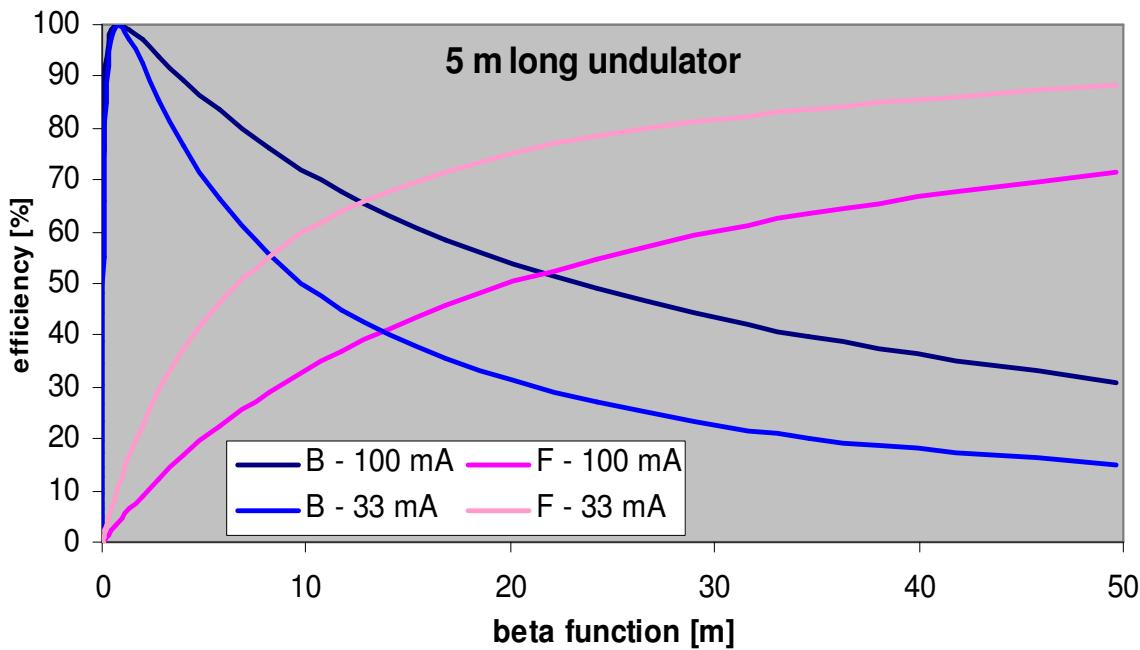
- Flux-driven experiments don't need small size and parallel beam – they need to be identified and dedicated IDs should be used for such applications

Brightness degradation due electron angular spread

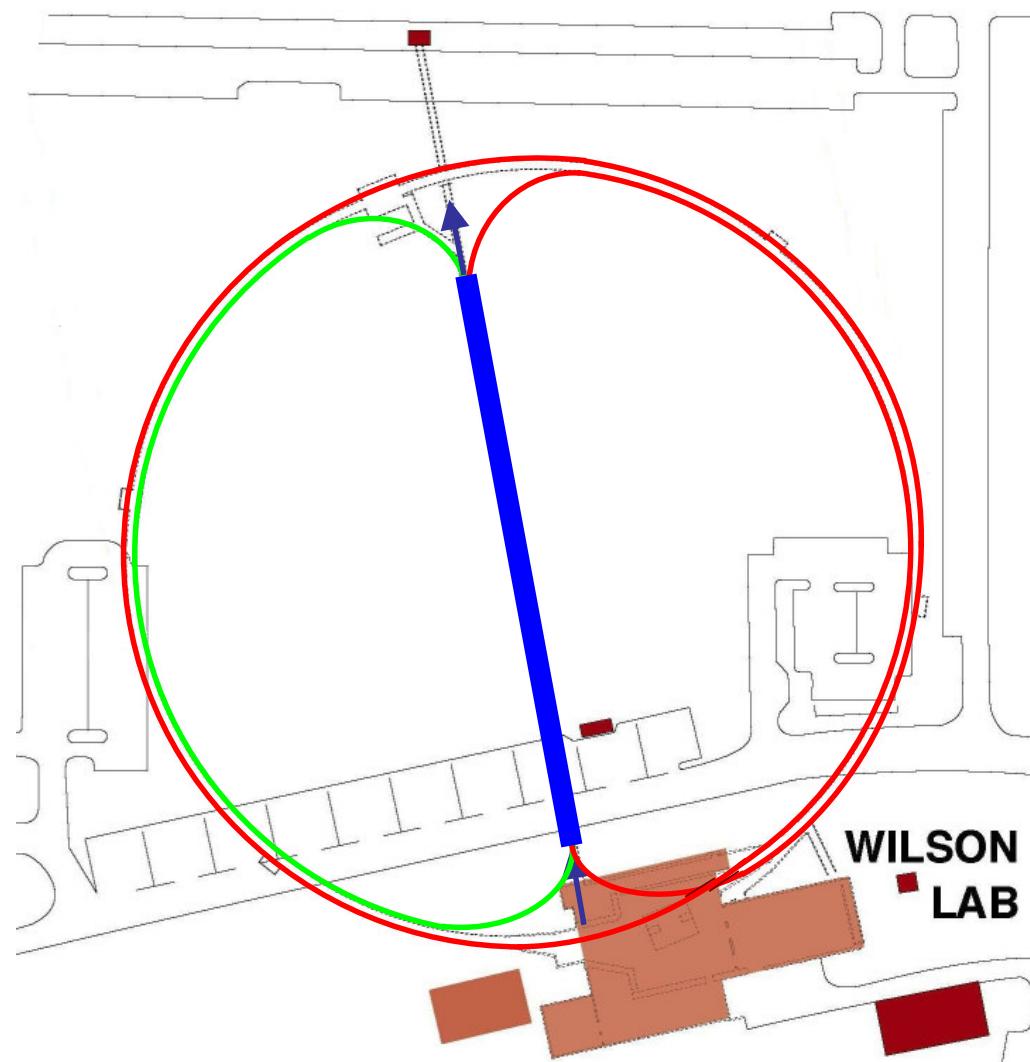


Beam matching for max brilliance and flux

$$\frac{dF_n}{d\Omega} = \frac{F_n}{2\pi\sqrt{\sigma_r^2 + \sigma_{x'}^2}\sqrt{\sigma_{r'}^2 + \sigma_{y'}^2}}$$



ERL in CESR tunnel II



Explaining tracking curve

