

Sept 4.

Quantum motion in phase space

Motivation

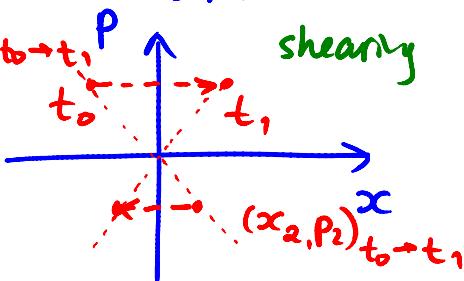
- ① class. mechanics: part. state (x, p)
+ Hamiltonian \Rightarrow full knowledge
(including past & future)

Time evolution: Hamilton's eqns (1D)

$$\begin{cases} \dot{p} = -\frac{\partial \mathcal{H}}{\partial x} \\ \dot{x} = \frac{\partial \mathcal{H}}{\partial p} \end{cases}$$

Ex1: simple motion (drift)

$$\mathcal{H}(x, p) = KE + PE = \frac{p^2}{2m}$$

$$\begin{cases} \dot{p} = 0 & (x_1, p_1)_{t_0 \rightarrow t_1} \\ \dot{x} = \frac{p}{m} & \end{cases}$$


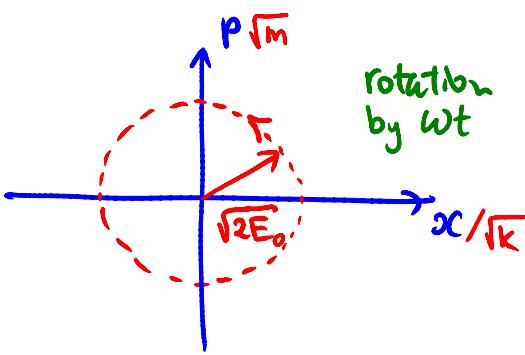
Ex2: simple motion (linear restor. force)

$$\mathcal{H} = \frac{p^2}{2m} + \frac{kx^2}{2} = E_0$$

$$\begin{cases} \dot{p} = -kx \\ \dot{x} = \frac{p}{m} \end{cases}$$

$$\ddot{x} + \frac{k''}{m}x = 0$$

$$\ddot{x} + \frac{k}{m}x = 0$$



② QM state: $|\psi\rangle$, \hat{X} = full knowledge

$$\begin{array}{ccc} \text{position} & & \text{momentum} \\ \psi(x) \equiv \langle x | \psi \rangle & \xleftarrow{\text{FT}} & \tilde{\psi}(p) \equiv \langle p | \psi \rangle \end{array}$$

$$\begin{array}{ccc} |\psi(x)|^2 dx & & |\tilde{\psi}(p)|^2 dp \\ p(x, x+dx) & & p(p, p+dp) \end{array}$$

probability density $W(x, p)$ - ?

$[\hat{x}, \hat{p}] = i\hbar \neq 0$: incompatible operators
 \Rightarrow no joint eigenbasis

But can still define $W(x, p)$ - [quasi] probability

③ Wigner distribution:

$$W(x, p) = \underbrace{\int_{-\infty}^{\infty} \langle \psi | x + \frac{x'}{2} \rangle \langle x + \frac{x'}{2} | p \rangle \langle p | x - \frac{x'}{2} \rangle \langle x - \frac{x'}{2} | \psi \rangle dx'}_A$$

Dirac notation: quantum trajectory from $x - \frac{x'}{2}$ to $x + \frac{x'}{2}$
of integrand with momentum p

$\int_{-\infty}^{\infty} \dots dx'$: superposition of all possible trajectories

Expanding A:

$$\langle x | p \rangle = \frac{1}{\sqrt{2\pi\hbar}} e^{\frac{ipx}{\hbar}} \leftarrow \text{HW1}$$

$$A: \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} e^{\frac{ip}{\hbar}(x+\frac{x'}{2})} e^{-\frac{ip}{\hbar}(x-\frac{x'}{2})} = \frac{1}{2\pi\hbar} e^{\frac{ipx'}{\hbar}}$$

$$W(x, p) = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} \psi^*(x + \frac{x'}{2}) \psi(x - \frac{x'}{2}) e^{\frac{ipx'}{\hbar}} dx'$$

④ Can write the same thing in momentum space

$$W(x, p) = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} \langle \psi | p + \frac{p'}{2} \rangle \langle p + \frac{p'}{2} | x \rangle \langle x | p - \frac{p'}{2} \rangle \langle p - \frac{p'}{2} | \psi \rangle dp'$$

$$W(x, p) = \frac{1}{2\pi\hbar} \int_{-\infty}^{\infty} \tilde{\Psi}(p + \frac{p'}{2}) \tilde{\Psi}(p - \frac{p'}{2}) e^{-\frac{ip'x}{\hbar}} dp'$$

where $\tilde{\Psi}(p) = \frac{1}{\sqrt{2\pi\hbar}} \int \psi(x) e^{-\frac{ipx}{\hbar}} dx$

$\psi(x) = \frac{1}{\sqrt{2\pi\hbar}} \int \tilde{\Psi}(p) e^{\frac{ipx}{\hbar}} dp$

FT

⑤ Properties of Wigner distribution function WDF

a) $W^*(x, p) = W(x, p)$ - or $W(x, p) \in \mathbb{R}$

b) $\iint W(x, p) dx dp = 1$ - normalization

c) $\int W(x, p) dx = |\tilde{\Psi}(p)|^2$ } marginals
 $\int W(x, p) dp = |\psi(x)|^2$

d) $|W(x, p)| \leq \frac{1}{\pi\hbar}$ - boundedness

⑥ Expectation values :

operator $\hat{A} \rightarrow A(x, p)$
 Wigner - Weyl transformation

$$\langle A \rangle = \iint A(x, p) W(x, p) dx dp$$

Some nice examples :

$$\hat{x}^n \rightarrow x^n$$

$$\hat{p}^n \rightarrow p^n$$

$$\frac{1}{2} (\hat{x}\hat{p} + \hat{p}\hat{x}) \rightarrow xp$$

⑦ Time evolution of WDF $W(x, p; t)$

$$\frac{\partial W}{\partial t} + \frac{p}{m} \frac{\partial W}{\partial x} + \frac{i}{\hbar} \left[V\left(x + \frac{i\hbar}{2} \frac{\partial}{\partial p}\right) - V\left(x - \frac{i\hbar}{2} \frac{\partial}{\partial p}\right) \right] W = 0$$

HW proof

$$\text{let } V(x) = V_0 - F_0 x + \frac{1}{2} k x^2$$

$$\text{or } F(x) = F_0 - kx ,$$

$$\Rightarrow \frac{\partial W}{\partial t} + \frac{p}{m} \frac{\partial W}{\partial x} + F \frac{\partial W}{\partial p} = 0$$

same as Liouville's eqn!

show that:

$$\begin{aligned} & \frac{i}{\hbar} \left[V_0 - F_0 \left(x + \frac{i\hbar}{2} \frac{\partial}{\partial p} \right) + \frac{1}{2} k \left(x + \frac{i\hbar}{2} \frac{\partial}{\partial p} \right)^2 \right. \\ & \quad \left. - V_0 + F_0 \left(x - \frac{i\hbar}{2} \frac{\partial}{\partial p} \right) - \frac{1}{2} k \left(x - \frac{i\hbar}{2} \frac{\partial}{\partial p} \right)^2 \right] \end{aligned}$$

$$= \frac{i}{\hbar} \left[-F_0 i \hbar \frac{\partial}{\partial p} + kx i \hbar \frac{\partial}{\partial p} \right]$$

$$= \left(F_0 - kx \right) \frac{\partial}{\partial p} = F(x) \frac{\partial}{\partial p} \quad \therefore$$

⑧ Implications

- * WDF evolves classically with time for linear forces
- * WDF can be measured using tomography method
(same as measuring $\Psi(x)$, not $|\Psi(x)|^2$!)
- * "Squeezed" states, "coherent" states