

• Photon Polarization States II

- ① Linear polarization analyzer and Projection amplitudes:
- ② Linear polarization analyzer loop / change of basis states:

## Recap

### III<sub>6</sub> Dirac Notation:

Dirac: chop bracket notation for inner product in two pieces:

bra:  $\langle \alpha |$   $\leftrightarrow$  instruction to integrate / row vector

ket:  $| \beta \rangle$   $\leftrightarrow$  state vector

state:

$$| \alpha \rangle = \sum_n | e_n \rangle \langle e_n | \alpha \rangle$$

call  $| e_n \rangle \langle e_n |$  projection operator

• Can expand state  $| \alpha \rangle$  in terms of different basis states:

$$| \alpha \rangle = \sum_n | e_n \rangle \langle e_n | \alpha \rangle = \sum_n | e_n' \rangle \langle e_n' | \alpha \rangle$$

• Unity operator:

$$\sum_n | e_n \rangle \langle e_n | = 1 \quad (\text{if it acts on } | \alpha \rangle, \text{ one gets } | \alpha \rangle)$$

## Recap

### III, Another example: Photon Polarization States

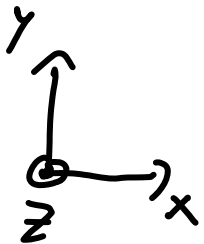
- Light

classically	QM picture
electromagnetic wave	photons

In the following: all waves / photons have same frequency and direction of motion (+ z-direction) ←

light with arbitrary polarization:

$$\vec{E}(z,t) = \text{Re}\{(a e^{i\phi_a} \hat{x} + b e^{i\phi_b} \hat{y}) e^{i(kz - \omega t)}\}$$
$$= a \cos(kz - \omega t + \phi_a) \hat{x} + b \cos(kz - \omega t + \phi_b) \hat{y}$$



Photon state:

$$|\psi\rangle$$

Recap

Classically

photon

1) linear polarized in x-direction:

$\Rightarrow b = 0$



$|x\rangle$

2) linear polarized in y-direction:

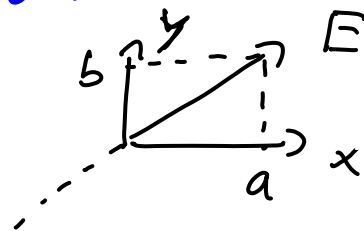
$\Rightarrow a = 0$



$|y\rangle$

3) linear polarized in  $a\hat{x} + b\hat{y}$  direction:

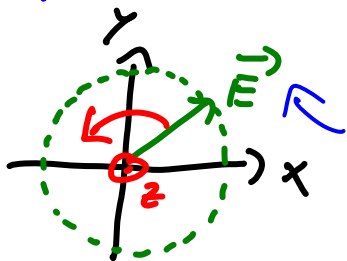
$\Rightarrow a, b \neq 0 \quad f_a = f_b$



$\frac{1}{\sqrt{a^2+b^2}} \{ a|x\rangle + b|y\rangle \}$

↑ normalize

4) left circular polarized:



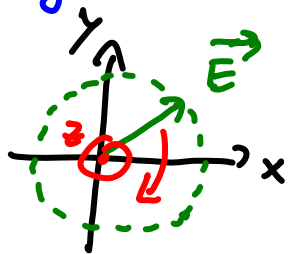
$|\vec{E}| = \text{const}$

$a = b$

$f_a = f_b - \frac{\pi}{2}$   
90°

$|L\rangle = \frac{1}{\sqrt{2}} \{ i|x\rangle + |y\rangle \}$

5) right circular polarized:



$a = b$

$f_a = f_b + \frac{\pi}{2}$

$|R\rangle = \frac{1}{\sqrt{2}} \{ (-i)|x\rangle + |y\rangle \}$

## Recap

① Linear polarization analyzer and Projection amplitudes:

$$E_x = a \cos(kz - \omega t + \phi_a) = \hat{x} \cdot \vec{E}_i$$

$$I_x \propto \frac{1}{2} a^2$$

$$E_y = b \cos(kz - \omega t + \phi_b) = \hat{y} \cdot \vec{E}_i$$

$$I_y \propto \frac{1}{2} b^2$$

detect.

$|x\rangle, \langle x|\psi\rangle$

detect.

$|y\rangle, \langle y|\psi\rangle$

x-y-analyzer

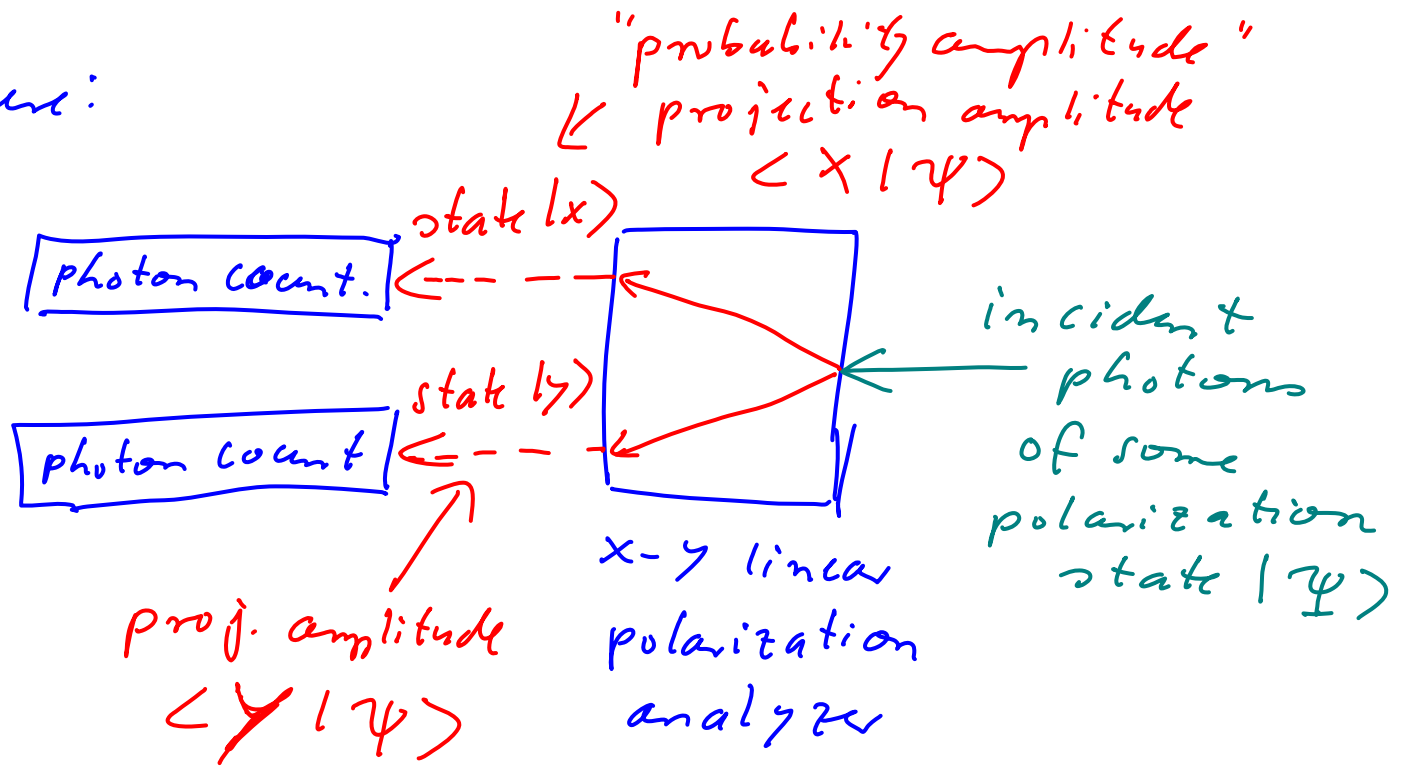
Projection!

$$\vec{E}_i = a \cos(kz - \omega t + \phi_a) \hat{x} + b \cos(kz - \omega t + \phi_b) \hat{y}$$

$$I_{in} \propto \overline{|\vec{E}|^2} = \frac{1}{2} a^2 + \frac{1}{2} b^2$$

separates light into linear  $\hat{x}$ -pol. and linear  $\hat{y}$ -pol. components

# Photon Picture:



• Input intensity / # of photons = total output intensity / total # of photons

⇒  $|x\rangle$  and  $|y\rangle$  linear polarization states give a complete set of basis states (only 2 linear indep. states)

$$\begin{aligned} \Rightarrow |\Psi\rangle &= c_x |x\rangle + c_y |y\rangle = |x\rangle \langle x|\Psi\rangle + |y\rangle \langle y|\Psi\rangle \\ &= \sum_n |\text{state}_n\rangle \langle \text{state}_n|\Psi\rangle \end{aligned}$$

•  $|x\rangle$  and  $|y\rangle$  polarization states are orthonormal to each other:

ortho-normal!

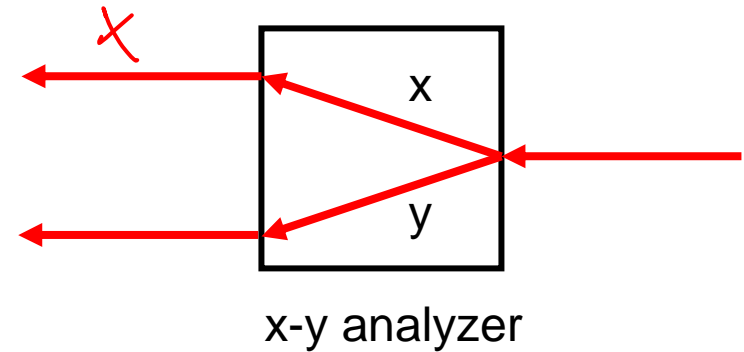
$$\begin{cases} \langle x|x\rangle = 1 \\ \langle y|x\rangle = 0 \\ \langle x|y\rangle = 0 \\ \langle y|y\rangle = 1 \end{cases}$$

↑ onto state    ↑ from state

X-Y analyzer

Suppose a single photon is sent through an x-y analyzer, and ends up in the x-channel.

Does that mean that the original photon was polarized in the x-direction?



A. yes

B. no

*example:*

$$|\psi\rangle_{in} = \frac{1}{\sqrt{2}} |x\rangle + \frac{1}{\sqrt{2}} |y\rangle$$



• not all photons behave the same way!

(some end up in the  $\hat{x}$ -detector, some in the  $\hat{y}$ -detector, even though they all start with the same polarization prior to the measurement!)

=> statistical behaviour in QM, probabilities only!

Example:  $|\Psi\rangle_{in} = \frac{1}{\sqrt{2}}|x\rangle + \frac{1}{\sqrt{2}}|y\rangle \Rightarrow$   $\left\{ \begin{array}{l} 50\% \text{ chance result is } |x\rangle \rightarrow \text{photon is in } |x\rangle \text{ state after measurement.} \\ 50\% \text{ chance result is } |y\rangle \rightarrow \text{photon is in } |y\rangle \text{ state after measurement.} \end{array} \right.$

=> Measurement can change the quantum state!  
(wave function "collapses")

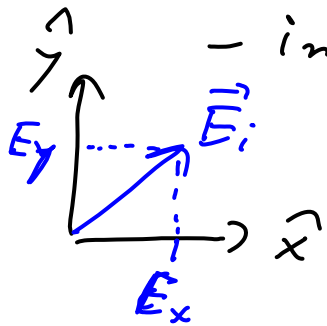
=> Measurements give only discrete results

( $x$ - $y$  analyzer: only  $\hat{x}$  or  $\hat{y}$  polarized, nothing in between)

=> Measurements = projection of initial state onto some basis state

## Projection amplitudes:

→ classically:



- initial:  $\vec{E}_i = a \cos(kz - \omega t + \phi_A) \hat{x} + b \cos(kz - \omega t + \phi_B) \hat{y}$

- after  $\hat{x}$ - $\hat{y}$  analyzer:

$\hat{x}$ -path:  $E_x = \hat{x} \cdot \vec{E}_i = a \cos(kz - \omega t + \phi_A)$

measurement = projection onto x-axis

Intensity along x  $\propto |E_x|^2 = |\hat{x} \cdot \vec{E}_i|^2$

$\hat{y}$ -path:  $E_y = \hat{y} \cdot \vec{E}_i = b \cos(kz - \omega t + \phi_B)$

Intensity  $\propto |E_y|^2 = |\hat{y} \cdot \vec{E}_i|^2$

→ Quantum picture:

- initial:  $|\Psi\rangle$

- after  $\hat{x}$ - $\hat{y}$ -analyzer:

$\hat{x}$ -path:  $|\text{photon}\rangle_{\text{after}} = |x\rangle$

define  $\langle x | \Psi \rangle =$  projection amplitude of incoming state  $|\Psi\rangle$  onto final pol. state  $|x\rangle$

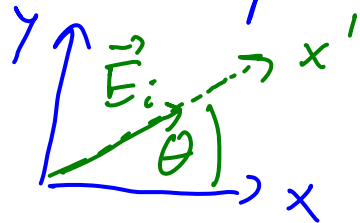
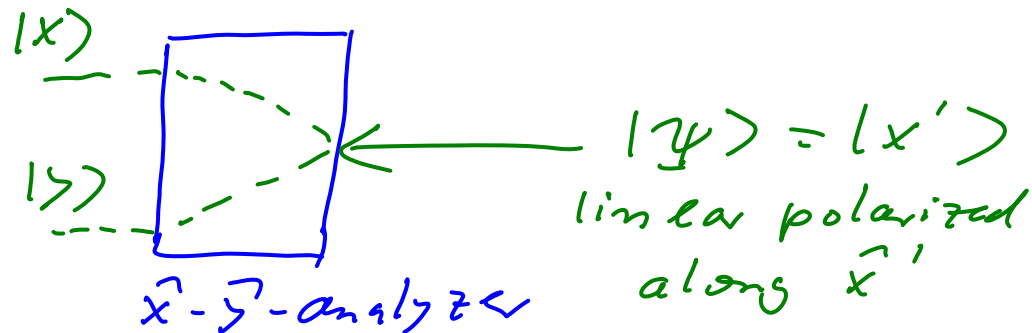
$\Rightarrow |C_x|^2 = |\langle x | \Psi \rangle|^2 =$  probability of finding photon in  $|x\rangle$  pol. state

similar for  $\hat{y}$ -path

$C_y = \langle y | \Psi \rangle$

$|C_y|^2 = |\langle y | \Psi \rangle|^2$

Example:  $|\psi\rangle = |x'\rangle$ : linear pol. along  $\hat{x}'$

→ Classically:  
- initial:

$$\vec{E}_i = E_0 \cos \theta \cos(kz - \omega t) \hat{x}' + E_0 \sin \theta \cos(kz - \omega t) \hat{y}'$$

$$= E_0 \cos(kz - \omega t) \hat{x}'$$

$$\text{Intensity} = I_i \propto \frac{1}{2} |E_0|^2$$

- after  $\hat{x}$ - $\hat{y}$  analyzer:

$$\hat{x}\text{-path: } E_x = \hat{x} \cdot \vec{E}_i = E_0 \cos \theta \cos(kz - \omega t)$$

$$I_x \propto |E_x|^2 = \frac{1}{2} E_0^2 \cos^2 \theta \rightarrow I_x = I_i \cos^2 \theta$$

$$\hat{y}\text{-path: } E_y = \hat{y} \cdot \vec{E}_i = E_0 \sin \theta \cos(kz - \omega t)$$

$$I_y \propto |E_y|^2 = \frac{1}{2} E_0^2 \sin^2 \theta \rightarrow I_y = I_i \sin^2 \theta$$

note:  $I_x + I_y = I_i$

→ Quantum picture:

$$+\text{sign} \langle x|x \rangle = 1$$

$\hat{x}$ -path: projection amplitude:  $\langle x|x' \rangle \stackrel{\downarrow}{=} \cos \theta$

$\hat{y}$ -path: " " "  $\langle y|x' \rangle = \sin \theta$

⇒ probability of measuring photon in x-pol. state  $|x\rangle$

$$|C_x|^2 = |\langle x|x' \rangle|^2 = \cos^2 \theta \propto \underbrace{I_x}_{\text{classically}}$$

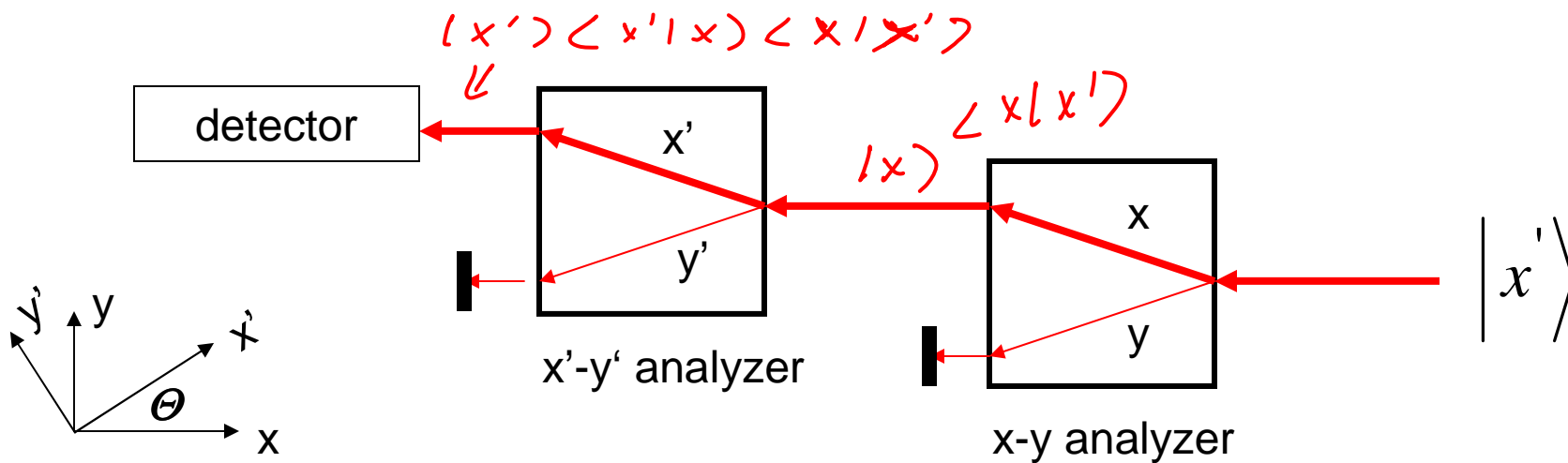
⇒ probability of measuring photon in y-pol. state  $|y\rangle$

$$|C_y|^2 = |\langle y|x' \rangle|^2 = \sin^2 \theta \propto I_y$$

Note:  $|C_x|^2 + |C_y|^2 = 1 = \sum_n |C_n|^2 = 1$

⇒ general: if  $|a\rangle, |b\rangle$  are linear polarization states, then

$$\langle a|b \rangle = \cos(\text{angle between polarizations})$$



What is the probability that one of the initial photons in linear polarization state  $|x'\rangle$  reaches the detector?

- A. 0
- B.  $\cos(\Theta)^2$
- C.  $\cos(\Theta)^4$**
- D.  $\cos(\Theta)^2 \sin(\Theta)^2$
- E. 1

total prob. of getting to detector

$$= |\langle x | x' \rangle|^2 |\langle x' | x \rangle|^2$$

$$= (\cos \Theta)^2 \cdot (\cos \Theta)^2$$

$$= \cos^4 \Theta$$