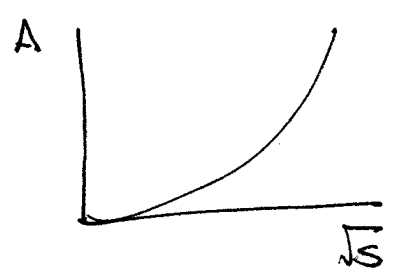


UVI : HIGGS A NGB (COMPOSITE)
 ↑
 not necessarily, but
 req. by naturalness.

$M_{WZ}^2 \neq 0$ ($M^2 \sim g^2 v^2$)
 ↑

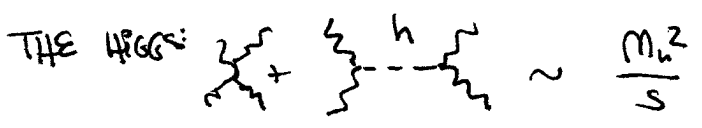
2 TRANSV + LONGITUDINAL
 ↓↓

$W_L W_L \rightarrow W_L W_L$ POORLY BEHAVES @ H⁺ E.



$A \sim \frac{s}{v^2} \leftarrow \sum_{W_L}^2$

for $\sqrt{s} \sim 2 \text{ TeV}$, $A \geq 1$? VIOLATES UNITARITY



↑
 IF THE COUPLINGS ARE CORRECT
 → STANDARD MODEL



SM: $V(H) = m_H^2 |H|^2 + \lambda |H|^4 \rightarrow \langle H \rangle = v$

from EFT perspective: $m_H^2 \sim \epsilon \Lambda_{UV}^2$ w/ $\epsilon \sim \mathcal{O}(1)$

$\epsilon \sim +1$: No EWSB X

$\epsilon \sim -1$: MEXICAN HAT, BUT
 $\langle H \rangle \sim \Lambda_{UV}$ X

HIERARCHY PROBLEM

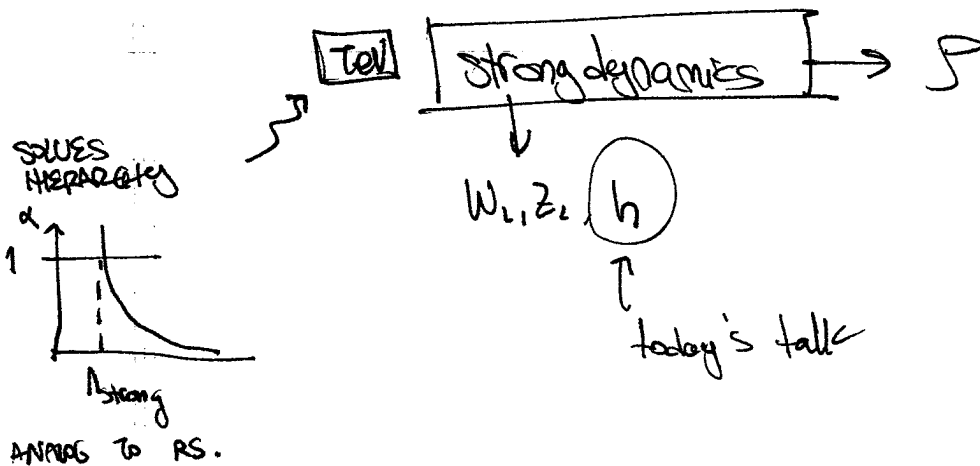
NEED: $\epsilon \sim -10^{-3} \Rightarrow \langle H \rangle = v = 246 \text{ GeV}$
 why is ϵ so small?

SOLUTIONS: SUSY, STRONG DYNAMICS

UNKNOWN DYNAMICS GENERATES $W_L, Z_L + (?)$

OLD DAYS: thing which unitarizes $W_L W_L$ scattering
 can be anything: eg ρ -like vector meson

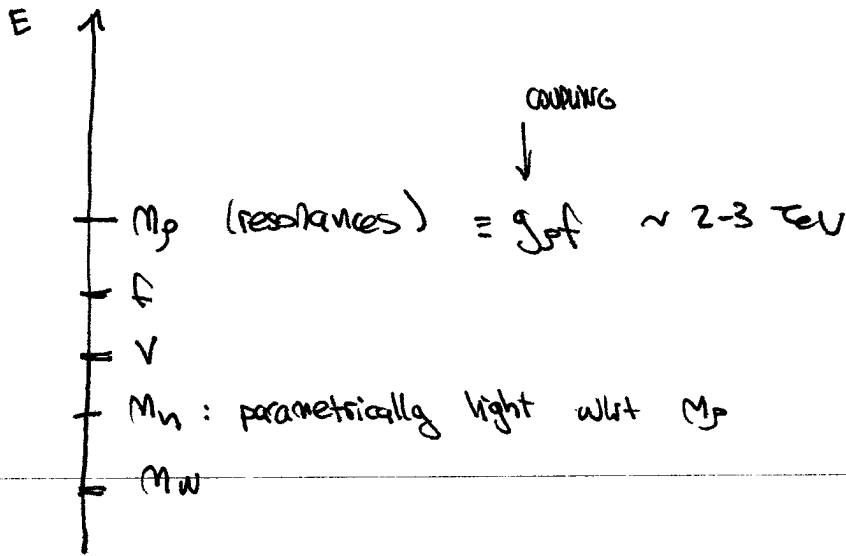
NOW: WE NEED A HIGGS (for 125 GeV state)



Remark: strong dynamics \neq XD
are not necessarily complementary P's.

THE HIGGS IN STRONG DYNAMICS

↑
LIGHT: NEED PROTECTION MECHANISM \rightarrow GOLDSTONE'S THEM
SHIFT SYMMETRY: $h \rightarrow h + \alpha$; eg MASS TERM FORBIDDEN



GOLDSTONES: $U = e^{i\pi/f}$ ← characteristic scale

recall: $\sim \frac{\partial^2}{v^2} \ll v$ HERE IS CHARAC. SCALE

Now for us: STRONG SECTOR GENERATES $\frac{H}{f} \leftarrow w, z, h$

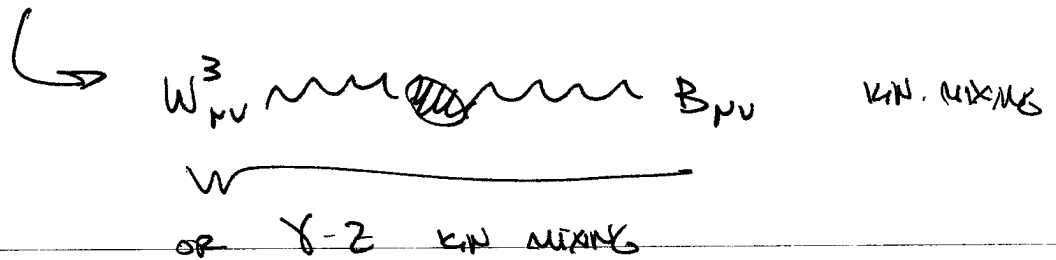
$$V(H) \Rightarrow \langle H \rangle = v \ll f$$

MUCH BETTER THAN TECHNICOLOR
SM OBSERVABLES CORRECTED
@ ORDER v^2/f^2

TECHNICOLOR: $v = f$
SM: $v \ll f$

TECHNICOLOR WAS ALREADY IN BAD SHAPE BEFORE LHC:

S-PARAMETER



$\hat{S} \sim \frac{g^2}{g_0^2} \left(\frac{v^2}{f^2} \right)$ ← 1 IN TECHNICOLOR

$\hat{S}_{exp} < 2 \times 10^{-3}$

→ TECHN. RES. $g \sim 4\pi$ → strongly coupled
calculable.

for us: $\Rightarrow \boxed{\frac{v^2}{f^2} \sim 0.1 - 0.3}$

this is our tuning.

ALL OBSERVABLES ARE MODIFIED BY $\Theta(v^2/\Lambda^2)$

→ IN PARTICULAR: HIGGS COUPLINGS

so: PREDICTIONS FOR HIGGS AS NGB (recent work)

$$\begin{aligned} \mathcal{L}_{\text{eff}} = & \frac{1}{2} (\partial h)^2 + \frac{M_W^2}{2} W W \left[1 + 2a \frac{h}{v} + b \frac{h^2}{v^2} \right] \\ & - m_i \bar{\psi}_i \left(1 + c \frac{h}{v} \right) \psi_i \\ & + \frac{1}{2} m_h^2 h^2 + d_3 \frac{1}{6} \left(\frac{3m_h^2}{v} \right) h^3 + d_4 \frac{1}{24} \left(\frac{3m_h^2}{v^2} \right) h^4 \\ & + c_g \frac{\alpha_s}{4\pi} \frac{h}{v} G_{\mu\nu}^2 + c_\gamma \frac{\alpha}{4\pi} \frac{h}{v} F_{\mu\nu}^2 \end{aligned}$$

SM: $a = b_1^c = d_3 = d_4 = 1$ $c_g = c_\gamma = 0$

NGB: $a, b_1^c, d_3, d_4 \sim \Theta(1)$ $c_g = c_\gamma \sim \frac{\alpha_i}{4\pi}$

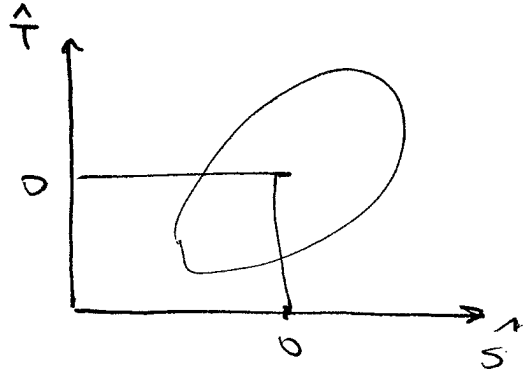
↳

the game: fit these parameters to LHC

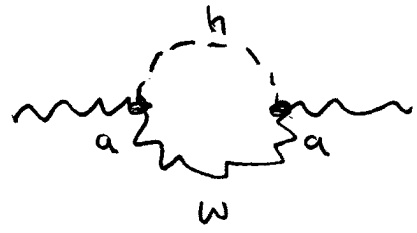
EXPERIMENT: a HAS BEEN MEASURED INDIRECTLY

\hat{T}	$W \sim W W - Z \sim Z$	MASS DIFFERENCE
\hat{S}	$\gamma \sim \gamma$	KIN. MIXING

LEP:

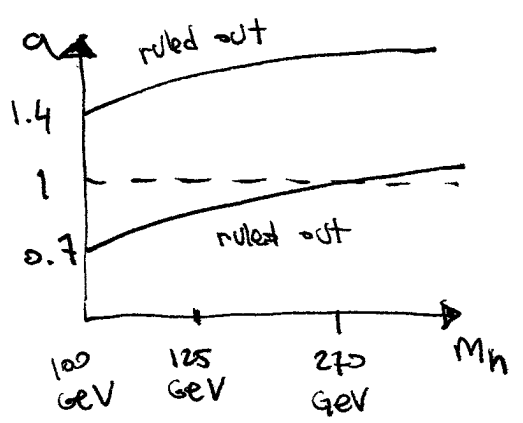
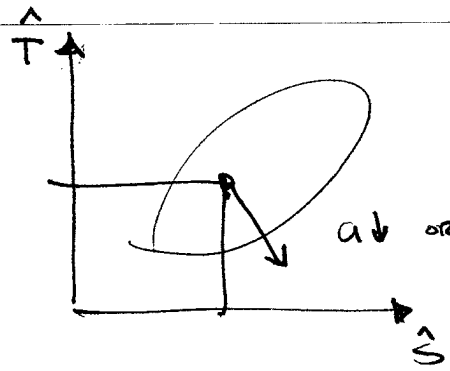


\hat{S} & \hat{T} ~~are~~ ARE CORRECTED BY HIGGS LOOPS:



$$\Delta \hat{T}, \Delta \hat{S} |_h \sim \frac{\#}{16\pi^2} (1-a^2) \log\left(\frac{\Lambda^2}{m_h^2}\right)$$

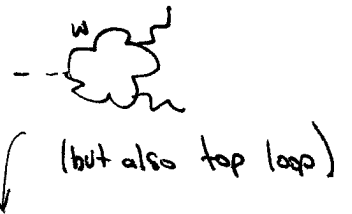
TRICK: SM $\Rightarrow a=1$



PREDICTIONS FROM $SO(5)/SO(4)$ COMPOSITE HIGGS MODEL

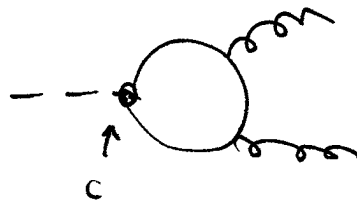
$$a = \sqrt{1 - \frac{v^2}{f^2}} \quad b = 1 - 2 \frac{v^2}{f^2} \quad \left. \vphantom{\frac{v^2}{f^2}} \right\} \text{kind of model independent!}$$

$$\textcircled{\text{IV}} \quad \left\{ \begin{array}{l} c = d_3 = \sqrt{1 - \frac{v^2}{f^2}} \\ \textcircled{\text{V}} \quad c = d_3 = \frac{1 - 2v^2/f^2}{\sqrt{1 - v^2/f^2}} \end{array} \right.$$



MEASURE : $\frac{\Gamma(h \rightarrow \nu\nu)}{\Gamma_{SM}(h \rightarrow \nu\nu)} = a^2 \quad \sim \frac{\Gamma(h \rightarrow \gamma\gamma)}{\Gamma_{SM}(h \rightarrow \gamma\gamma)}$

$$\frac{\Gamma(h \rightarrow gg)}{\Gamma_{SM}(h \rightarrow gg)} = \frac{\Gamma(h \rightarrow t\bar{t})}{\Gamma_{SM}(h \rightarrow t\bar{t})} = c_t^2$$



\neq when $c_g \neq 0$

$$\frac{\Gamma(h \rightarrow f\bar{f})}{\Gamma_{SM}(h \rightarrow f\bar{f})} = c_f^2$$

OBSERVE:

$$a < 1$$

$h \rightarrow \gamma\bar{\gamma}$ REDUCED

(but can push up by
reducing $h \rightarrow b\bar{b}$)

$$d < 1$$

$h \rightarrow t\bar{t}$ REDUCED

↑
"now it's a game"

Dean's 2: COMPOSITE ZHM w/ NGB HIGGS?
yes - Jui has done the theory
but not the phenomenology

REMARK ON BREAKING PATTERNS:

$so(5)/so(4)$: GIVES W_L, Z_L, h
(minimal composite higgs)

HIGGS AS A NGB

$$G \rightarrow H$$

$$NGB = G/H$$

REQUIREMENTS :

1. $G > SU(2)_L \times U(1)_Y$
2. $H > \underbrace{so(4) \cong SU(2) \times SU(2)}_W$

AVOIDS TREE-LEVEL
CORRECTIONS TO \uparrow

HIGGS VEV: $SU(2)_L \times SU(2)_R \xrightarrow{\langle H \rangle} SU(2) \cong SO(3)$

$H = \begin{pmatrix} \vdots \\ \vdots \\ \vdots \end{pmatrix} \} SO(4)$

$\langle H \rangle = \begin{pmatrix} 0 \\ 0 \\ 0 \\ v \end{pmatrix} \} \left. \begin{array}{l} \text{ROTATIONS OF} \\ W = \begin{pmatrix} W^1 \\ W^2 \\ W^3 \end{pmatrix} \end{array} \right\}$

3) $G/H \supset (\underline{2}, \underline{2})$ of $SU(2)_L \times SU(2)_R$
 $\underbrace{\hspace{1cm}}_{\text{quantum \#s of HIGGS}}$

WE HAVE $SU(2)_L \times SU(2)_R$ IN SM:
 WRITE HIGGS AS 4-VECTOR ϕ
 $V = m_H^2 \vec{\phi} \cdot \vec{\phi} + \lambda (\vec{\phi} \cdot \vec{\phi})^2$

THE SMALLEST COSETS THAT CAN HAVE HIGGS AS NGB

	G	H	N_G	rep [SU(2) _L × SU(2) _R]
worked out	SO(5)	SO(4)	4	(2, 2) ← HIGGS
being worked on	SO(6)	SO(5)	5	(1, 1) + (2, 2) ↖ extra singlet (easy to hide)
only theory, no analysis	SO(6)	SO(4) × SO(2)	8	2 × (2, 2) 2HDM
	SO(7)	G ₂	7	(1, 3) + (2, 2) ↑ easier to hide

WHICH OF THESE CAN WE GET LIKE QCD?

↑ sym of techniquarks
ie breaking by fermions

MINIMUM WHICH CONTAINS A HIGGS

techniquark w/ 4 flavors: (SU(4))_F

$$\tilde{\psi}_\alpha = \begin{pmatrix} y \\ x \\ x \\ x \end{pmatrix}_\alpha$$

↑ strongly coupled → HAS TO BE sp(N)

technig-technig
BOUND STATE
↓
CF QCD
WHICH IS
β-outig
BOUND
STATE

~~QCD~~ $\langle \tilde{\psi}_\alpha \tilde{\psi}_\beta \rangle \neq 0 \Rightarrow$

$$\begin{matrix} \text{SU}(4) & \rightarrow & \text{Sp}(4) \\ \parallel 15 & & \parallel 15 \\ \text{SO}(6) & \rightarrow & \text{SO}(5) \end{matrix}$$

$\int_{\alpha\beta} [\text{PAIR MATRIX}]$

↓
minimal model in terms
of techniquarks

RECENT PAPER: COMPOSITE SCALAR DM ψ

→ $SO(6)/SO(5)$ MODEL whose singlet is DM

↳ COUPLINGS: $\left(\frac{W_L \partial W_L \psi \partial \psi}{f^2} \right)$

PARTLY REQ TWO ψ 'S.

BUT NEED TO WORRY THAT STRONG
SECTOR RESPECTING THE PARITY

↳ $SO(6)/O(5)$
RESPECTS PARITY

EXPLICIT
BREAKING

COUPLING TO FERMIONS: $\bar{\psi}_L H \psi_R \frac{\eta^2}{f^2} \rightarrow$ controls
PHENOM.

ANOMALIES? $\psi G_{\mu\nu} \tilde{G}^{\mu\nu}$

↳ like QCD: $\alpha \ll 1$, "strong CP problem"
can be done, (can make it zero)



SCALES w/ E
SO CAN HOPE
FOR SIGNAL @ 14 TeV

THESE KINDS OF COUPLINGS PROBE COMPOSITENESS OF ψ .

Cuidado:

the point is that once you embed your nice model into strong dynamics (as in sisg) there are predictions for how the model must behave.

↳ analogous to sisg \Rightarrow zham.
