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Fundamental processes in III-V photocathodes; application for high-brightness photoinjectors



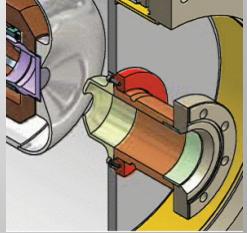
Contents

- Motivation
- NEA photoemission
- Some practical aspects
- Study cases: GaAs, GaAsP, GaN
- Summary



Why are we interested?

 Photoinjectors: a photocathode in high electric field (>> MV/m), either DC or RF

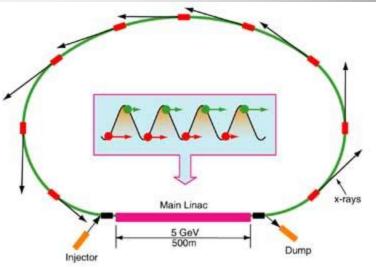


- Relativistic electrons can be further accelerated in a linac (linear accelerator) without degradation of beam brightness:
 - CW ultra-bright x-ray sources; high power FELs
 - Electron-ion colliders and ion coolers
 - Ultrafast electron diffraction, etc.



Energy recovery linac

- *Energy recovery linac*: a new class of accelerators in active development
- Essentially removes the average current limitation typical to linacs (i.e. $P_{beam} >> P_{wall plug}$)



• Average currents 10's to 100's of mA can be efficiently accelerated (and de-accelerated)



• QE and photon excitation wavelength

$$i(\text{mA}) = \frac{\lambda(\text{nm})}{124} \times P(\text{W}) \times \text{QE}(\%)$$

- E.g. 1W of 775 nm (Er-fiber $\lambda/2$) \Rightarrow 6.2 mA/% 520 nm (Yb-fiber $\lambda/2$) \Rightarrow 4.2 mA/% 266 nm (Nd-glass $\lambda/4$) \Rightarrow 2.1 mA/%

• Transversely cold (thermalized) electron distribution

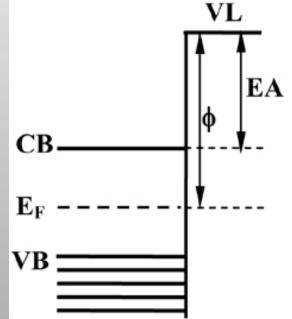
 Directly sets the solid angle of the emitted electrons; an upper limit on achievable beam brightness



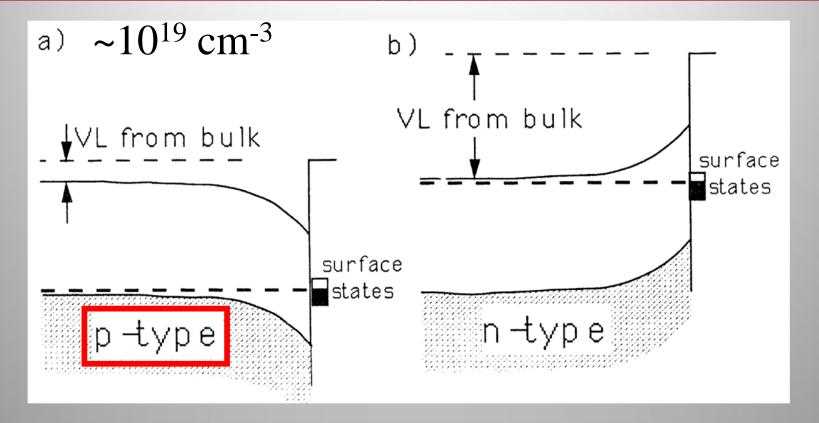
- Prompt response time
 - A picosecond response is essential to take advantage of the space charge control via laser pulse shaping
- Long lifetime and robustness
 - Extraction of many 100's to 1000's of C between activations are necessary to make the accelerator practical



- Defined as vacuum level E_{vac} relative to the conduction band minimum
- Negative affinity: the vacuum level lies below the CBM
 ⇒ very high QE possible
- NEA:
 - 1) band bending
 - 2) dipole layer



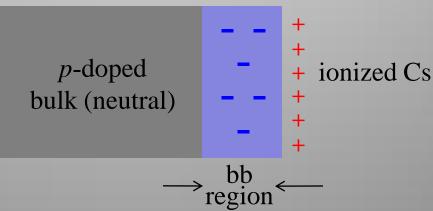
Strong p-doping



 Alperovich et al., Phys. Rev. B 50 (1994) 5480: clean p-doped GaAs has Fermi level unpinned and shows little band bending

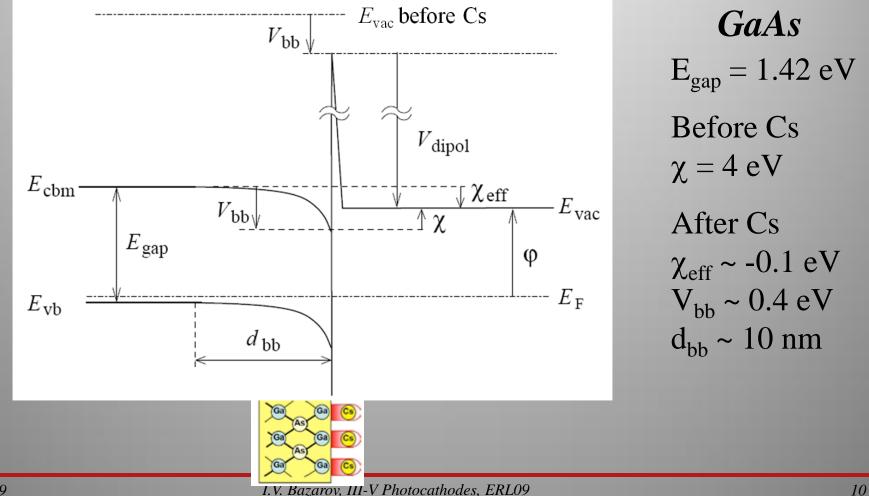


- Cs was found to play a larger role for NEA instead:
 1) band bending through donor surface states, and
 - 2) dipole surface layer from polarized Cs adatoms
 - Cs-induced donor-like surface states contribute their electrons to the bulk
 - Hole depleted region (negatively charged acceptors) lead to band bending region



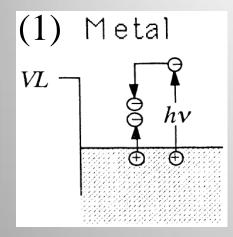
NEA: ~Cs monolayer (contd.)

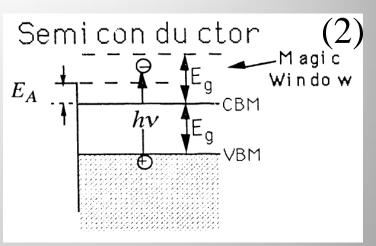
• Majority of Cs atoms become only polarized (not ionized), forming a dipole layer (e- Cs+)





Spicer's magic window

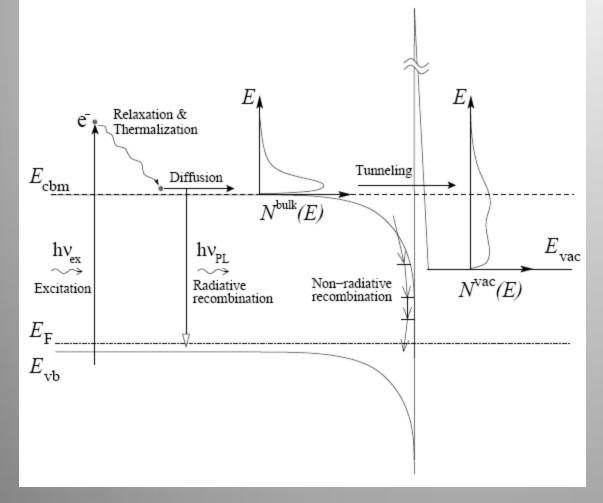




- (1) *electron-electron* scattering: typical of metals, large energy loss per collision
- (2) *electron-phonon* scattering: slowly depletes excessive energy of excited electron (LO phonons in GaAs ~ 35 meV) "*Magic window*": in semiconductors, one needs excess $KE > E_{gap}$ for e–/e– scattering. Thus, electrons excited with $E_{vac} < KE < E_{VBM} + 2E_{gap}$ have excellent chances of escape

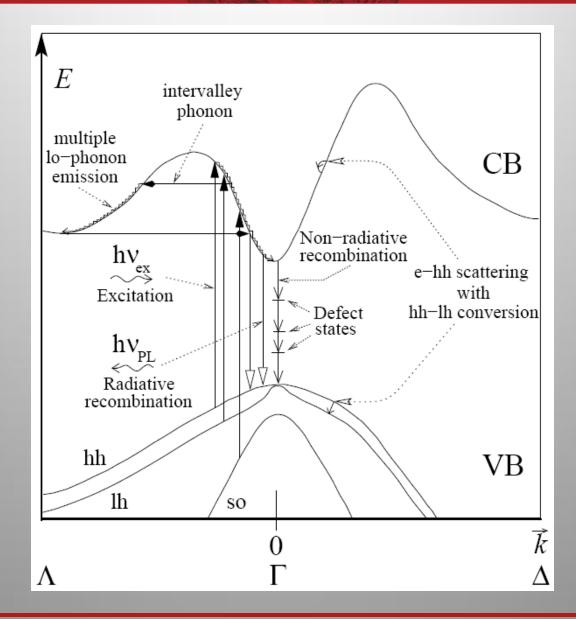


Electron transport processes



CBM thermalization time: 0.1-1ps Electron-hole recombination: ~ns **Emission time:** $\propto 1/(\alpha^2 D)$ strong wavelength dependence

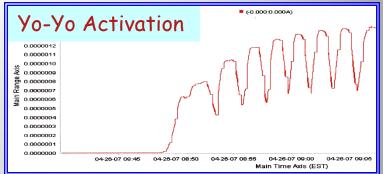
Energy vs. momentum

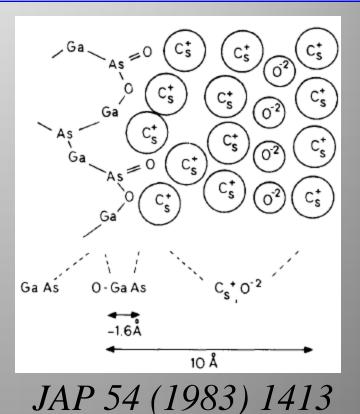




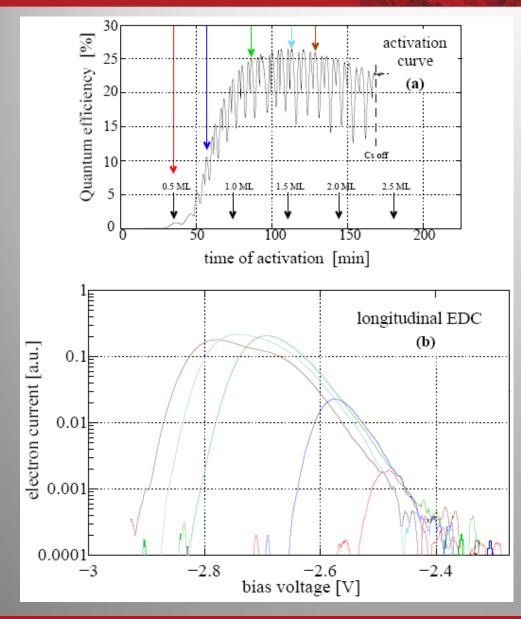
Role of fluorine/oxygen

- Routine "yo-yo" activation employs O₂ or NF₃
- Further reduction of affinity consistent with a double dipole model
- Stabilizes Cs on the surface; no lifetime or otherwise apparent advantage for either gas
- Bonded unstable nitrogen is found on Cs-NF₃ activated surfaces (APL 92, 241107)





GaAs: Optimal Cs coverage



laser wavelength: 670 nm

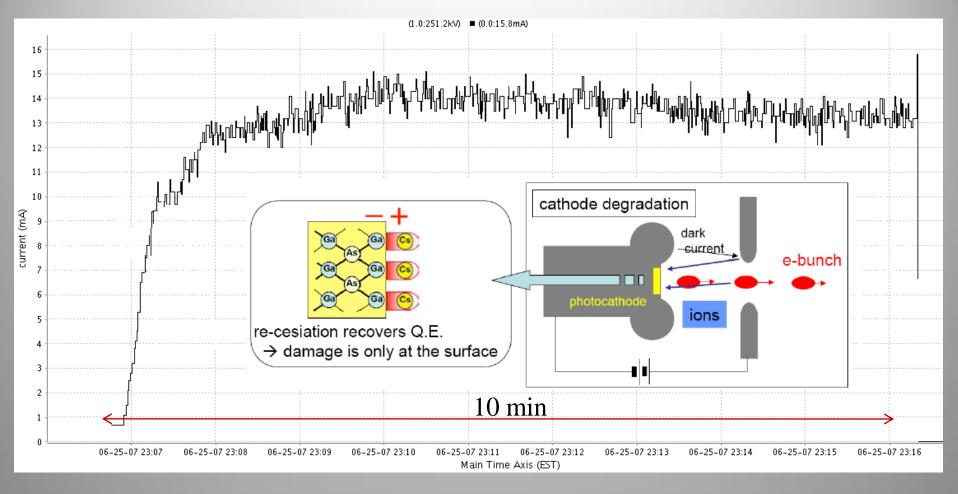
Ugo Weigel, PhD thesis



- H₂O, CO₂ and O₂ can lead to chemical poisoning of the activated layer
- Low current (~ 1µA) 1/e lifetimes ~ 100 hours typical in our prep chambers
- 3-5 times better in the DC gun (low 10⁻¹² Torr vacuum)
- High average current (mA's) lifetime limited by ion backbombardment



14 mA



~5 hour lifetime (limited by gas backstreaming from the beam dump), i.e. 20 hours 1/e for 5 mA



Study cases

- Our group has been evaluating III-V photocathodes
 - Transverse energy of electrons (thermal emittance)
 - Measure the photoemission response time
- Materials studied so far
 - GaAs @ 450-850nm: JAP 103, 054901; PRST-AB 11, 040702
 - GaAsP @ 450-640nm: Ibid
 - GaN @ 260nm: JAP 105, 083715



Diffusion model

$$\frac{\partial c(h,t)}{\partial t} = D \frac{\partial^2 c(h,t)}{\partial h^2}$$
subject to:

$$c(h,t=0) = c_0 e^{-\alpha h}$$

$$c(h=0,t) = 0$$

$$I(t) \propto \frac{\partial}{\partial t} \int_{0}^{\infty} c(h,t) dh.$$

$$I(\kappa) \propto \frac{1}{\sqrt{\pi\kappa}} - \exp(\kappa) \operatorname{erfc}(\sqrt{\kappa})$$

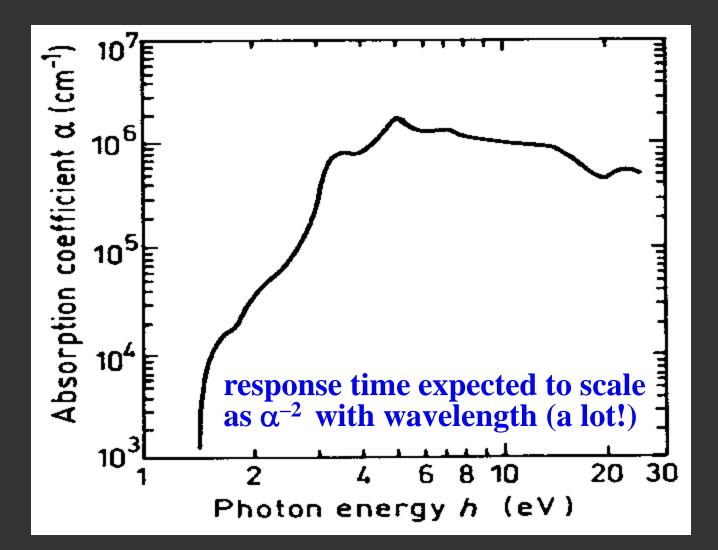
$$\kappa \equiv t/\tau, \text{ where } \tau \equiv \alpha^{-2}D^{-1}$$

$$\int_{0}^{1} \frac{1}{\sqrt{\pi\kappa}} \int_{0}^{\infty} \frac{1}{\sqrt{\pi\kappa}}$$

NEA photocathodes meas.







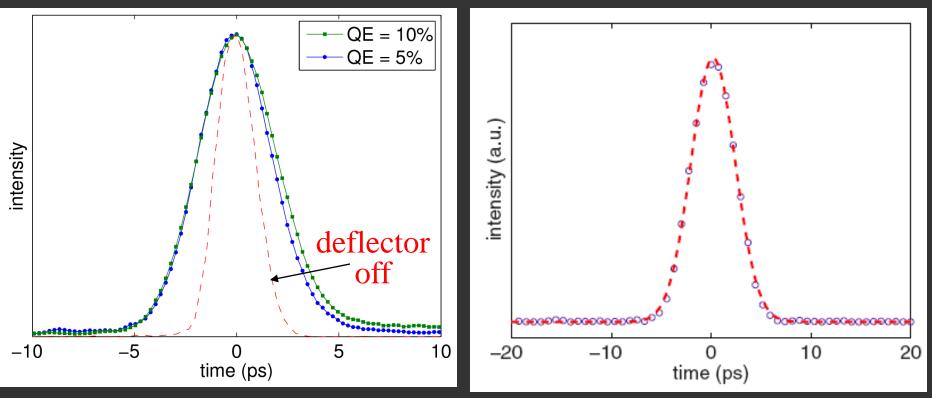
NEA photocathodes meas.



Prompt emitters

GaAs @ 520 nm





Measurements done by transverse deflecting RF cavity Limited by 1.8 ps rms resolution dominated by laser to RF synchronization

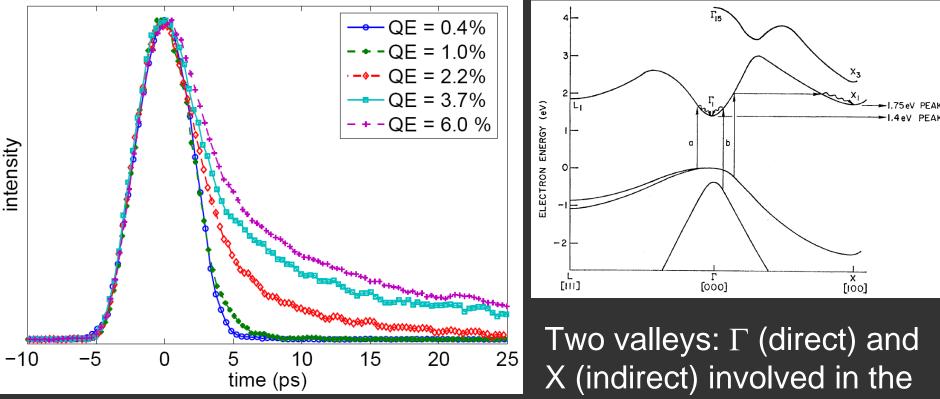
NEA photocathodes meas.



Diffusion tail

GaAsP @ 520 nm

P concentration 45%

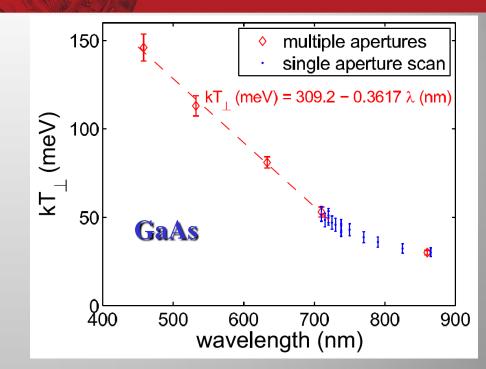


Strong QE dependency

process

Transverse energy distributions

- Cornell University
- No surprises for bulk GaAs: cold electrons with a near band-gap excitation
- Surprisingly large transverse energy spread for GaN and GaAsP:



- GaAsP: $kT_{\perp} = 130-240$ meV for photons 0-780 meV photons above the band-gap
- GaN: $kT_{\perp} = 0.9$ eV for photons with 1.4 eV above the band-gap



- Transverse energy of photoelectrons remain poorly understood for III-V semiconductors (other than GaAs)
- More carefully controlled experimental data on transverse energy distributions/time response needed
- Predictive codes and models need to be developed and benchmarked with experiments
- This will allow photocathode engineering with the desired characteristics such as cold electrons with a ps response



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