

## Phase Measurements

Suppose a particle is offset at a point in the ring at  $(x_0, y_0)$ . This displacement will propagate as:

$$x_{ij} = A_i \cos(j2\pi Q_x + \phi_x^i) \quad (1)$$

$$y_{ij} = B_i \cos(j2\pi Q_y + \phi_y^i) \quad (2)$$

where the scripts  $ij$  indicate the  $i^{th}$  BPM on the  $j^{th}$  turn. We assume a large number of turns ( $\sim 1000$  or more) are used.

Define  $C_x^i, S_x^i$  to be:

$$C_x^i = \sum_{j=1}^N x_{ij} \cos(j2\pi Q_x) \quad (3)$$

$$S_x^i = \sum_{j=1}^N x_{ij} \sin(j2\pi Q_x) \quad (4)$$

with similar definitions for  $C_y^i, S_y^i$ . Further inspect the expression for  $C_x^i$ :

$$\begin{aligned} C_x^i &= \sum_{j=1}^N A_i \cos(j2\pi Q_x + \phi_x^i) \cos(j2\pi Q_x) \\ &= A_i \sum_{j=1}^N \left( \cos(j2\pi Q_x) \cos \phi_x^i - \sin(j2\pi Q_x) \sin \phi_x^i \right) \cos(j2\pi Q_x) \\ &= A_i \sum_{j=1}^N \left( \cos^2(j2\pi Q_x) \cos \phi_x^i - \sin(j2\pi Q_x) \cos(j2\pi Q_x) \sin \phi_x^i \right) \\ &= A_i \sum_{j=1}^N \left( \frac{1}{2} (1 + \cos(j4\pi Q_x)) \cos \phi_x^i - \frac{1}{2} \sin(j4\pi Q_x) \sin \phi_x^i \right) \\ &= \frac{1}{2} A_i \left( N \cos(\phi_x^i) + \sum_{j=1}^N \left( \cos(j4\pi Q_x) \cos \phi_x^i - \sin(j4\pi Q_x) \sin \phi_x^i \right) \right) \\ &= \frac{1}{2} A_i N \left( \cos(\phi_x^i) + \frac{1}{N} \sum_{j=1}^N \cos(j4\pi Q_x + \phi_x^i) \right) \end{aligned}$$

Cosine is constrained to be  $\in (-1, 1)$  and we explicitly choose the tunes  $Q$  such that we are not on-resonance, therefore as  $N \rightarrow$  very large, the remaining sum's contribution to  $C_x^i$  goes to zero. Therefore,

$$C_x^i \approx \frac{1}{2} A_i N \cos(\phi_x^i) \quad (5)$$

We attain similar expressions for  $S_x^i, C_y^i$ , and  $S_y^i$ , where the  $S$  terms are negative. Finding the phase advance is now relatively trivial:

$$\phi_x^i = \tan^{-1} \left( \frac{S_x^i}{C_x^i} \right) \quad (6)$$

$$\phi_y^i = \tan^{-1} \left( \frac{S_y^i}{C_y^i} \right) \quad (7)$$

Note that the arctan values returned by Fortran are constrained to be  $\in (-\pi, \pi)$ . In order to correctly propagate the phase, we must add factors of  $2\pi$  appropriately whenever the phase at a BPM “rolls over” from  $\pi$  to  $-\pi$ .

This only leaves an ambiguity in a constant phase offset from the design values. This can be resolved in the same way CESRv resolves the problem: find the average difference between the “measurement” and the design, and add this to the measured values. This sets the average difference between measured and design to zero.

## Coupling Measurements

Although everything in the above section is technically correct, it only holds for the zero-coupling scenario. In the more general case, we have to also include terms arising from coupling:

$$x_{ij} = A_a^i \cos(j2\pi Q_x + \phi_{ax}^i) + A_b^i \cos(j2\pi Q_y + \phi_{bx}^i) \quad (8)$$

$$y_{ij} = B_b^i \cos(j2\pi Q_y + \phi_{by}^i) + B_a^i \cos(j2\pi Q_x + \phi_{ay}^i) \quad (9)$$

We must now introduce new sums  $(C_{ax}^i, S_{ax}^i, C_{ay}^i, S_{ay}^i, C_{bx}^i, S_{bx}^i, C_{by}^i, S_{by}^i)$  to fully describe the system. These sums are described in the same fashion as in equations (3,4), and the results for  $C_{ax}^i, S_{ax}^i, C_{by}^i, S_{by}^i$  are the same as those we originally found. The new cross-terms are what will allow us to find the coupling measurements.

The  $\bar{C}$  elements of interest are defined by:

$$\bar{C}_{12} = \frac{B_a}{A_a} \sin(\phi_{ay} - \phi_{ax}) \quad (10)$$

$$\bar{C}_{22} = \frac{B_a}{A_a} \cos(\phi_{ay} - \phi_{ax}) \quad (11)$$

$$\bar{C}_{11} = \frac{A_b}{B_b} \cos(\phi_{bx} - \phi_{by}) \quad (12)$$

We can similarly use the b-mode measurements to define  $\bar{C}_{12}$ . To obtain the amplitudes  $A, B$  from our measurement, consider the sums  $C_{ax}, S_{ax}$ :

$$\sqrt{C_{ax}^2 + S_{ax}^2} = \frac{N}{2} A_a^i \quad (13)$$

We can repeat this method to find all four amplitudes.