

# David Marsh <sup>by</sup> SUGRA ~~or~~ PHENOS

Motivation : susy, strings

- cosmo :  $\Lambda$ , inflation, reheating, baryogen.
- PHENO : ~~susy~~ in fact, all rel to ~~susy~~
- GAUGE/GRAVITY : ~~today~~

~~susy~~ in SUGRA

$$\mathcal{L} = \tilde{g} + \sqrt{2} \Delta \psi + F_{\mu\nu} \theta^2$$

relevant ops

$$\mathcal{L}_{\text{soft}} = \frac{1}{2} M_{\text{soft}} |\tilde{\phi}|^2 + \frac{1}{2} B \phi^2 + \text{c.c.} + \frac{1}{3!} A_{ijk} \tilde{\phi}^i \tilde{\phi}^j \tilde{\phi}^k + \dots + \text{c.c.} + M_{\lambda} \lambda \lambda + \text{c.c.}$$

SOFT MASS

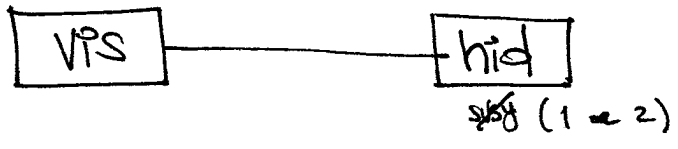
"HOMOGENEOUS" TERMS

GAUGINO MASSES

low e phenomenology determined by these  
 GEN. of these from UV theory is  
 story of ~~susy~~

## Susy Scenarios : How to gen.

- |                      |   |                 |                           |
|----------------------|---|-----------------|---------------------------|
| 1. NON-REN. OPS      | → | GRAV MEDIATION  | } <sup>mix:</sup> anomaly |
| 2. LOOP INTERACTIONS | → | GAUGE MEDIATION |                           |



# 1. GRAVITY MEDIATION

CONSIDER FUNDAMENTAL THEY w/

$$K = \frac{1}{M_{Pl}^2} X^\dagger X Q^\dagger Q$$

$$\uparrow$$
$$X = x + \theta^2 F_x$$

$$\mathcal{L} = \int d^4\theta \frac{1}{M_{Pl}^2} X^\dagger X Q^\dagger Q$$
$$= \frac{1}{M_{Pl}^2} F_x^2 \tilde{q}^\dagger \tilde{q}$$

↑ similar for B ~~is~~ term

A  $\sim$   $M_x$  from similar game in superpot.

eg if  $F_x \sim (10^{10} \text{ GeV})^2$

then acceptable soft masses

How to generate this in SUGRA?

$$\mathcal{L} = (C_{ij}) \int d^4\theta \frac{1}{M_{Pl}^2} X^\dagger X \underbrace{Q_i^\dagger Q_j}_{\uparrow \text{ flavors}}$$

↑  
mass matrix

FCNC  $\Rightarrow$   $C_{ij}$  wants to be diagonal

BUT "GRAVITY DOESN'T RESPECT GLOBAL SYM" so

No flavor games to play.

# Conformal compensator formalism

motiv: promote on-shell action to off-shell

idea: start w/ SUPERCONFORMAL SUGRA

≠ break to non-conformal SUGRA

used by Randall & Sundrum - out of this world susy

$$\frac{1}{\sqrt{-g}} \mathcal{L} = \int d^4\theta \Phi^\dagger \Phi f(\mathcal{Q}^\dagger, \mathcal{Q}) + \int d^2\theta (\Phi^3 W(\mathcal{Q}) + \tau W_\alpha W^\alpha) - \frac{1}{6} R \mathcal{E}^4 + \text{fermions}$$

$\Phi$  is conformal comp. :  $(\phi, \psi, F_\phi)$   
 $F_\phi$  <sup>WILL BE</sup> ~~is~~ AUX COMP OF GRAVITY MULTIPLY

$$[A] = 2 \\ R[\phi] = 0 \\ f = -3M_p^2 e^{-k/3m^2}$$

$$(e^a_\mu, \psi_{a\mu}, P_\mu)$$

$\langle \phi \rangle$  BREAKS CONFORMAL SYM

$(F_\phi, P_\mu)$  IS AUX DOF OF GRAV MULTIPLY

$\langle F_\phi \rangle \rightarrow$  VEV IS THE SOURCE OF GRAVITY MEDIATION

STANDARD GAUGE CHOICE :  $\Phi = 1 + \theta^2 F_\phi$

note :  $[F_\phi] = 0$ , convention

$$R[F_\phi] = 2/3$$

→ 'anomaly mediation'  
Randall & Sundrum: out of this world stuff

LET'S SEE WHAT KIND OF SOFT MASSES WE GET

on-shell formalism

← all after canonical norm.

$$(M^2_{\text{soft}})_{\bar{a}\bar{b}} = \underbrace{(F^2 - 2M_{3/2}^2)}_{\substack{\uparrow \\ F^2 = 3M_{3/2}^2}} K_{\bar{a}\bar{b}} - \frac{R_{\bar{a}\bar{b}c\bar{d}} F^c F^{\bar{d}}}{\text{for } \Lambda \sim 0}$$

danger: will give off-diag soft terms (even there in off shell formalism)

↳ need extra ingredients to suppress these

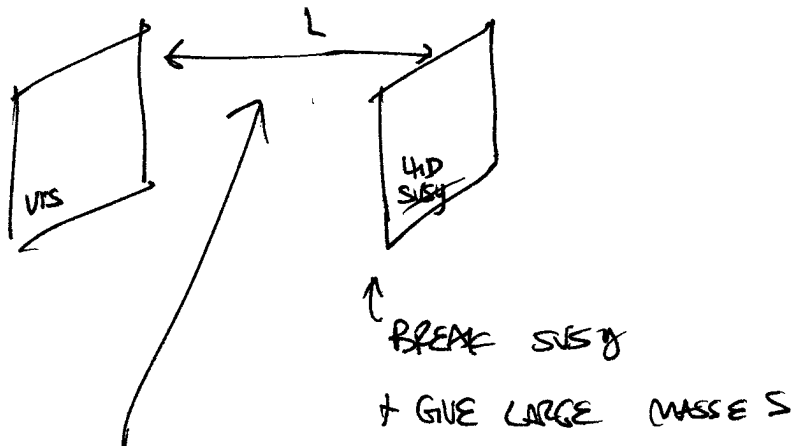
RS: suppose we have  $W = W_{\text{vis}} + W_{\text{hid}}$   
 $\uparrow \qquad \qquad \uparrow$   
 $W_{\text{vis}}(\phi) \qquad V_{\text{vis}}(x)$   
 (SEPARABLE!)

also:  $f = f_{\text{vis}} + f_{\text{hid}}$   
 $\tau = \tau_{\text{vis}} + \tau_{\text{hid}}$

then in  $\frac{1}{\Lambda^2} L$ , VIS & HD SECTORS DECOUPLE.

$\Rightarrow$  SEQUESTERING.

XD picture



$\Rightarrow$  HAS TO OCCUR THROUGH  $F_{\star}$  (gravity mediation)

BUT:  $\exists$  KK MODES  $m \sim \frac{1}{L}$   
 $G_F^{kk}(L) \leftarrow$  GREEN'S FUNK FOR KK MODE

$$\hookrightarrow = e^{-M_{KK} L} \quad \left. \begin{array}{l} \uparrow \\ \frac{1}{L} \end{array} \right\} \text{not small!}$$

PROBLEM!

hand wave = embed into ~~sets~~ strings  
 $\uparrow$  ignore string-scale-ish  
KK modes.

$$\frac{1}{f_g} \mathcal{L} = \int d^4x \left( \frac{1}{2} f(\alpha^2) + \int d^2\theta (\Phi^3 W + \tau W^2) - \frac{1}{6} \tilde{\tau} \tilde{\Phi}^3 + \dots \right)$$

## Soft Masses Explicitly

ASSUME WE HAVE SEQUESTERING

↑ THAT WE'VE INTEGRATED OUT MID SEC.

$$(F_\Phi \neq 0)$$

$$\begin{aligned} \frac{1}{f_g} \mathcal{L} &= |F_\Phi|^2 f(\theta, \bar{\theta}) + F_\Phi \bar{F}_\Phi \hat{\partial}_f^{\text{④}} + \text{c.c.} \\ &\quad + |F_\Phi|^2 \partial_\theta \bar{\partial}_{\bar{\theta}} f + \text{c.c.} \\ &\quad + 3F_\Phi W + F_\Phi \partial_\theta W + \text{c.c.} \end{aligned}$$

②

solve for  $F_g$ :

$$F_g = \left( \frac{\partial \bar{W} + \partial f F_\Phi}{2 \partial f} \right) \rightarrow \text{plug back in}$$

bottom line of sequestering

$$V_{\text{vis}} = \frac{|\partial W + \partial f \bar{F}_\Phi|^2}{2 \partial f} - 6 \text{Re}(W F_\Phi) - f_{\text{vis}} |F_\Phi|^2$$

③ + ④                      ②                      ①

RS: "obviously scalar masses cancel"  
but this isn't obvious or true!

ASSUMED

$$f = g g^{\chi} + (\text{Planck suppressed})$$

$$W = \frac{1}{2} h g^2 + \frac{\lambda}{6} g^3$$

$$V_{\text{vis}} = \left\{ \begin{array}{l} \underbrace{h g + \frac{\lambda}{2} g^2}_{\sim} + \underbrace{g^{\chi} |F_{\phi}|^2}_{\sim} \\ - 6 \text{Re} \left( \left( \frac{1}{2} g^2 + \frac{\lambda}{6} g^3 \right) F_{\phi} \right) \\ - \underline{g g^{\chi} |F_{\phi}|^2} \end{array} \right.$$

underlined terms cancel (RS obs.)

but: wiggly terms also contribute  
+ supersym masses from  $h$  term.

$$= \underbrace{|h|^2 |g|^2 + (h \bar{F}_{\phi} g^2 + \text{c.c.}) - 3 \text{Re}(h g^2 F_{\phi})}_{\text{SUSYING SOFT MASSES}} + \text{SUSY}$$

↳ but these terms don't exist  
for SM fermions

POINT: still no general  $\int$  masses!

So where do the soft masses come from?

### Anomaly Mediation

$\Phi$  couples to all things w/ mass dim.  
 $\xi$  couples equally

so: introducing a cutoff  $\Lambda_{UV}$ ,  
it really comes w/  $(\Lambda_{UV} \xi)$

$\Lambda$  APPEARS IN LOOPS

→  $\Phi$  APPEARS IN LOOPS

→  $F_\Phi$  ———  $\Pi$

so: we've argued that the tree-level terms don't show up for our models of interest

so these loop terms are now the leading terms to soft masses.



why?

breaking out

$$M_{\pm} = \frac{\beta(g)}{2g} F_{\pm} = b_0 \frac{g^2}{16\pi^2} F_{\pm}$$

$$M^2 = \pm \frac{g^4}{(16\pi^2)^2} |F_{\pm}|^2$$

loop factors (flavor indep)

→ come w/  $F_{\pm} \sim M_{\pm}$

Anomaly? → Konishi  
SuperWeg

- UP TO INTERPRETATION.

left  
fact

still other contributions @ loop level  
made some assumptions st. these  
flavor universal terms would dominate

(analogous to tree level terms)

→ eg.  $\Delta M^2 = F_i \partial_i \ln |\det Z| \int \text{Kogut-Popitz} + F_i \partial_i K$

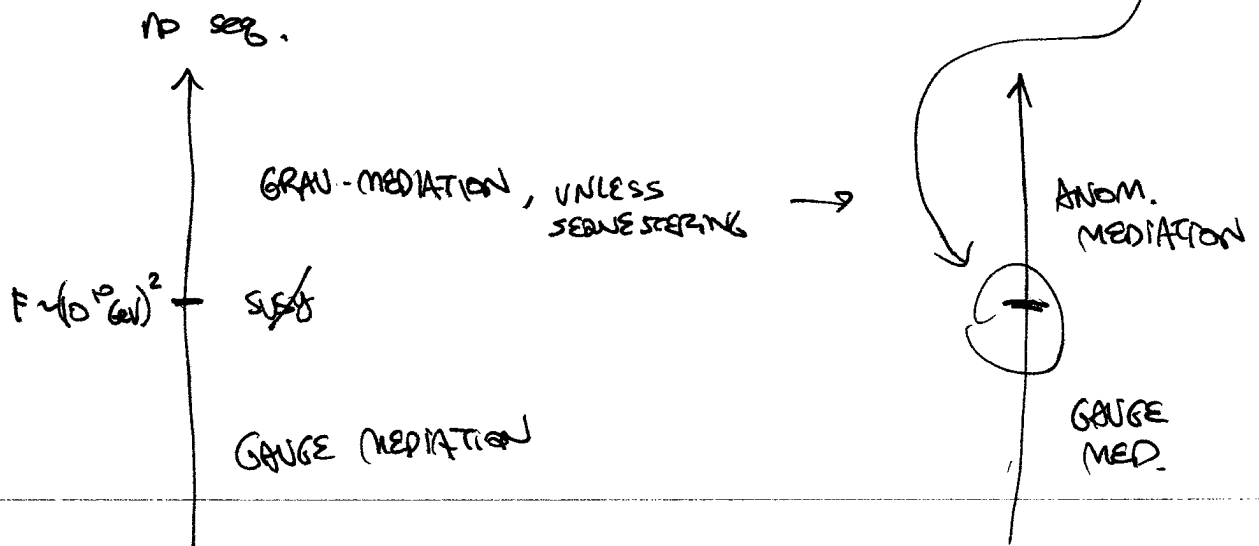
assump. about vevs in hidden sector,  $K_{uv}$

note :  $m^2$  has  $\neq$

IR free  $\rightarrow$  -  
 SEPTONS  $\rightarrow$  tachyonic masses!  $\times$

HARD TO FIX w/M ANOMALY MED.  
 so try to balance tree contributions  
 w/ ~~gravity~~ mediation contributions.  
 gauge

How to get rid of additional KK stuff  
 $\hookrightarrow$  WARPING  $\hookrightarrow$  conformal SEQUESTERING



eg: Anom med gives  $M_X$ , A terms, ...  
 that GAUGE MED DOES NOT.